

Exclusive π^0 studies for calorimeter resolutions

Maxime Defurne - CEA Saclay

April 11, 2020

Contents

1	Introduction	1
2	Electron kinematics	1
3	π^0 kinematics	2
4	Recoil proton kinematics	3
5	Discussion	3

1 Introduction

Considering the very small mass of neutral pion with respect to the energies at EIC, these kinematics studies are the same for DVCS. In the following studies, I have applied the following cuts for the electron plots :

- $Q^2 < 100 \text{ GeV}^2$.
- $\theta_e < 140^\circ$.
- $W^2 > 4 \text{ GeV}^2$.
- $y > 0.01$ and $y < 0.95$.

with in addition the following cut on the π^0 -momentum for the π^0 and the recoil proton plots :

- $2 \text{ GeV} < p_{\pi^0} < 80 \text{ GeV}$.

Since one of the main observable is the t -dependence of the cross sections, I have looked at the π^0 -momentum for $t=t_{min}$, -0.5 GeV^2 and -1 GeV^2 . In the case where t_{min} is lower than the requested t -value, then t_{min} is used y default. If t_{max} is greater than the requested t -value, then t_{max} is used y default.

Following the recommendations of the accelerator group, I have studied the configurations :

- Electron beam at 5 GeV and Proton beam at 41 GeV.
- Electron beam at 5 GeV and Proton beam at 100 GeV.
- Electron beam at 10 GeV and Proton beam at 100 GeV.
- Electron beam at 18 GeV and Proton beam at 275 GeV.

2 Electron kinematics

Here are the electron momentum/rapidity distributions as a function of Q^2 and x_B . I have required to have $y > 0.01$ and $y < 0.95$. I have assumed that the maximal scattering angle we can detect for the electron is 140 degrees. These plots have been produced without applying any cut on the π^0

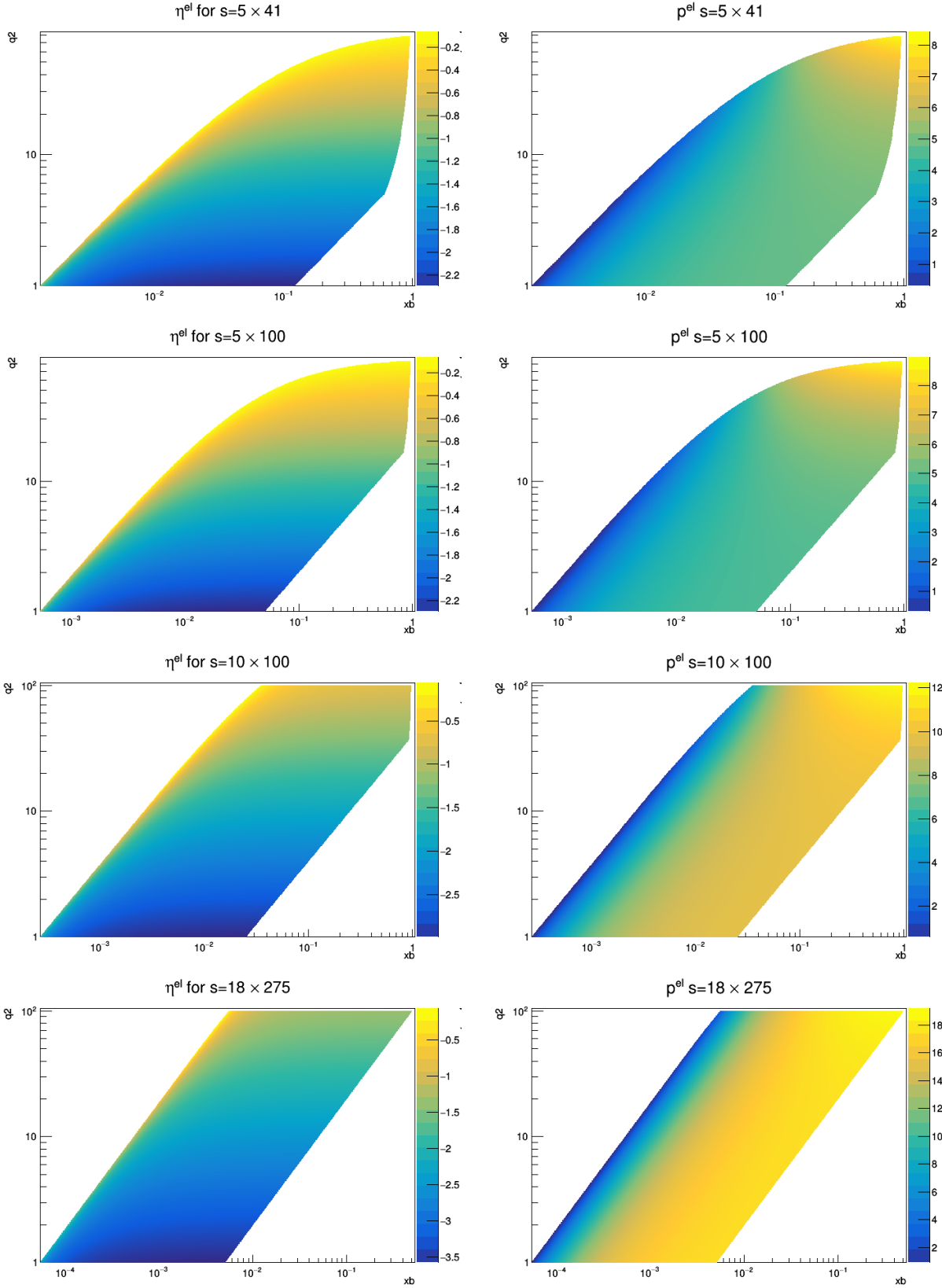


FIGURE 1 – Rapidity (Left) and momentum (right) of scattered electron as function of Q^2 and x_B .

3 π^0 kinematics

The $\text{tmin-}\pi^0$ is generated by having it collinear to the virtual photon exchanged between the proton and the electron. As Q^2 increases, the π^0 is more backward and its momentum slightly increases. The π^0 -momentum has a minimum as function of x_B for Q^2 less than 600 GeV^2 for the highest $\sqrt{s}$, and then becomes a monotonic increasing function of x_B for higher Q^2 values.

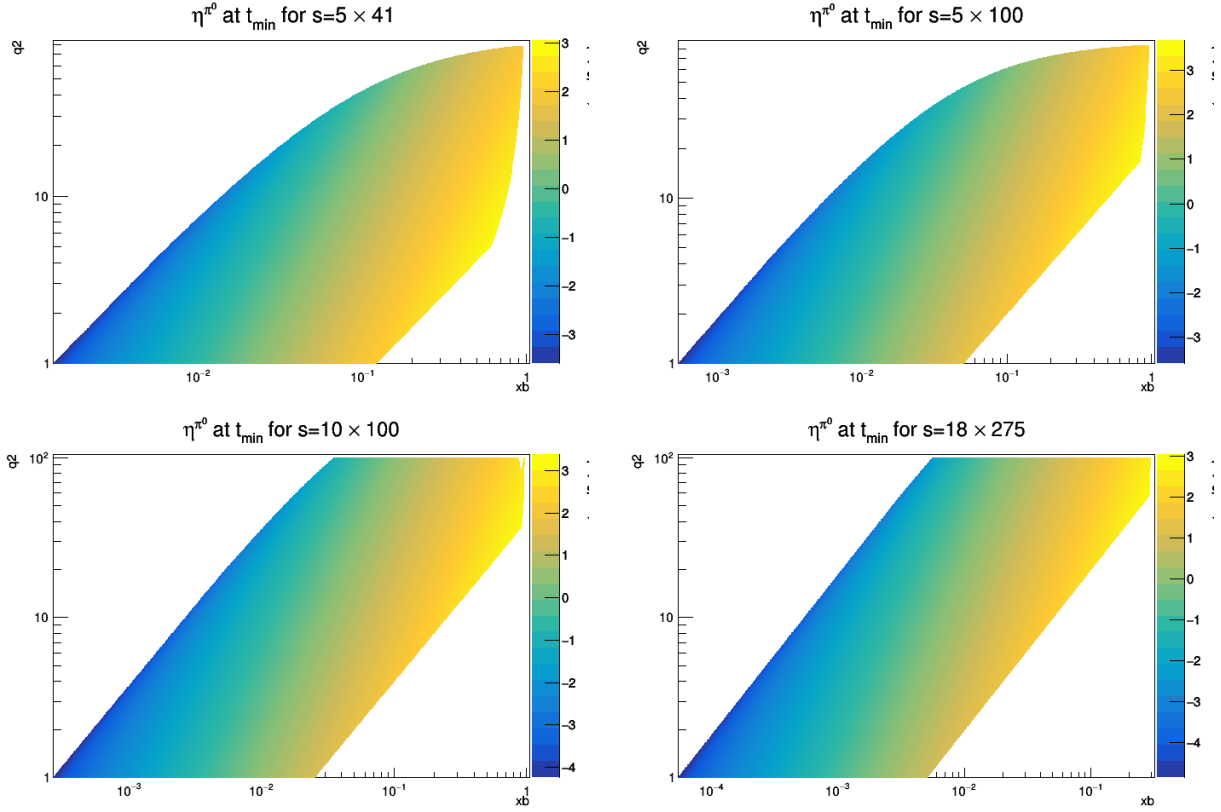


FIGURE 2 – π^0 -rapidity at $t=t_{min}$ as function of Q^2 and x_B for the four beam configurations.

Taking into account the possible limitation of the calorimeter on the lowest photon energy at 1 GeV, I have applied a cut at 2 GeV on the π^0 momentum for three t -values which are t_{min} , -0.5 and -1 GeV².

For all $\sqrt{s}$ -configuration, a significant phase space is lost at lower Q^2 and high x_B for t ranging from t_{min} to -1 GeV².

4 Recoil proton kinematics

For the moment, no constrain are applied on the proton side. Here are the proton kinematics in the same configurations as it was presented for the π^0 in the previous section.

5 Discussion

In order to determine the most suitable configurations for a GPD extraction with DVCS data, we must go further with pseudo-phenomenological analysis. To cover $x_B > 0.1$, it seems that the configuration 5×100 is more suitable than 5×41 because of 2 GeV cut on the π^0 -momentum. Although, because of the cut on $y > 0.1$, we are losing low Q^2 phase space for x_B above 0.2. Would we get enough statistics at 5×100 compared to 5×41 for x_B above 0.2 despite the higher Q^2 -values? Or is there any technical reasons behind the cut $y > 0.01$? Otherwise we could even recover these low Q^2 -points.

Finally the region x_B from 10^{-3} to 10^{-2} for Q^2 smaller than 3 GeV² will not be accessible for DVCS studies because the π^0 cannot be studied. I remind that the π^0 cross sections increases with x_B and decreases with Q^2 .

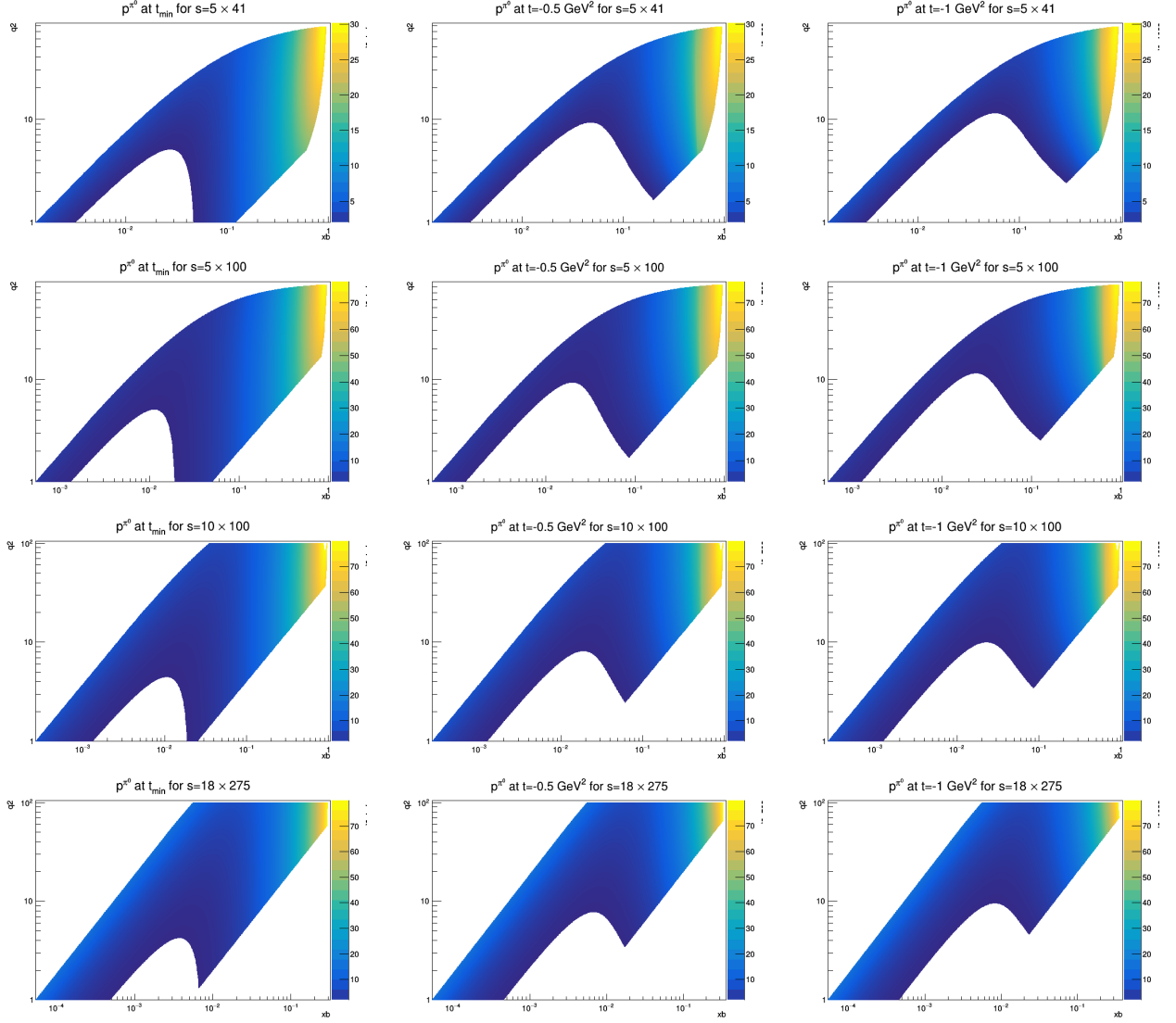


FIGURE 3 – π^0 -momentum as function of Q^2 and x_B for $t=t_{\min}$, -0.5 and -1 GeV^2 .

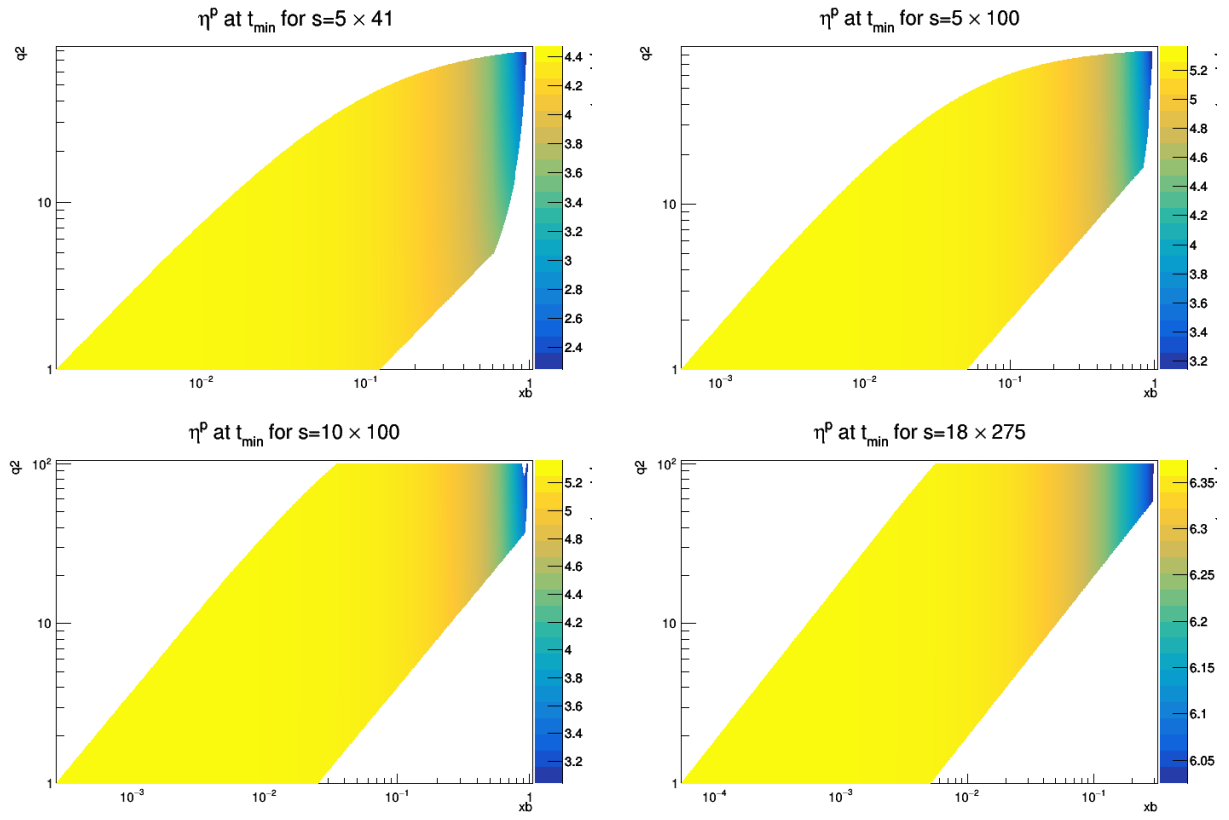


FIGURE 4 – Proton rapidity at $t=t_{\min}$ as function of Q^2 and x_B for the four beam configurations.

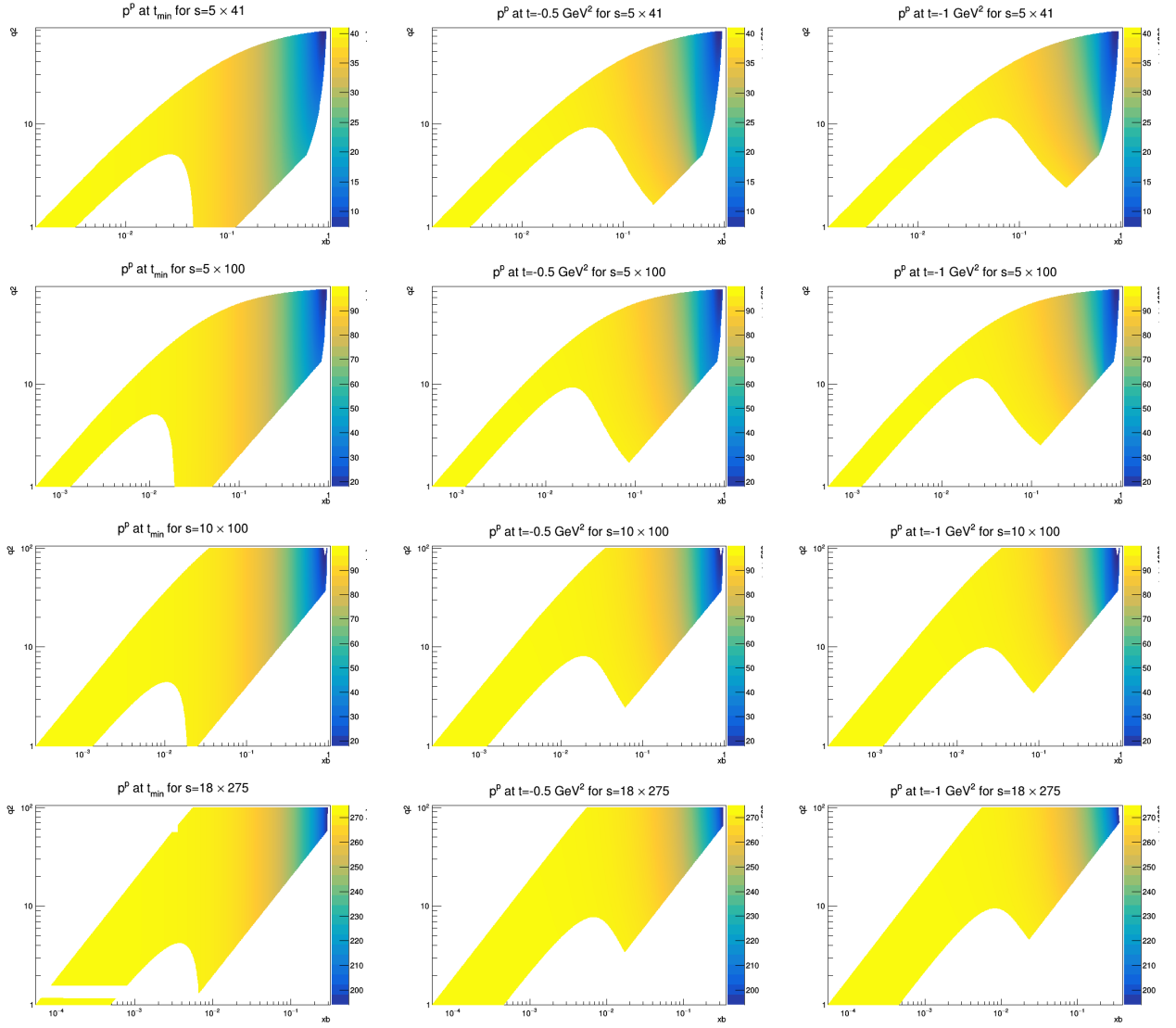


FIGURE 5 – Proton momentum as function of Q^2 and x_B for $t=t_{\min}$, -0.5 and -1 GeV^2 .