

Determination of the Hyperon Induced Polarization and Polarization–Transfer Coefficients for Quasi–Free Hyperon Photoproduction off the Bound Neutron

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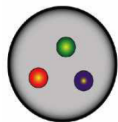
Overview

- Baryon Spectroscopy and Jefferson Lab
- Strangeness ($K\Lambda$) photoproduction
- Preliminary results for $\vec{\gamma}d \rightarrow K^0\vec{\Lambda}(p)$
 - 1 Comparison to theoretical predictions
 - 2 Comparison to free proton data
 - 3 Dependence on the neutron momentum

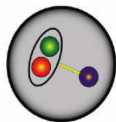
Baryon Spectroscopy

Provides a way to measure the excited nucleon (N^*) spectrum

- Excited nucleon states \rightarrow understanding of nucleon
- Map N^* spectrum to learn about the internal structure of nucleons
- Goal is to provide information about the underlying degrees of freedom in the non-perturbative regime



CQM



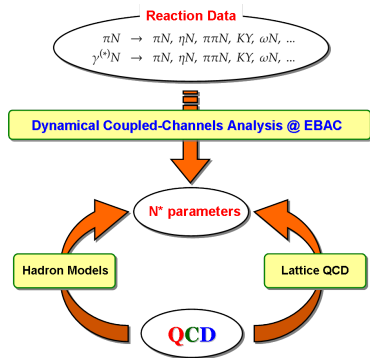
Quark-diquark clustering

Constituent quark models three valence quarks

Di-quark models bound quark pair \rightarrow less degrees of freedom

Lattice QCD numerical solution to QCD

Many other approaches...



<http://ebac-theory.jlab.org/>

Jefferson Lab, Hall B, and CLAS



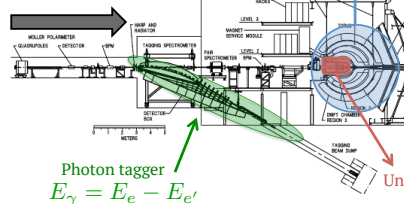
- CEBAF accelerated e up to 6 GeV

Jefferson Lab, Hall B, and CLAS

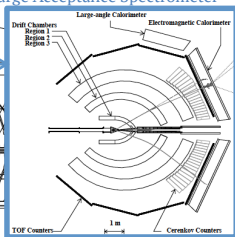


B.A. Mecking et al., Nucl. Instr. and Meth. A 503, 513 (2003)

Polarized electron beam- g13a
energies up to $E_e = 2.6$ GeV



CEBAF Large Acceptance Spectrometer
(CLAS)



- $\approx 80\%$ e polarization
- $P_\gamma \approx 20 - 95\% P_e$
- $0.9 < E_\gamma < 2.55$ GeV

Motivation for $\vec{\gamma}d \rightarrow K^0 \vec{\Lambda}(p)$

- Need more data
 - Majority of data is πN scattering or final state
 - Some resonances couple weakly to these channels while having significant branching ratios to $K\Lambda$
- Most strangeness data from free proton
 - $\gamma p \rightarrow K^+ \Lambda$ moving $N(1900)\frac{3}{2}^+$ from $\star\star$ to $\star\star\star$
 - $\gamma n \rightarrow K^0 \Lambda$ sensitive to $\star\star N(2150)\frac{3}{2}^-$ and $\star\star\star N(1875)\frac{3}{2}^-$
- How do data from the proton and bound neutron compare to each other?

Particle	J^P	Status		Status as seen in —							
		overall	πN	γN	$N\eta$	$N\sigma$	$N\omega$	ΛK	ΣK	$N\rho$	$\Delta\pi$
N	$1/2^+$	****									
$N(1440)$	$1/2^+$	****	****	****		***				*	***
$N(1520)$	$3/2^-$	****	****	****	***					***	***
$N(1535)$	$1/2^-$	****	****	****	****					**	*
$N(1650)$	$1/2^-$	****	****	****	***			***	**	**	***
$N(1675)$	$5/2^-$	****	****	***	*			*		*	***
$N(1680)$	$5/2^+$	****	****	****	*	**				***	***
$N(1685)$	$?^?$	*									
$N(1700)$	$3/2^-$	***	***	**	*			*	*	*	***
$N(1710)$	$1/2^+$	***	***	***	***		**	***	***	*	**
$N(1720)$	$3/2^+$	****	****	***	***			**	**	**	*
$N(1860)$	$5/2^+$	**	**							*	*
★ $N(1875)$	$3/2^-$	***	*	***			**	***	**		***
$N(1880)$	$1/2^+$	**	*	*		**		*			
$N(1895)$	$1/2^-$	**	*	**	**			**	*		
★ $N(1900)$	$3/2^+$	***	**	***	**		**	***	***	*	**
$N(1990)$	$7/2^+$	**	**	**					*		
$N(2000)$	$5/2^+$	**	*	**	**			**	*	**	
$N(2040)$	$3/2^+$	*									
$N(2060)$	$5/2^-$	**	**	**	*				**		
$N(2100)$	$1/2^+$	*									
★ $N(2150)$	$3/2^-$	**	**	**				**			**
$N(2190)$	$7/2^-$	****	****	***		*	**			*	
$N(2220)$	$9/2^+$	****	****								
$N(2250)$	$9/2^-$	****	****								
$N(2600)$	$11/2^-$	***	***								
$N(2700)$	$13/2^+$	**	**								

Polarization Observables in $K\Lambda$ Photoproduction

- Can't “bump hunt” on the energy evolution of the cross-section
- 16 Polarization observables are derived from the matrix elements of the scattering matrix

Polarization Observables in $K\Lambda$ Photoproduction

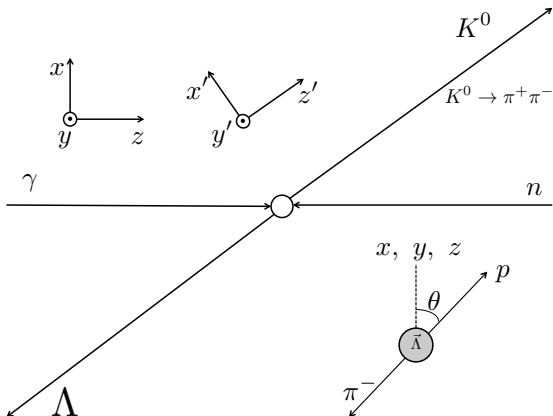
- Can't “bump hunt” on the energy evolution of the cross-section
- 16 Polarization observables are derived from the matrix elements of the scattering matrix

Unpolarized Cross Section	σ_0			
Single		P	Σ	T
Beam-Recoil	C_x	C_z	O_x	O_z
Target-Recoil	T_x	T_z	L_x	L_z
Beam-Target	E	F	G	H

- 8 carefully chosen observables are needed to determine the full scattering amplitude

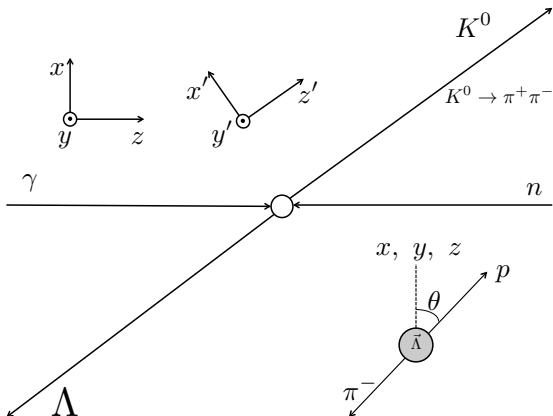
$$\frac{d\sigma}{d\Omega} = \sigma_0 [1 - P_{lin} \Sigma \cos 2\phi - \alpha \cos \theta_x (P_{lin} O_x \sin 2\phi + P_{circ} C_x) - \alpha \cos \theta_y (-P + P_{lin} T \cos 2\phi) - \alpha \cos \theta_z (P_{lin} O_z \sin 2\phi + P_{circ} C_z)]$$

Axis Conventions for $\gamma n \rightarrow K^0 \Lambda$



① $\vec{\gamma} n \rightarrow N^*$

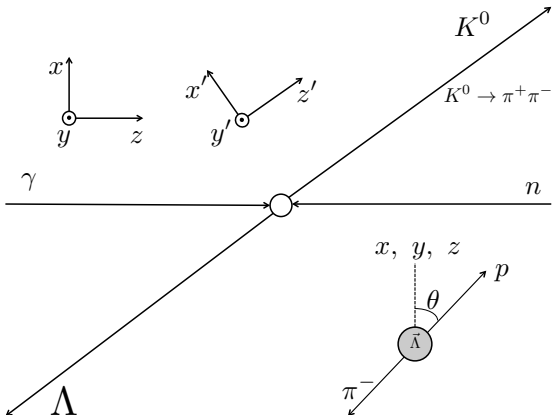
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Axis Conventions for $\gamma n \rightarrow K^0 \Lambda$

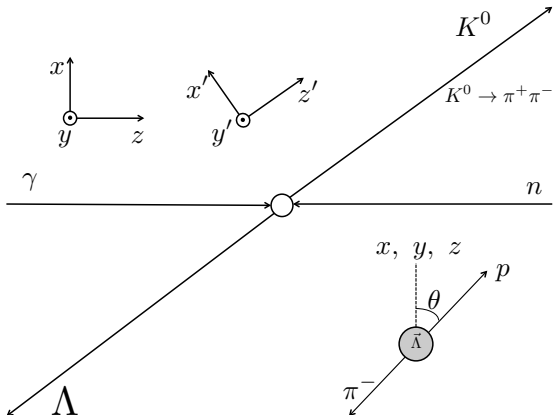


① $\vec{\gamma} n \rightarrow N^*$

② $N^* \rightarrow K^0 \Lambda$

③ $K^0 \rightarrow \pi^+ \pi^-$ and $\Lambda \rightarrow p \pi^-$

Axis Conventions for $\gamma n \rightarrow K^0 \Lambda$



- 1 $\vec{\gamma} n \rightarrow N^*$
- 2 $N^* \rightarrow K^0 \Lambda$
- 3 $K^0 \rightarrow \pi^+ \pi^-$ and $\Lambda \rightarrow p \pi^-$
- 4 Using information about θ , we can calculate the polarization transfer from $\vec{\gamma}$ to Λ

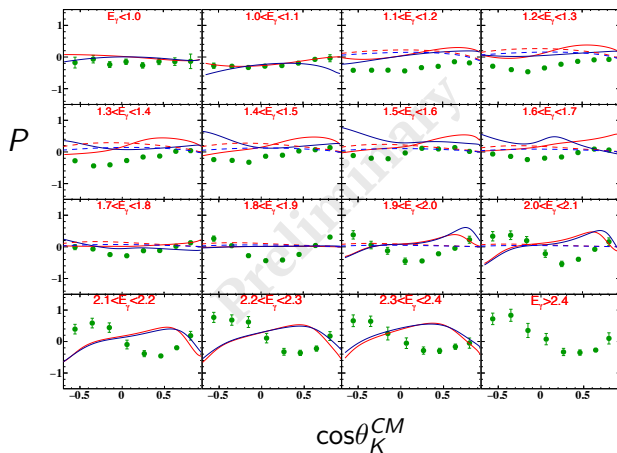
$$\frac{d\sigma}{d\Omega} = \sigma_0 [1 - \alpha \cos \theta_{x'} P_{\text{circ}} C_{x'} + \alpha \cos \theta_{y'} P - \alpha \cos \theta_{z'} P_{\text{circ}} C_{z'}]$$

Preliminary Results

Will present and discuss the first (preliminary) estimates for C_x and C_z for $\vec{\gamma}d \rightarrow K^0 \vec{\Lambda}(p)$

- 1 Compare to models using the primed axis convention (z along K^0 momentum).
- 2 Compare to data off the free proton in the unprimed axis convention (z along γ momentum).
- 3 Dependence on the momentum of the bound neutron.

Comparison to Models- P



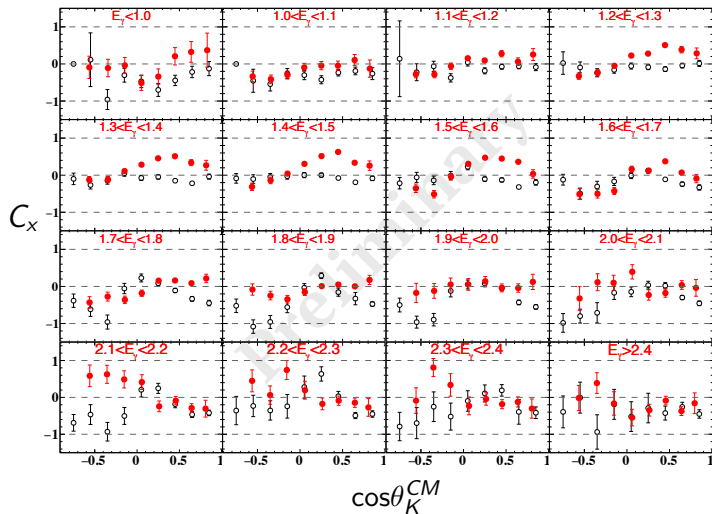
Two Models

- 1 Waluyo (solid), CCA
 - 2 Kaon-MAID (dashed), PWA
- No D13(1900)
● With D13(1900)

1) A. Waluyo, Ph.D. Thesis, The George Washington University, (2005).

2) F.X. Lee, T. Mart, C. Bennhold, H. Haberzettl, L.E. Wright, Nucl. Phys. A695 (2001) 237

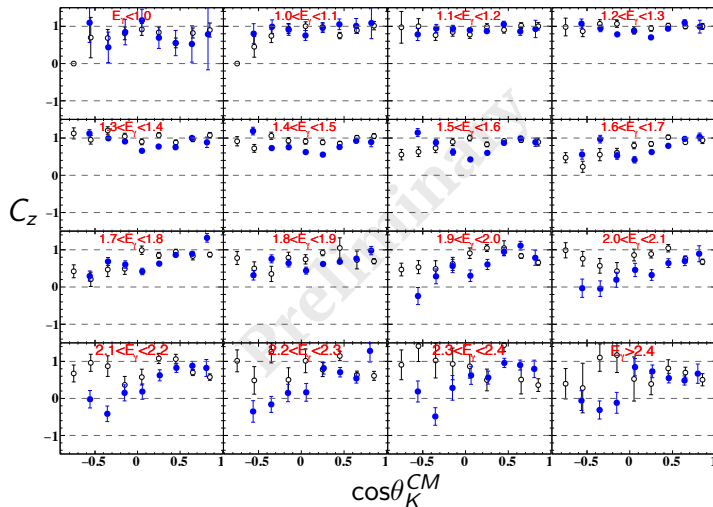
C_x : Comparison of $\gamma d \rightarrow K^0 \Lambda(p)$ to $\gamma p \rightarrow K^+ \Lambda$



$\gamma d \rightarrow K^0 \Lambda(p)$
 $\gamma p \rightarrow K^+ \Lambda$

R. K. Bradford et al. (CLAS Collaboration), Phys. Rev. C 75, 035205

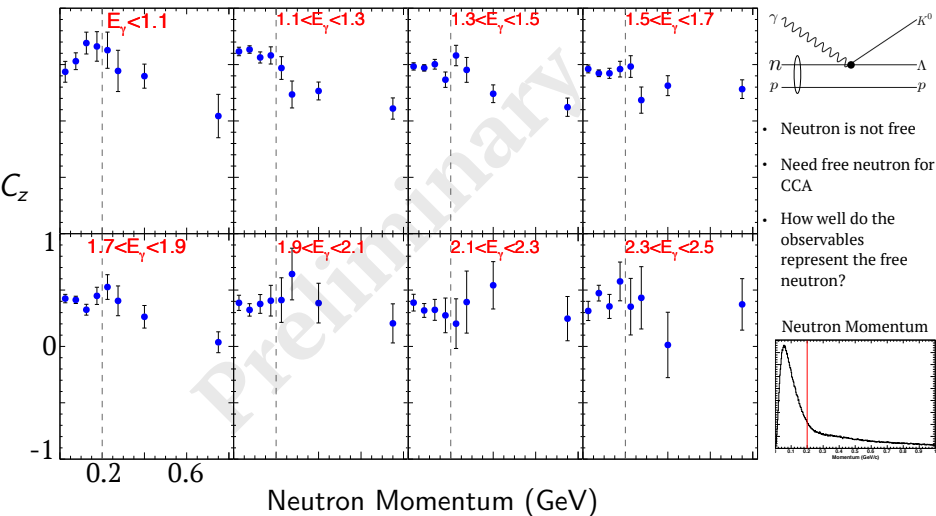
C_z : Comparison of $\gamma d \rightarrow K^0 \Lambda(p)$ to $\gamma p \rightarrow K^+ \Lambda$



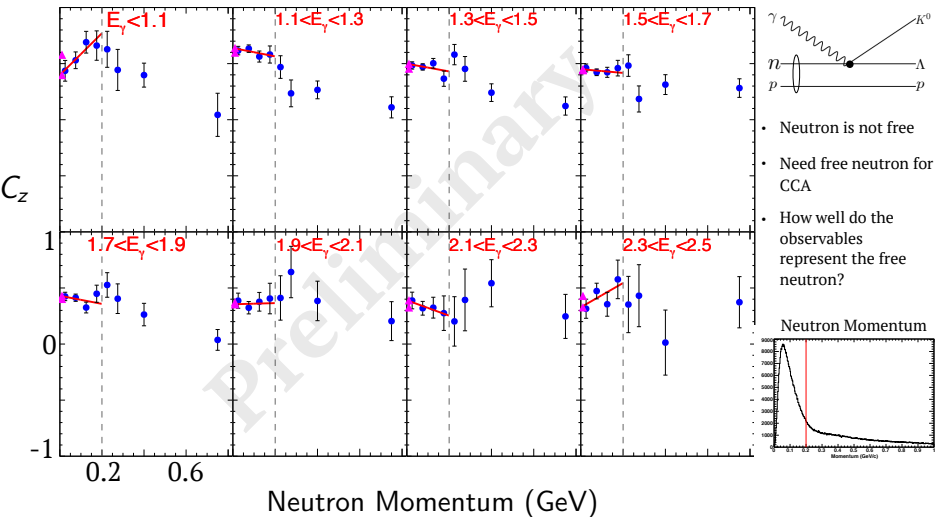
$\gamma d \rightarrow K^0 \Lambda(p)$
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R. K. Bradford et al. (CLAS Collaboration), Phys. Rev. C 75, 035205

Dependence on Neutron Momentum— C_z



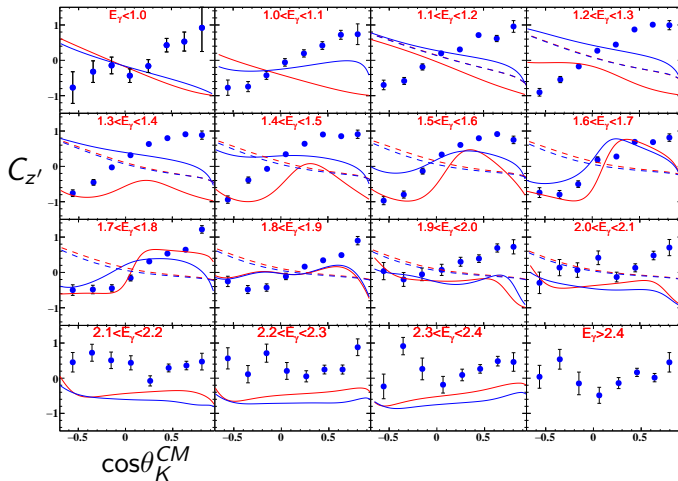
Dependence on Neutron Momentum– C_z



Conclusion

- Many resonant states predicted by constituent quark models have yet to be observed
- Hyperon channels have a strong coupling to some of these resonances
- Very first (preliminary) estimates of polarization observables for data off the bound neutron have been obtained.
- **Differences with model calculations from are observed. Data are expected to have an impact.**
- **Differences between data off free proton and quasi-free neutron are observed that could be due to reaction dynamics or other resonances.**
- **C_z is dependent on the neutron momentum, but below 0.2 GeV it is mostly constant.**

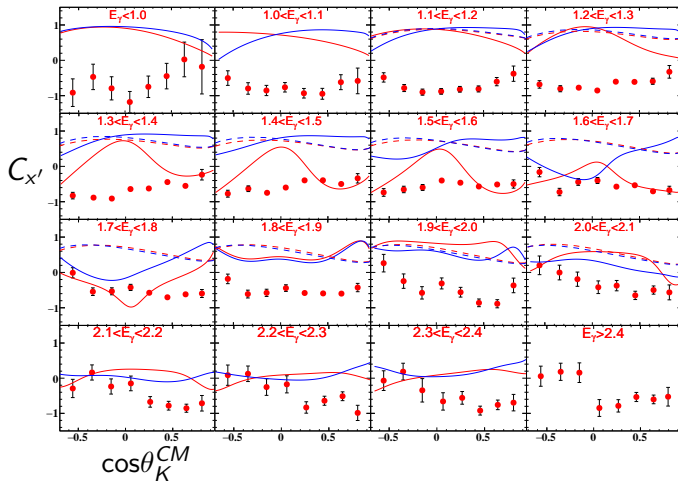
Comparison to Models- $C_{Z'}$



Two Models

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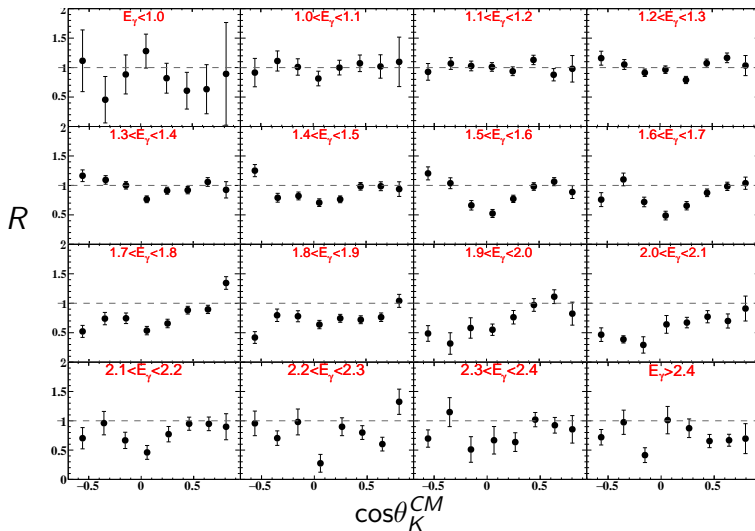
Comparison to Models- $C_{X'}$



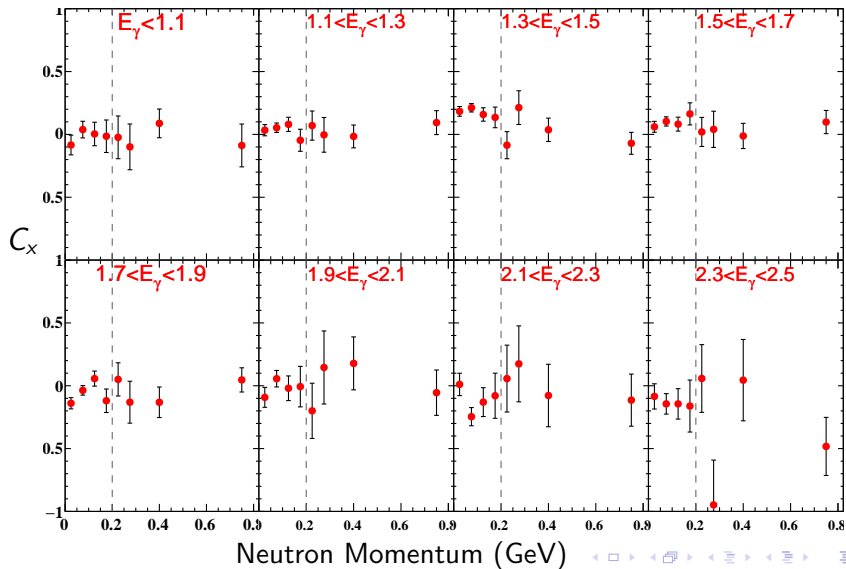
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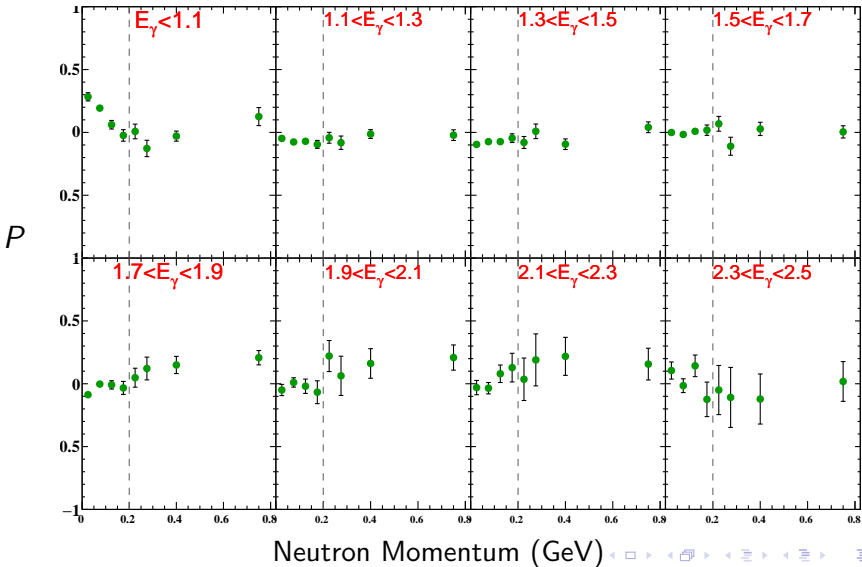
$$R = \sqrt{C_x^2 + C_z^2 + P^2} \leq 1 \text{--Total Polarization Transfer}$$



Dependence on Neutron Momentum– C_x

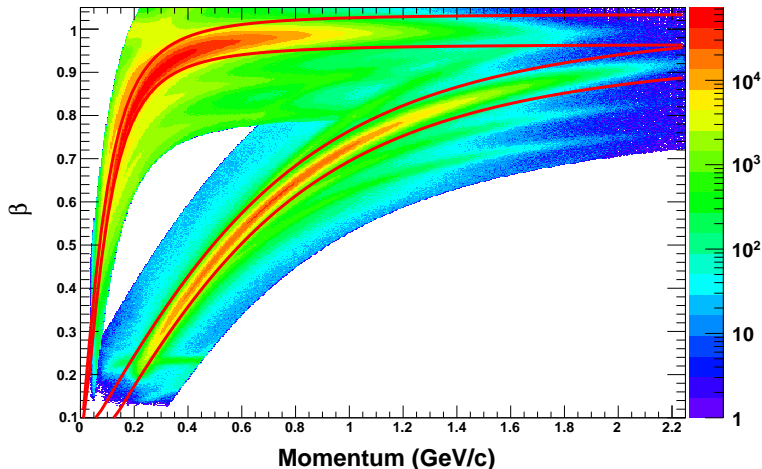


Dependence on Neutron Momentum- P



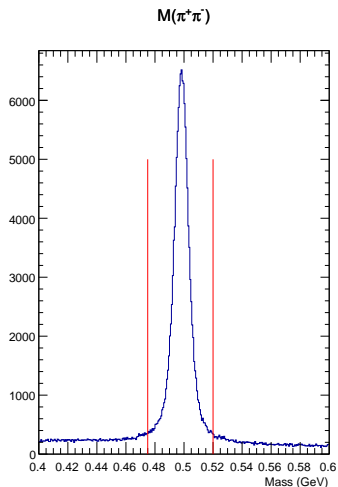
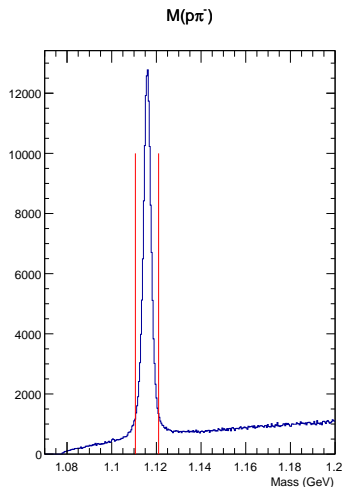
Particle Identification

Particles are identified based off their β and momentum in CLAS
 β vs. Momentum



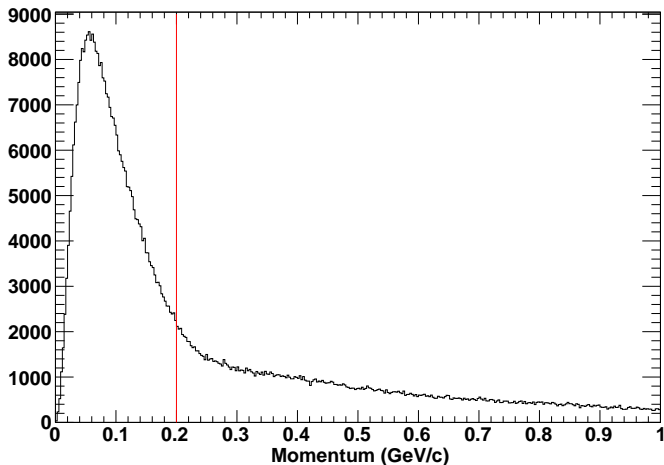
Reaction Identification

The Λ and K^0 need to be reconstructed since $\Lambda \rightarrow p\pi^-$ and $K^0 \rightarrow \pi^+\pi^-$



Quasi-free Selection

The neutron momentum can be reconstructed
Missing Momentum p_x



Final-state Identification

The missing proton is reconstructed

M_x

