

# Timelike Compton Scattering

I. Albayrak<sup>1</sup>, V. Burkert<sup>2</sup>, E. Chudakov<sup>2</sup>, M. Guidal<sup>3,\*</sup>, V. Guzey<sup>2</sup>, K. Hicks<sup>4</sup>, T. Horn<sup>1,\*</sup>, C. Hyde<sup>5</sup>, Y. Ilieva<sup>6</sup>, F.J. Klein<sup>2</sup>, C. Munoz Camacho<sup>3</sup>, P. Nadel-Turonski<sup>2,\*\*</sup>, M. Osipenko<sup>7</sup>, R. Paremuzyan<sup>3,8</sup>, B. Pire<sup>9</sup>, F. Sabatie<sup>10</sup>, C. Salgado<sup>11</sup>, S. Stepanyan<sup>2,\*</sup>, L. Szymanowski<sup>12</sup>, J. Wagner<sup>12</sup>, C. Weiss<sup>2</sup>.

\* spokespersons

\*\* contact person

- 1) The Catholic University of America, Washington, DC 20064
- 2) Thomas Jefferson National Accelerator Facility, Newport News, VA 23606
- 3) Institut de Physique Nucleaire d'Orsay, IN2P3, BP 1, 91406 Orsay, France
- 4) Ohio University, Athens, OH 45701
- 5) Old Dominion University, Norfolk, VA 23529
- 6) University of South Carolina, Columbia, SC 29208
- 7) INFN, Sezione di Genova e Dipartimento di Fisica dell'Universita, 16146 Genova, Italy
- 8) Yerevan Physics Institute, 375036 Yerevan, Armenia
- 9) CPhT Ecole Polytechnique, F 91128 Palaiseau CEDEX, France
- 10) CEA, Centre de Saclay, Irfu/Service de Physique Nucleaire 91191 Gif-sur-Yvette, France
- 11) Norfolk State University, Norfolk, VA 23504
- 12) National Center for Nuclear Research, Warsaw, Poland

# Timelike Compton Scattering, or timelike DVCS

$$\gamma p \rightarrow \gamma^* p \rightarrow e^+ e^- p$$

- Important addition to the CLAS GPD program at 12 GeV
- Complementary to (spacelike) DVCS
- Will test universality of factorization, which makes it possible to probe GPDs in exclusive processes
- Excellent experimental tool for accessing the real part

# LOI11-106: $e^+e^-$ pair production with CLAS12 at 11 GeV

- Physics objectives: Timelike Compton (DVCS),  $J/\Psi$  on nucleon and nuclei

## **PAC recommendation**

“The physics addressed in this proposal is very relevant for the JLab 12 GeV program. The PAC encourages the development of a full proposal.”

## **PAC 39 Proposal: Timelike Compton Scattering**

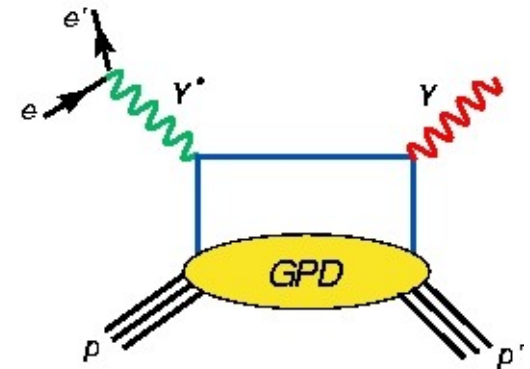
- Important contribution to the GPD program
- Theory rapidly developing – much progress
- $J/\Psi$  cross section will be measured ( $\rightarrow$  normalization)

## **$J/\Psi$ : will be submitted as separate proposal(s)**

- Developing physics case: Gluonic structure of nucleon, non-perturbative color forces.
- Workshop Temple University, March 26-28, 2012
  - <http://quarks.temple.edu/~npcfiqcd>

# Compton scattering – overview

- Real Compton Scattering
- Deeply Virtual Compton Scattering (DVCS)
  - Outgoing photon is real
  - Simplest probe of GPDs
- Timelike Compton Scattering (TCS)
  - Incoming photon is real
  - Complementary to DVCS
- Double DVCS
  - Both photons are virtual
  - Experimentally challenging



$$\gamma^* + p \rightarrow \gamma^* + p$$

GPDs can be extracted from Helicity Amplitudes or Compton Form Factors

# Why TCS?

$$\gamma p \rightarrow \gamma^* p \rightarrow e^+ e^- p$$

## Theory

- Straightforward access to real part of amplitudes/CFFs/GPDs
- Universality of GPDs extracted from exclusive processes
  - Spacelike – Timelike comparison as for DIS and Drell-Yan
- Dispersion relations – how important is the real part at large  $x$ ?
  - Model-independent analysis
- Impact on global fits for Compton Form Factors (Guidal, Sabatie)
- Interesting behavior of NLO corrections
  - Sensitive to GPD models and potentially to gluons

## Experiment

- Challenging to get real part from beam charge asymmetries (electron – positron)
  - And currently no such facility

# Probing GPDs through Compton scattering

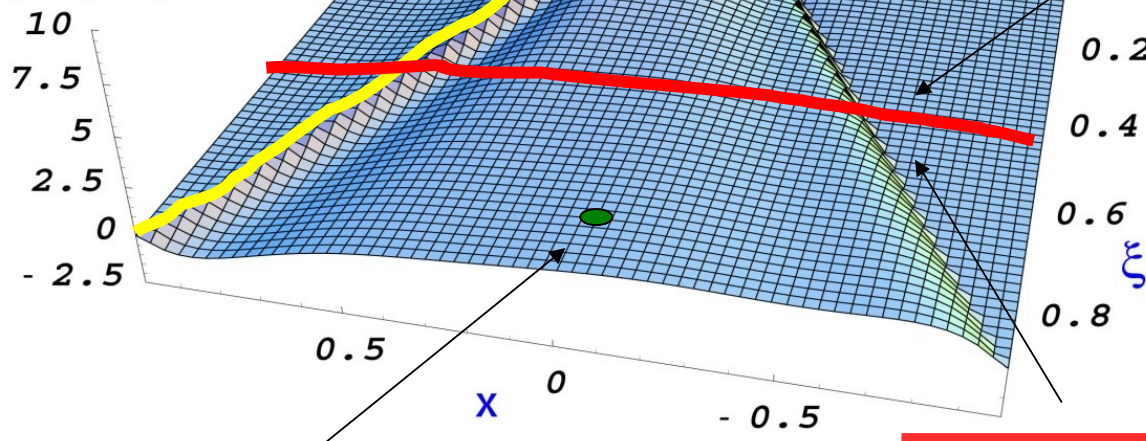
$(\text{Im}, x=\xi)$

DVCS: spin asymmetries  
HERMES, CLAS, Hall A

$(|\text{Re}|)$

TCS: azimuthal asymmetry  
CLAS  
DVCS: charge asymmetry  
HERMES

$H(x, \xi, 0)$



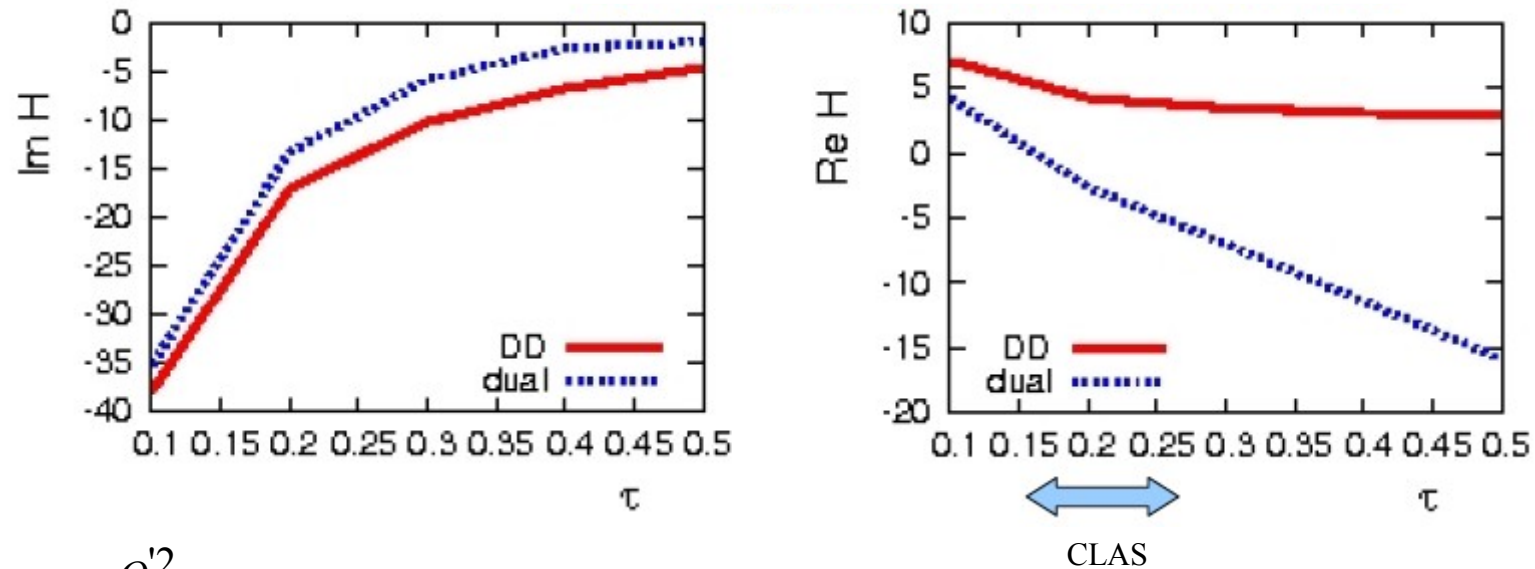
$(\text{Im}, x \neq \xi, x < |\xi|)$

DDVCS

$(|\text{Re}|^2)$

DVCS: cross sections  
H1, Hall A

# GPD models sensitive to real part at large x

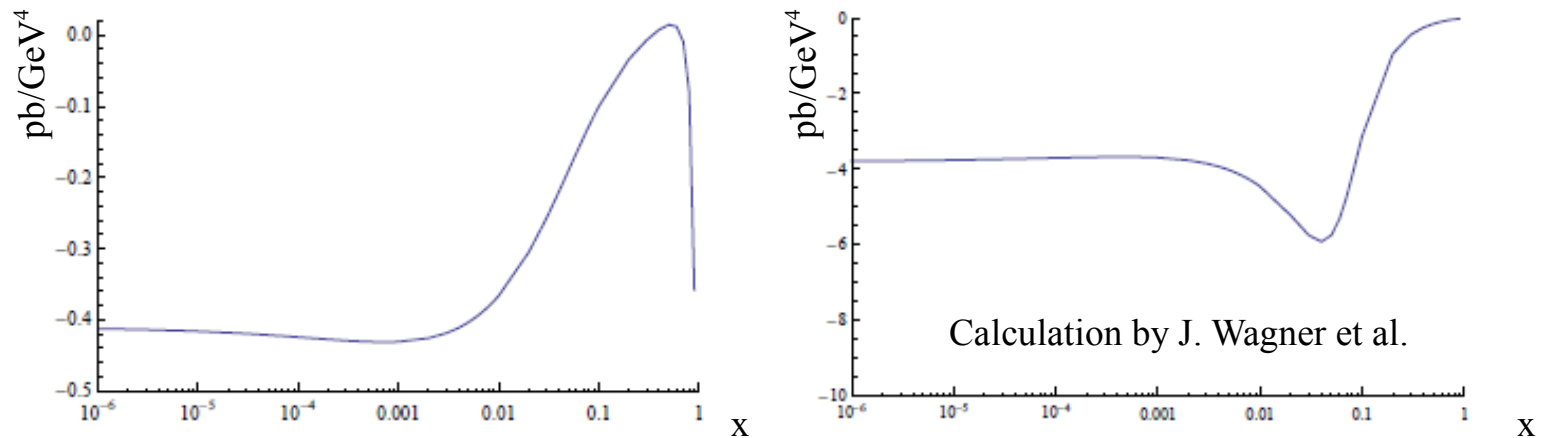


$\tau = \frac{Q'^2}{s - M_p^2}$  is the equivalent of Bjorken  $x$ ,  
hard scale is given by  $Q^2 = M_{e^+e^-}^2$

LO cCalculation by V. Guzey  
for  $Q^2 = 5 \text{ GeV}^2$  and  $t = 0$

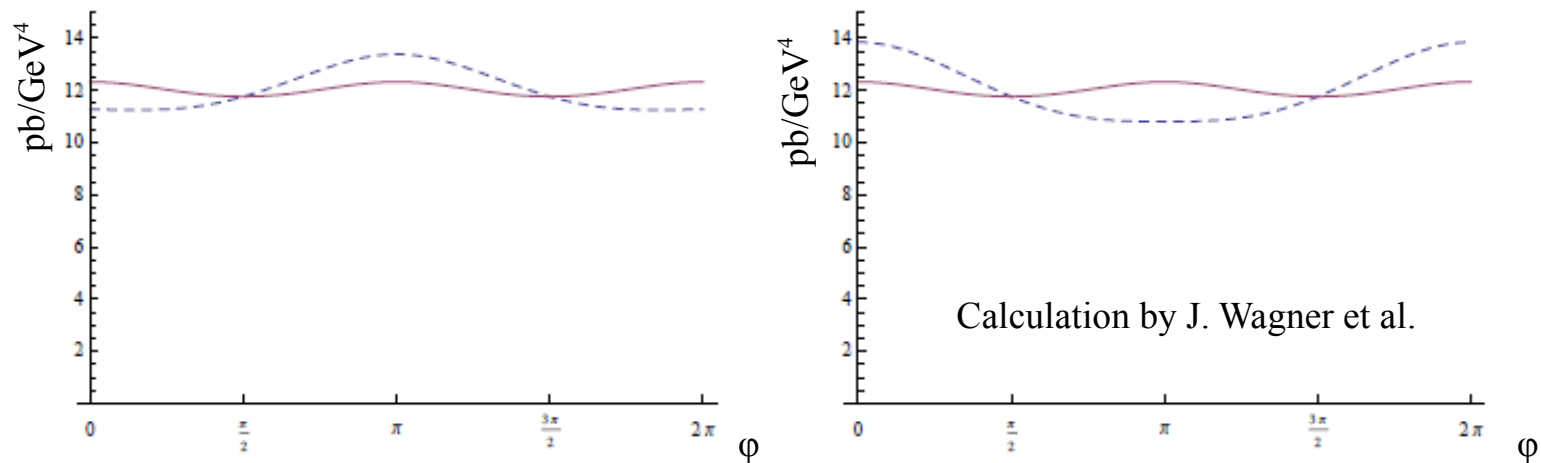
- Model predictions similar for  $\text{Im } H$ , but large differences for  $\text{Re } H$
- Reliable measurements of real part are needed!

# NLO corrections – a crucial test for theory



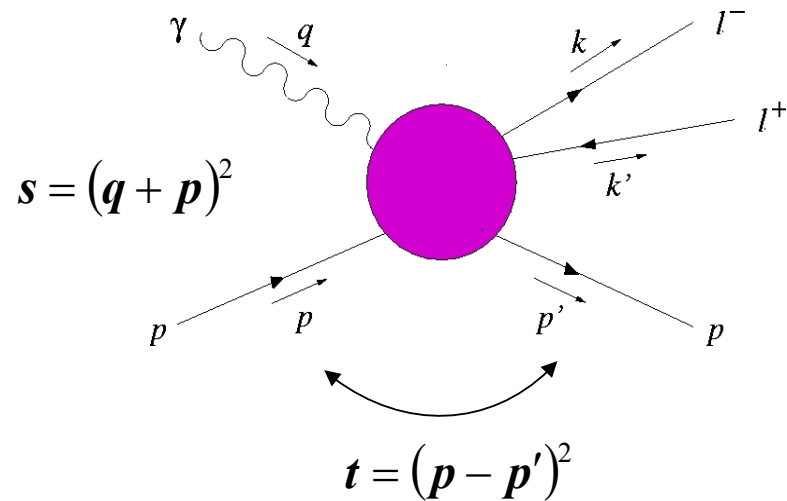
Ratio of NLO correction to the Born term for the imaginary (left) and real (right) part of the CFF H for the Kroll-Goloskokov model, for  $Q^2 = 4 \text{ GeV}^2$ ,  $t = 0$ , and  $\mu_F = Q$ .

The correction on the right is almost entirely due to gluons.

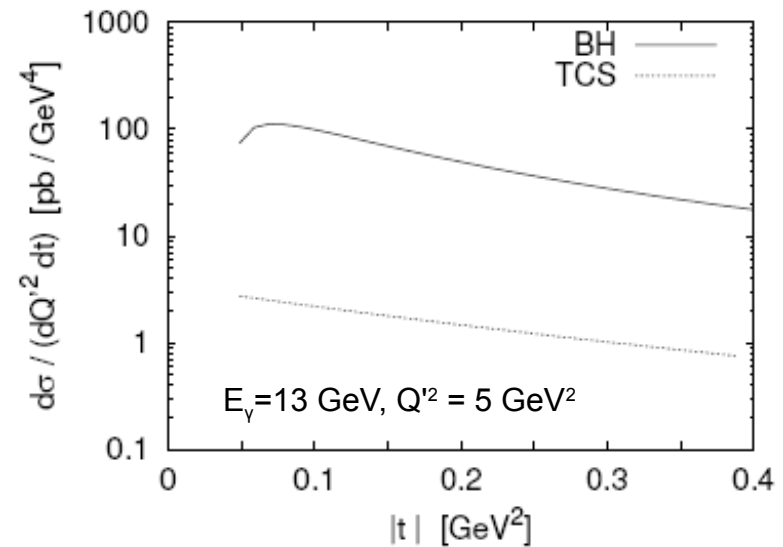


LO (left) and NLO (right) differential cross sections for B-H (solid) and B-H + INT (dashed) for  $E_\gamma = 11 \text{ GeV}$ ,  $Q^2 = 5 \text{ GeV}^2$ , and  $t = -0.1 \text{ GeV}^2$ .

# Photoproduction of lepton pairs

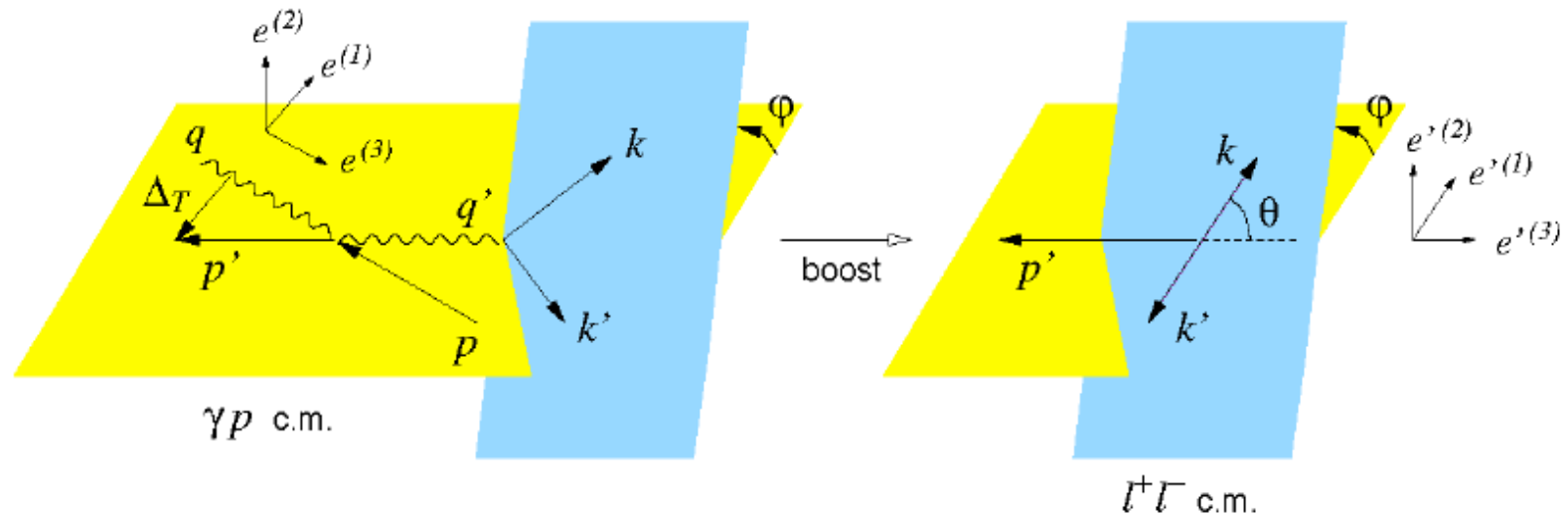


E. Berger *et al.*, Eur. Phys. J. C23, 675 (2002)



- TCS cross section is small compared with Bethe-Heitler (B-H) for all kinematics
  - cannot be accessed directly
- The interference term is, however, enhanced by the B-H and easy to isolate

# Kinematics

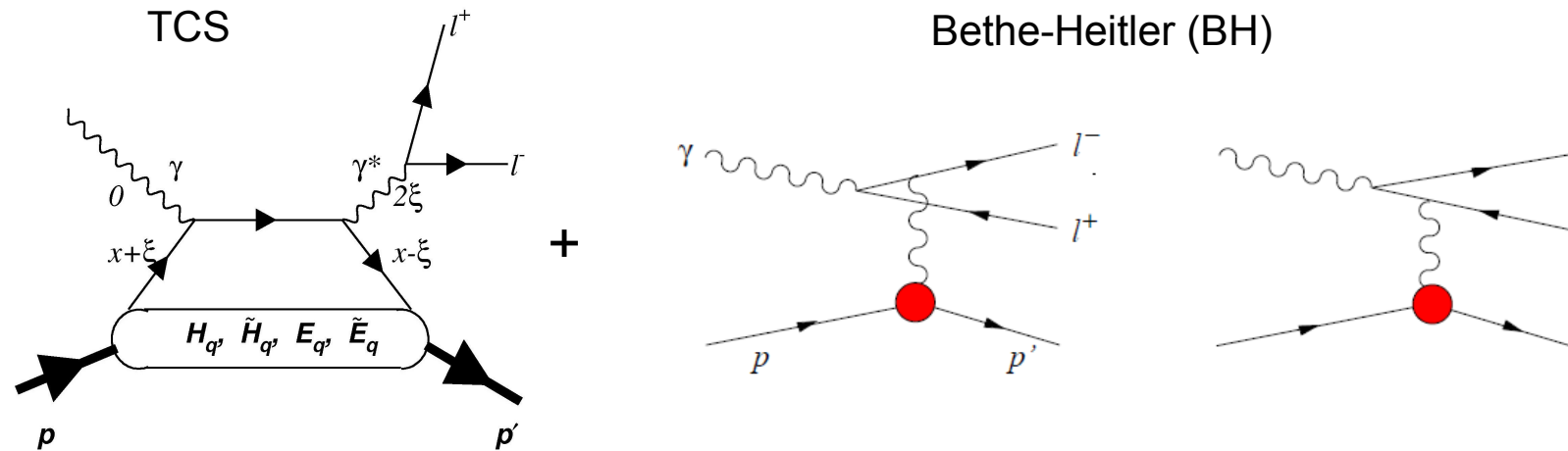


- $k, k' =$  momentum of  $e^-, e^+$
- $\theta =$  angle between the scattered proton and the electron
- $\phi =$  angle between lepton scattering- and reaction planes

$$\frac{d\sigma_{BH}}{dQ'^2 dt d\cos\theta} \approx 2\alpha^3 \frac{1}{-tQ'^4} \frac{1 + \cos^2\theta}{1 - \cos^2\theta} \left( F_1(t)^2 - \frac{t}{4M_p^2} F_2(t)^2 \right)$$

- For small  $\theta$ , B-H becomes large. A cut is usually applied.

# Observables sensitive to the interference term



$$\frac{d^4\sigma}{dx_B dQ^2 dt d\varphi} \approx |T^{BH}|^2 + 2T^{BH} \bullet \text{Re}(T^{VCS}) + |T^{VCS}|^2$$

- Under lepton charge conjugation:
  - Compton and BH amplitudes are *even*
  - Interference term is *odd*
  - Observables that change sign project out *only the interference term*
- Example of observable: azimuthal angular distribution ( $\varphi$ ) of the lepton pair

# TCS cross section and the interference term

$$\frac{d\sigma_{TCS}}{dQ'^2 d\Omega dt} \approx \frac{\alpha^3}{8\pi} \frac{1}{s^2} \frac{1}{Q'^2} \left( \frac{1 + \cos^2 \theta}{4} \right) 2(1 - \xi^2) |\mathcal{H}(\xi, t)|^2$$

$$\frac{d\sigma_{INT}}{dQ'^2 dt d\cos\theta d\varphi} = - \frac{\alpha_{em}^3}{4\pi s^2} \frac{1}{-t} \frac{M}{Q'} \frac{1}{\tau\sqrt{1-\tau}} \cos\varphi \frac{1 + \cos^2 \theta}{\sin\theta} \text{Re } \tilde{M}^{--}$$

$$\tilde{M}^{--} \approx \frac{2\sqrt{t_0 - t}}{M} \frac{1 - \xi}{1 + \xi} [F_1(t)\mathcal{H}(\xi, t)]$$

$$\mathcal{H}(\xi, t) = \sum_q e_q^2 \int_{-1}^1 dx \left( \frac{1}{\xi - x + i\epsilon} - \frac{1}{\xi + x + i\epsilon} \right) H^q(x, \xi, t)$$

# Full interference term with polarized beams

To leading order, in terms of helicity amplitudes:

$$\begin{aligned} \frac{d\sigma_{INT}}{dQ'^2 dt d(\cos\theta) d\varphi} = & -\frac{\alpha_{em}^3}{4\pi s^2} \frac{1}{-t} \frac{M}{Q'} \frac{1}{\tau\sqrt{1-\tau}} \frac{L_0}{L} \left[ \cos\varphi \frac{1+\cos^2\theta}{\sin\theta} \text{Re } \tilde{M}^{--} \right. \\ & \left. - \cos 2\varphi \sqrt{2} \cos\theta \text{Re } \tilde{M}^{0-} + \cos 3\varphi \sin\theta \text{Re } \tilde{M}^{+-} + O\left(\frac{1}{Q'}\right) \right], \\ & -\nu \frac{\alpha_{em}^3}{4\pi s^2} \frac{1}{-t} \frac{M}{Q'} \frac{1}{\tau\sqrt{1-\tau}} \frac{L_0}{L} \left[ \sin\varphi \frac{1+\cos^2\theta}{\sin\theta} \text{Im } \tilde{M}^{--} \right. \\ & \left. - \sin 2\varphi \sqrt{2} \cos\theta \text{Im } \tilde{M}^{0-} + \sin 3\varphi \sin\theta \text{Im } \tilde{M}^{+-} + O\left(\frac{1}{Q'}\right) \right] \end{aligned}$$

$\nu$ : circular polarization of incoming photon also gives access to imaginary part

$$\begin{aligned} \frac{1}{2} \sum_{\lambda, \lambda'} |M^{\lambda', \lambda-}|^2 = & (1 - \eta^2) (|\mathcal{H}_1|^2 + |\tilde{\mathcal{H}}_1|^2) - 2\eta^2 \text{Re}(\mathcal{H}_1^* \mathcal{E}_1 + \tilde{\mathcal{H}}_1^* \tilde{\mathcal{E}}_1) \\ & - \left(\eta^2 + \frac{t}{4M^2}\right) |\mathcal{E}_1|^2 - \eta^2 \frac{t}{4M^2} |\tilde{\mathcal{E}}_1|^2, \end{aligned}$$

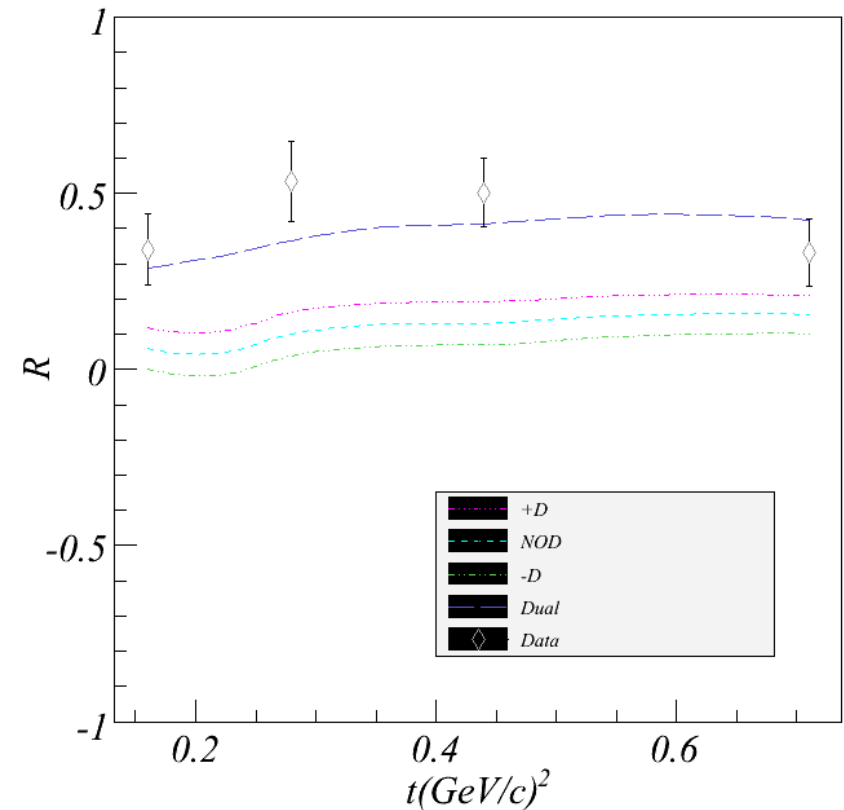
# Example: first data from 6 GeV

- Cosine moment of weighted cross sections

$$\frac{dS}{dQ'^2 dtd\varphi} = \int \frac{L(\theta, \varphi)}{L_0(\theta)} \frac{d\sigma}{dQ'^2 dtd\varphi d\theta} d\theta$$

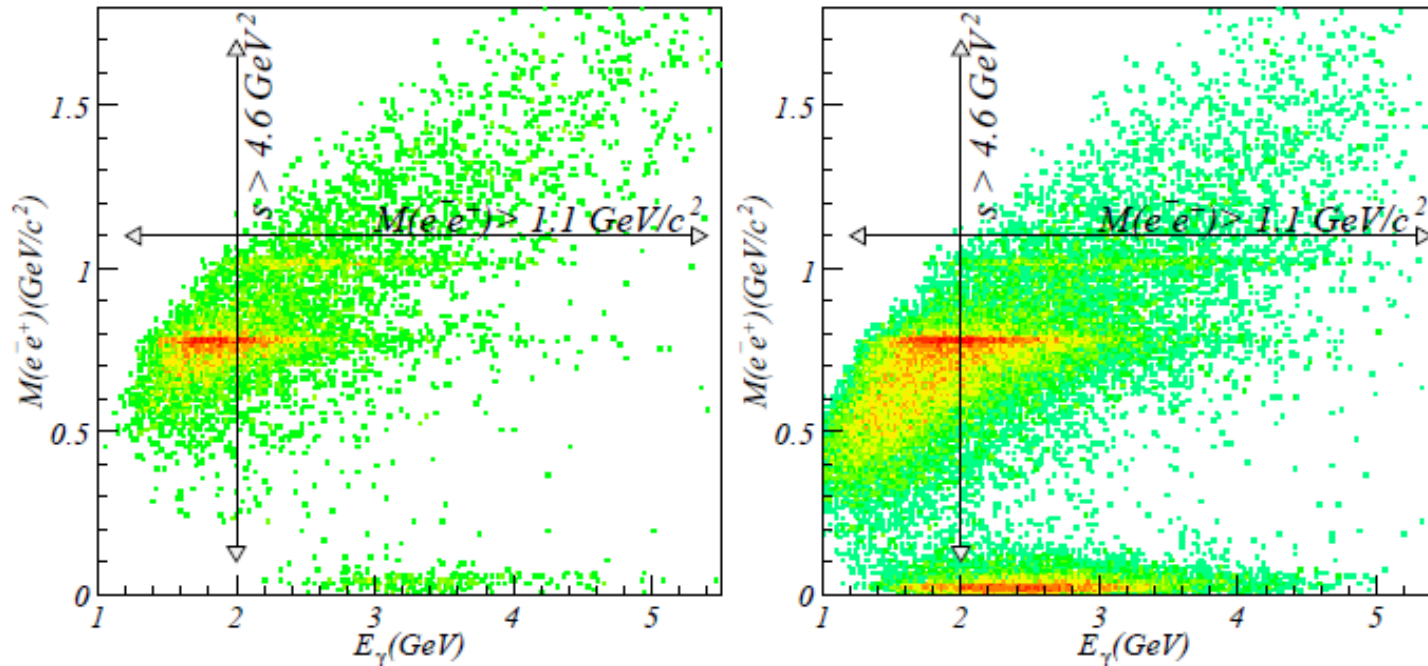
$$R = \frac{2 \int_0^{2\pi} d\varphi \cos \varphi \frac{dS}{dQ'^2 dtd\varphi}}{\int_0^{2\pi} d\varphi \frac{dS}{dQ'^2 dtd\varphi}}$$

- Numerator is proportional to  $\text{Re } M^-$ 
  - $\cos \varphi$  part of interference term
- R can be compared directly with GPD models even in experiments with limited statistics
- Sensitive to Polyakov-Weiss D-term?



Comparison of results from e1-6/e1f with LO calculations by V. Guzey.

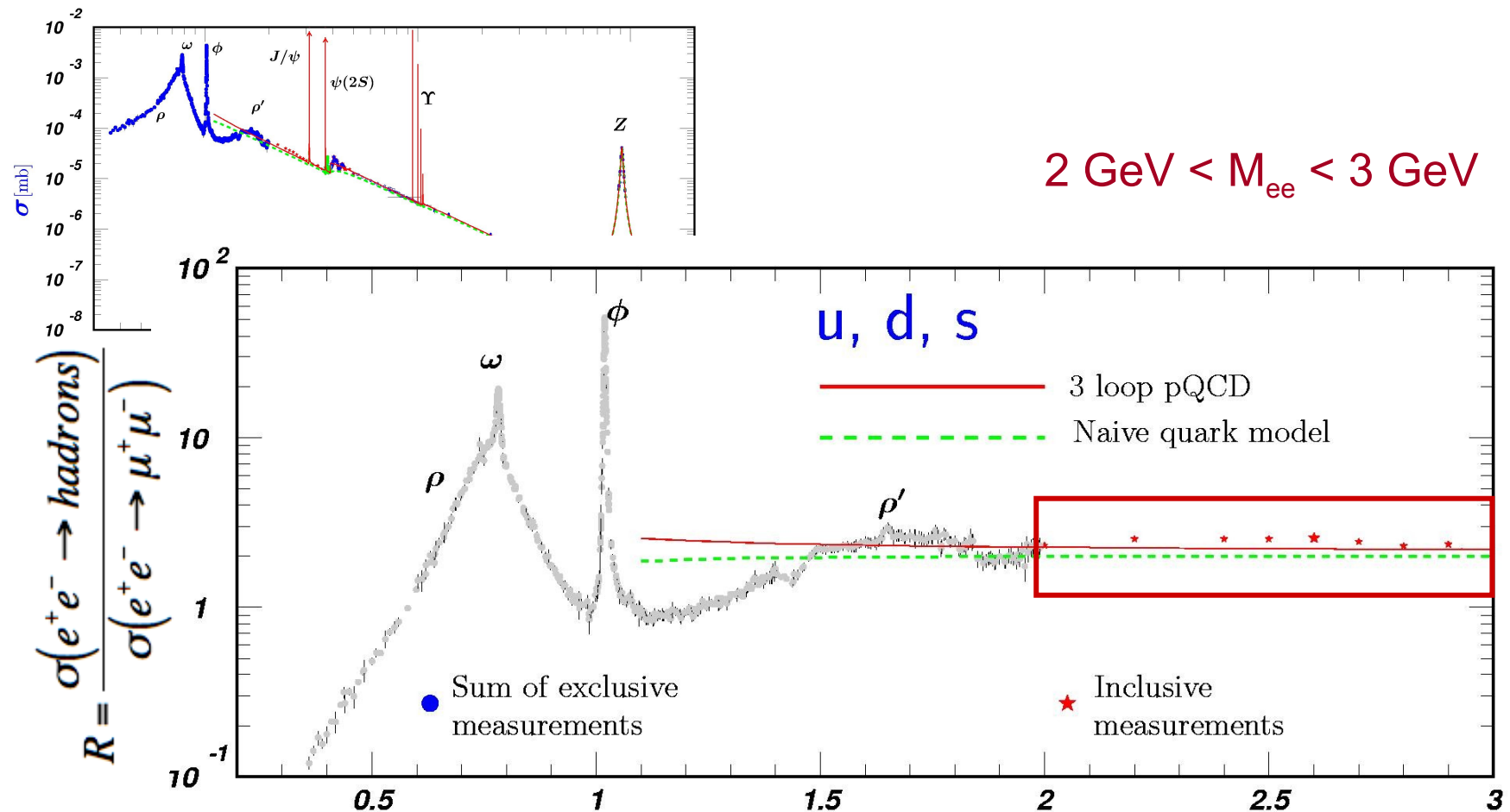
# Limited coverage and statistics at 6 GeV



Data from e1-6 (left), e1f (right), and g12 have limited coverage in  $s$  and  $M_{ee} = Q'$ .

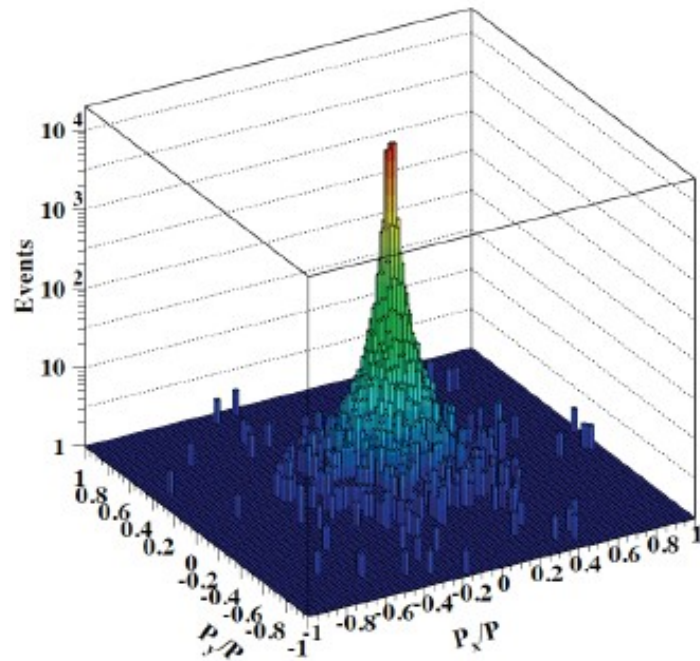
- Pilot experiments at 6 GeV are important for developing methods.
- 12 GeV will provide
  - A much larger reach in  $Q^2$  (factorization, x range)
  - More statistics for multi-dimensional binning
  - A possibility to avoid resonances

# Resonance-free region at 12 GeV

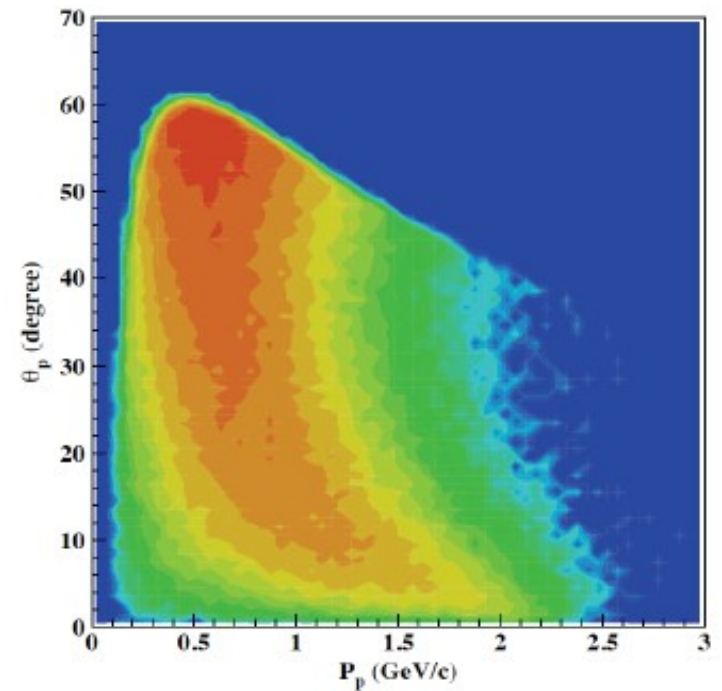


- JLab 12 GeV kinematics are ideally suited for TCS
- Data can be taken in the resonance-free region between  $\rho'$  and  $J/\psi$

# Beam electron and recoil proton kinematics

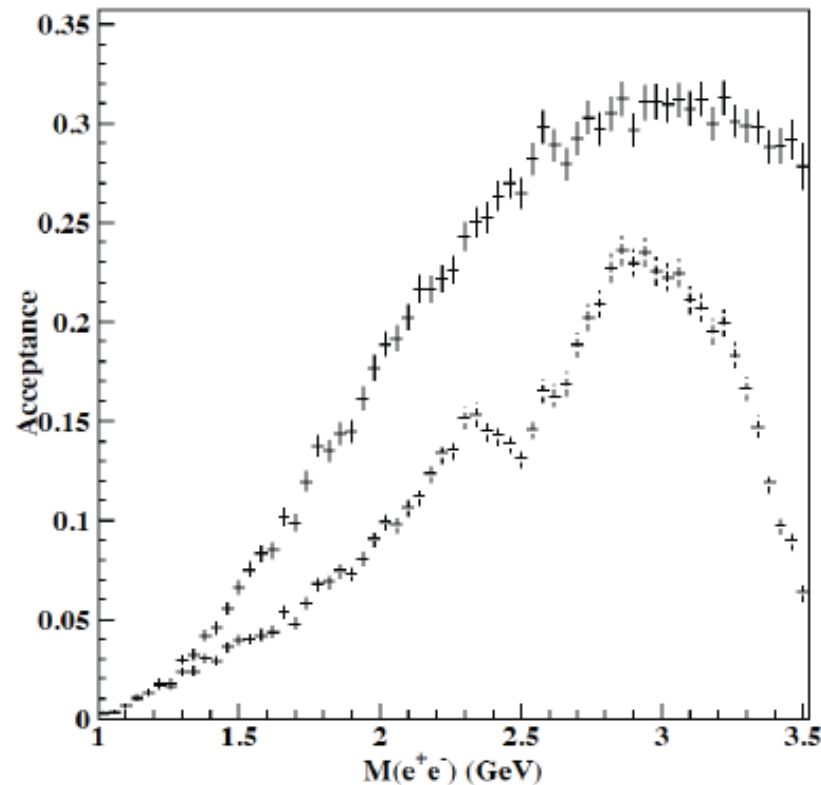


Low- $Q^2$  events are reconstructed by cuts on the transverse momentum of the missing beam electron



Protons from 9.5-10.5 GeV photons, for  $M_{ee} > 1.5$  GeV

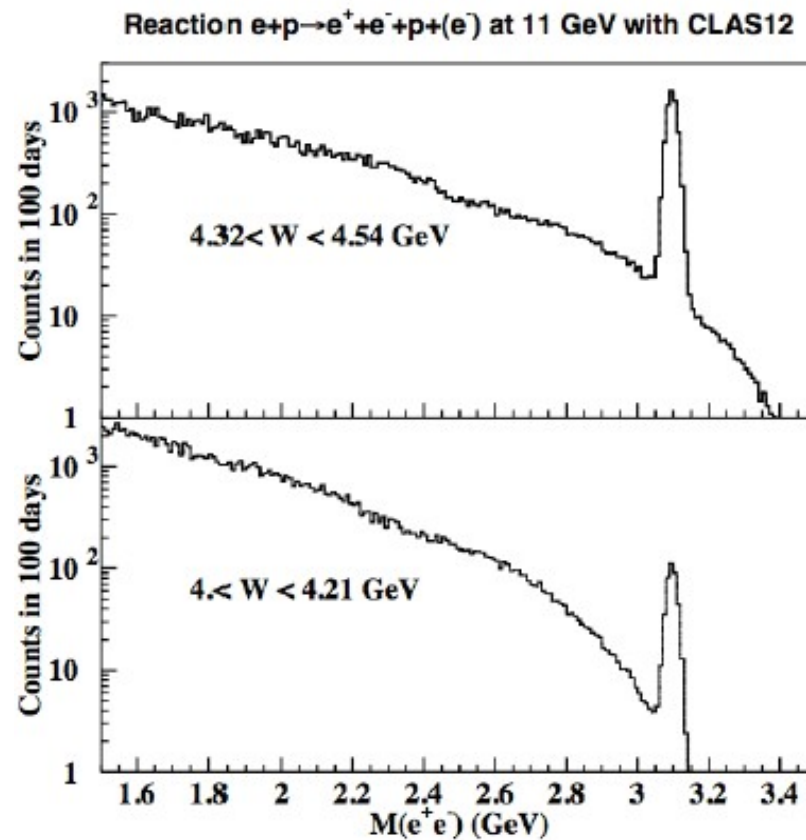
# Acceptance (from fast Monte Carlo)



- (Top): only  $e^+e^-$  detected in CLAS
- (Bottom):  $e^+e^-p$  detected in CLAS
- With an untagged beam of quasi-free photons we need all three for complete event kinematics

- Acceptance is good for the most interesting events at high  $M_{ee}$  ( $Q'$ ) due to the large lepton opening angle.

# Rate estimate



- Counts in 100 days as function of  $M_{ee} (Q')$
- For  $2 < Q' < 3 \text{ GeV}$ , the upper  $W = \sqrt{s}$  bin will have 20k events, and the lower 25k events.
- For analysis, the binning in  $W$  may be coarser, but the data will also be binned in  $t$  and  $\phi$ , and a cut will be applied in  $\theta^*$ .

# Beam time already approved for CLAS12

Proposal	Physics	Contact	Rating	Days	Group	needed equipment	Energy	Group	Target
E12-07-104	Neutron magnetic form factor	G. Gilfoyle	A-	30	90	Neutron detector RICH IC	11	A	liquid D <sub>2</sub> target
PR12-11-109 (a)	Dihadron DIS production	Avakian		D					
E12-09-007a	Study of partonic distributions in SIDIS kaon production	K. Hafidi	A-	56					
E12-09-008	Boer-Mulders asymmetry in K SIDIS w/ H and D targets	M. Contalbrigo	A-	TBA					
11-003	DVCS on neutron target	S. Niccolai	A	90					
E12-06-108	Hard exclusive electro-production of $\pi^0, \eta$	P. Stoler	B	80	119	RICH IC Forward tagger	11	B	liquid H <sub>2</sub>
E12-06-112	Probing the proton's quark dynamics in Semi-Inclusive pion production	H. Avakian	A	60					
E12-06-119	Deeply Virtual Compton Scattering	F. Salatie	A	80					
E12-09-103	Excitation of nucleon resonances at high Q <sup>2</sup>	R. Golhe	B+	40					
E11-005	Hadron spectroscopy with forward tagger	M. Battaglieri	A-	119					
PR12-11-103	DVMP of $\rho, \omega, \phi$	M. Guidal		D					
E12-06-106	Color transparency in exclusive vector meson electroproduction off nuclei	K. Hafidi	B+	60	60		11	C	Nuclear targets
E12-06-117	Quark propagation and hadron formation	W. Brooks	A-	60	60		11	D	Nuclear
E12-10-102	Free Neutron structure at large x	S. Bueltman	A	40	40	Radial TPC	11	E	Gas D <sub>2</sub>
E12-06-109	Longitudinal Spin Structure of the Nucleon	S. Kuhn	A	80	170	Polarized target RICH IC	11	F	NH <sub>3</sub> ND <sub>3</sub>
E12-06-119(b)	DVCS on longitudinally polarized proton target	F. Salatie	A	120					
E12-07-107	Spin-Orbit Correl. with Longitudinally polarized target	H. Avakian	A-	103					
PR12-11-109 (b)	Dihadron studies on long. polarized target	H. Avakian		D					
E12-09-007(b)	Study of partonic distributions using SIDIS K production	K. Hafidi	A-	110					
E12-09-009	Spin-Orbit correlations in K production w/ pol. targets	H. Avakian	B+	103					
PR12-11-109	SIDIS on transverse polarized target	M. Contalbrigo		C2		Transverse target	11	G	HD
TOTAL run time				1231	539				

# Running conditions and beamtime request

## Running conditions

- The TCS proposal requires 11 GeV beam and a LH2 target.
  - There are no other restrictions so beamtime can easily be shared.

## Beam time request

- The TCS proposal will request about 100 days at full luminosity
  - Simulations are still ongoing

## Notes

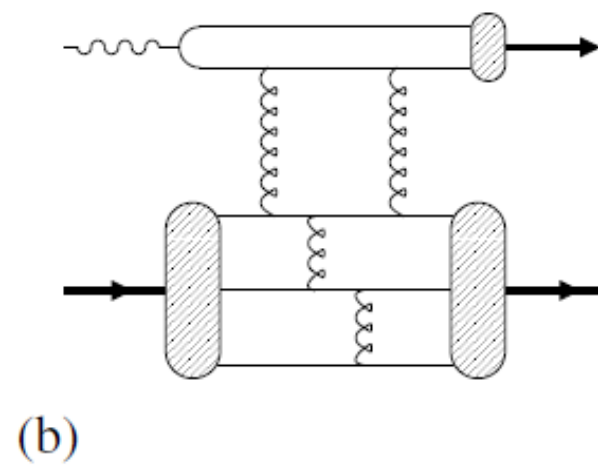
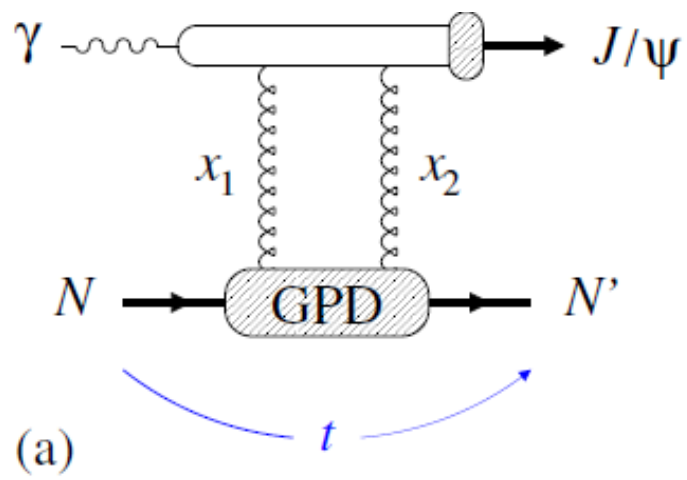
- It would be possible to study TCS on other targets, but there is no obvious motivation for doing so at this point
- The upcoming  $J/\Psi$  proposal for the nucleon will request additional time with a reversed field for systematics checks.

A photograph of a child sitting on a small, colorful polka-dot stool and playing a white piano. The piano and stool are also covered in the same colorful polka-dot pattern. The walls and floor of the room are entirely covered in a dense, multi-colored polka-dot pattern, creating a vibrant and playful environment. The text "Thank you!" is overlaid on the right side of the image in a stylized, yellow, italicized font with a black outline.

*Thank you!*

# Backup

$J/\psi$

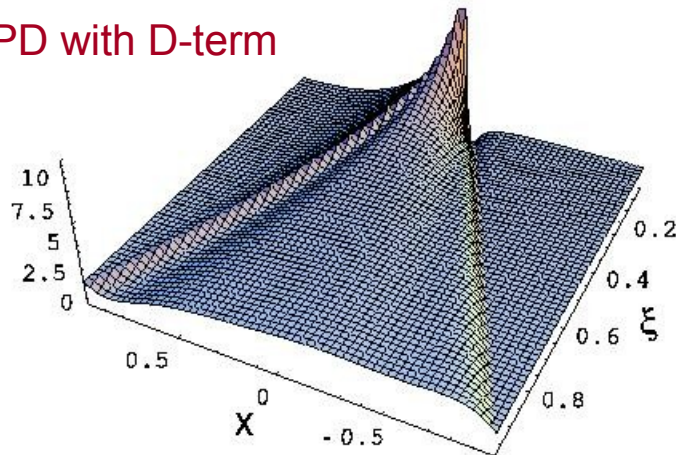


## D-term in DD-parametrization of GPDs

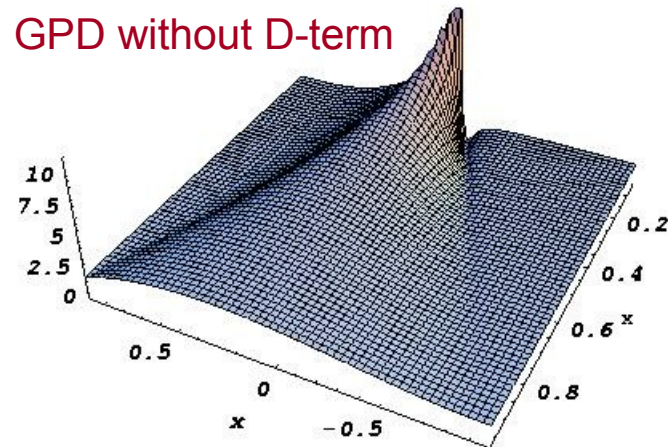
$$H^q(x, \xi) = H_{DD}^q(x, \xi) + \theta(\xi - |x|) \frac{1}{N_f} \mathbf{D}\left(\frac{x}{\xi}\right)$$

D-term – allows to satisfy polynomiality of Mellin moments of GPDs

GPD with D-term



GPD without D-term



**Real part of the Compton amplitude is sensitive to the D-term**