

The g_2^p Experiment

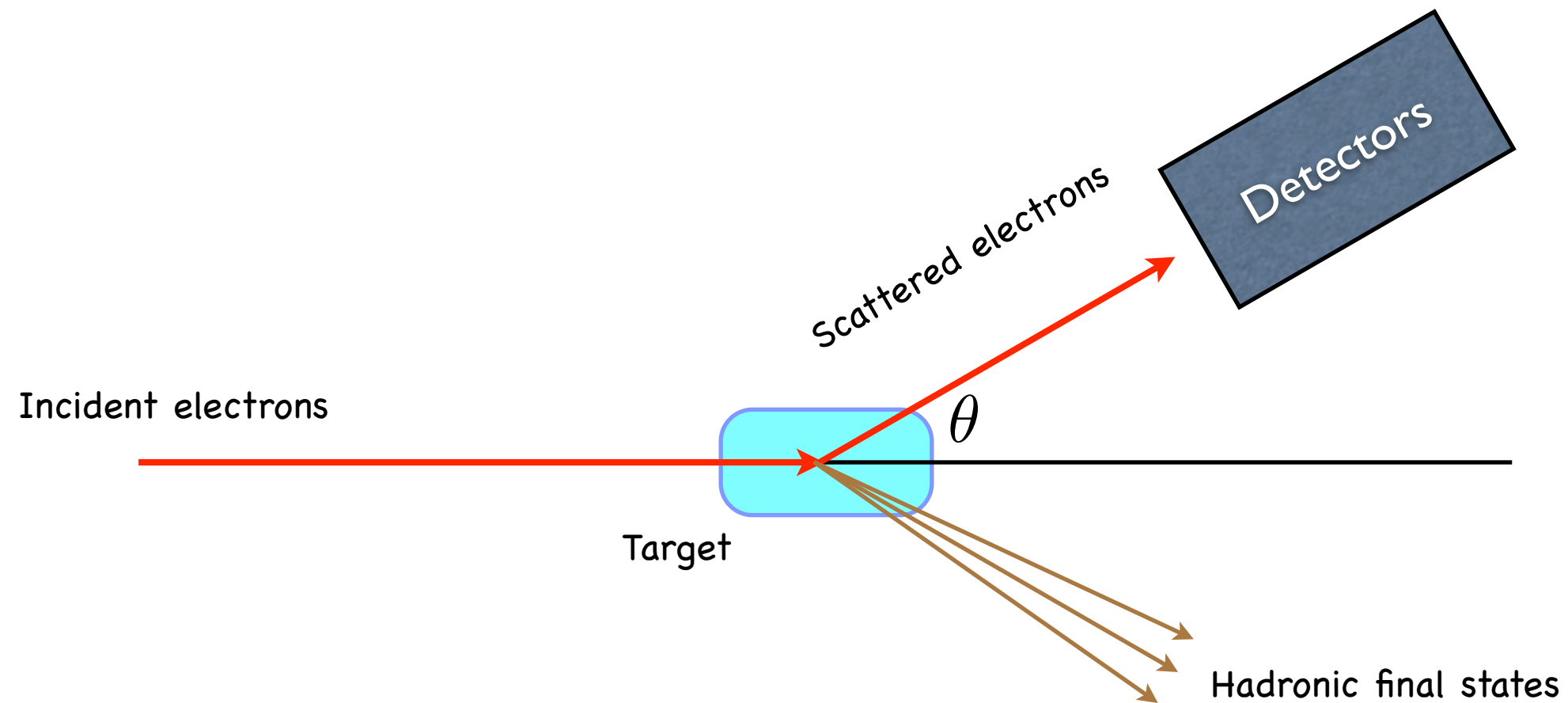
Chao Gu

Chiral Dynamics Workshop, Aug 2012

Outline

- Review of physics motivation
- Brief review of experiment setup
- Status of experiment run

Inelastic scattering

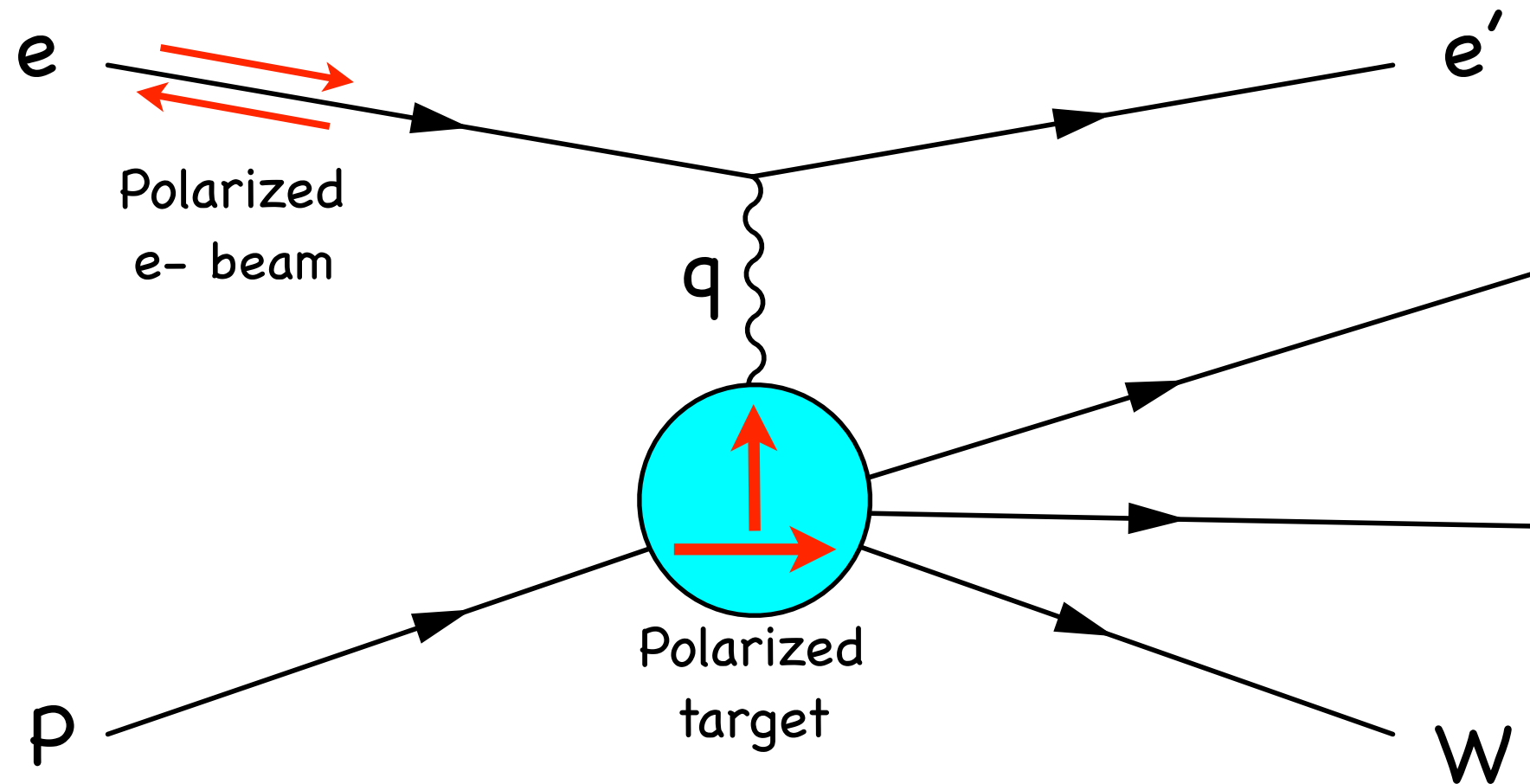


- Inclusive **unpolarized** cross section:

$$\frac{d^2\sigma}{d\Omega dE'} = \sigma_{\text{Mott}} \left[\frac{1}{\nu} F_2(x, Q^2) + \frac{2}{M} F_1(x, Q^2) \tan^2 \frac{\theta}{2} \right]$$

Structure Function
which indicates the
parton distribution

Inelastic scattering



- Inclusive **polarized** cross section:

$$\frac{d^2\sigma}{d\Omega dE'} = \sigma_{\text{Mott}} \left[\frac{1}{\nu} F_2(x, Q^2) + \frac{2}{M} F_1(x, Q^2) \tan^2 \frac{\theta}{2} + \gamma g_1(x, Q^2) + \delta g_2(x, Q^2) \right]$$

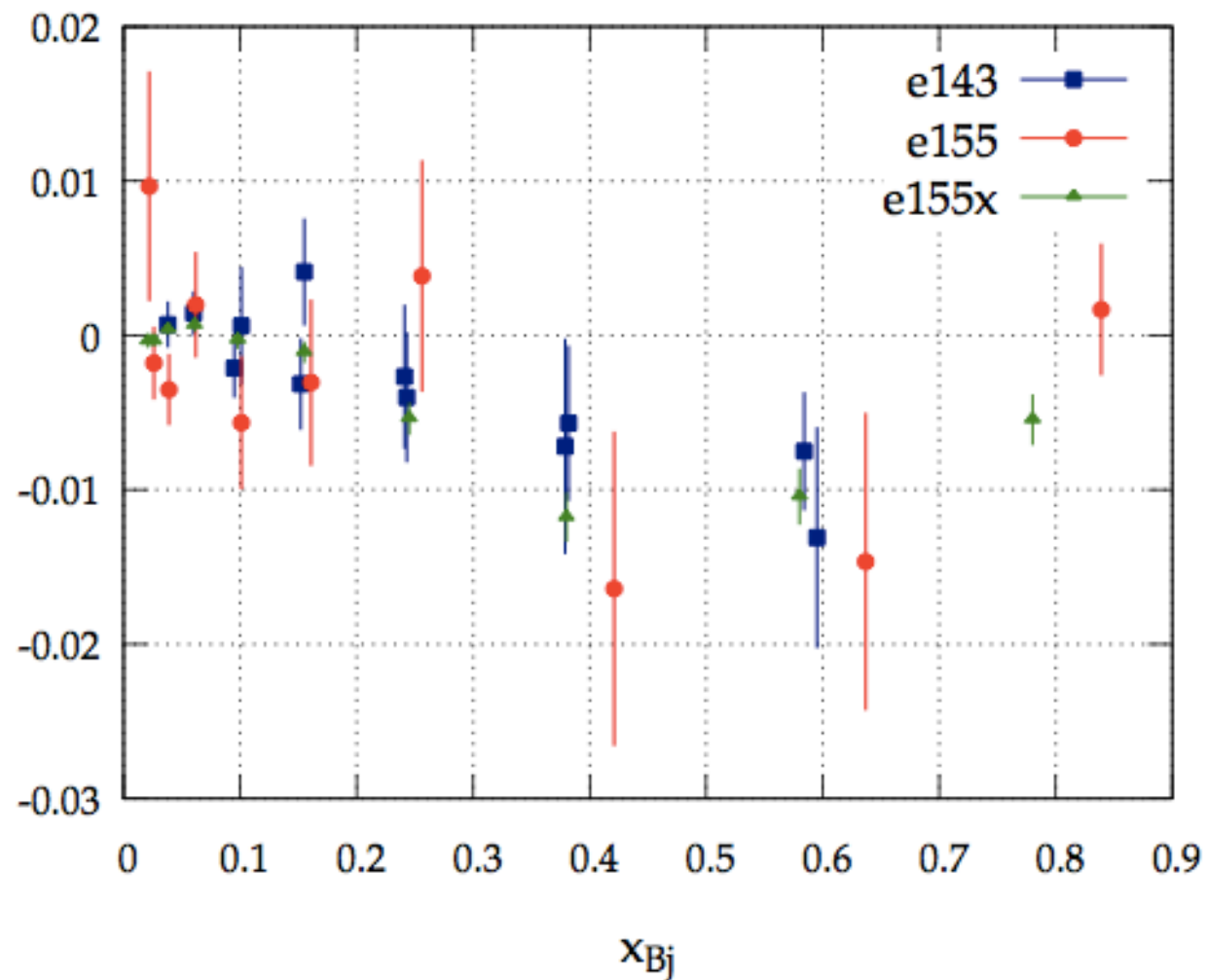
2 addition Structure Function which related to the spin distribution

Motivation

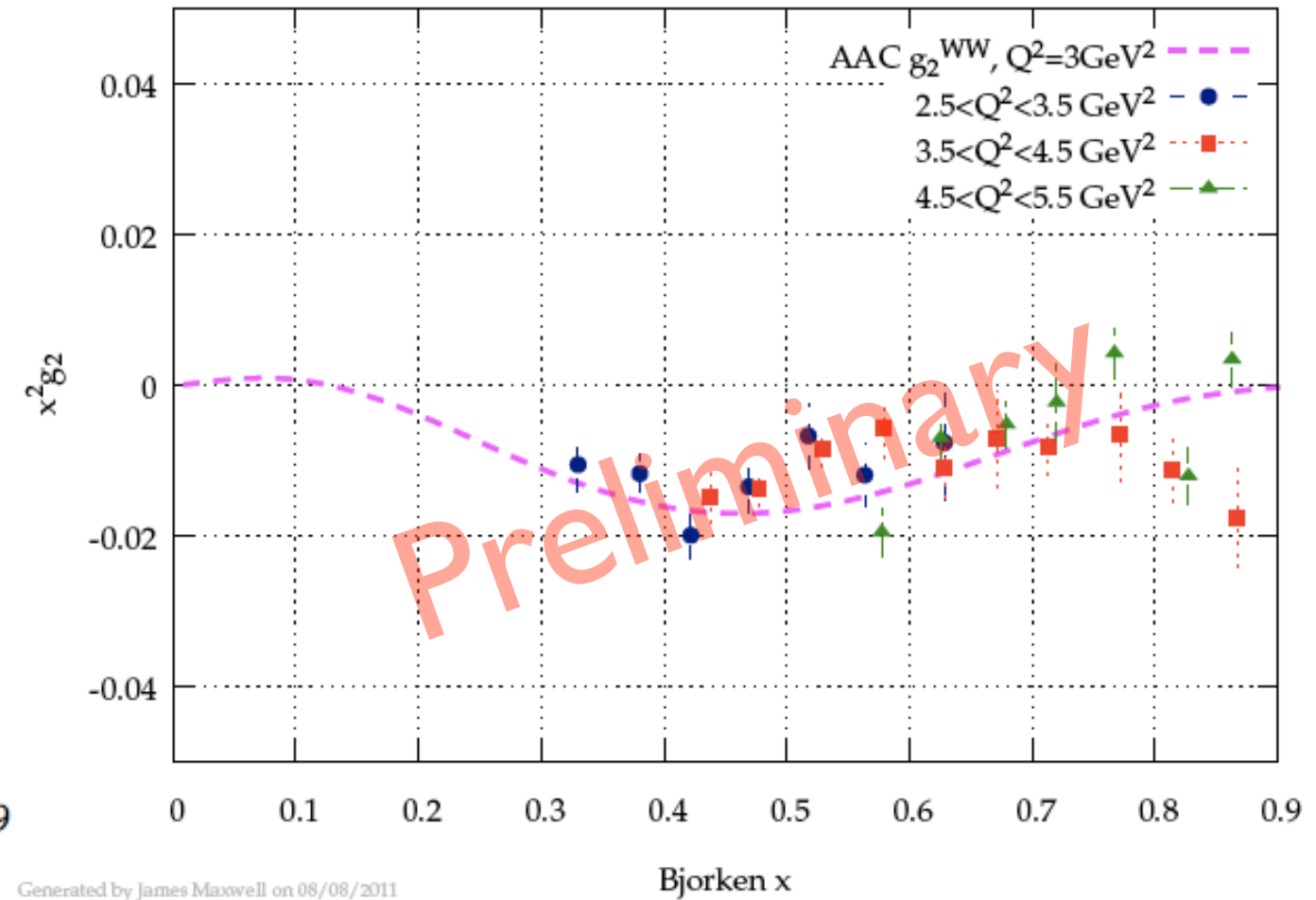
- Measure proton g_2 structure function at low Q^2 region ($0.02\text{--}0.2\text{GeV}^2$) for the first time
- Will help to clarify several puzzles:
 - Test the Burkhardt–Cottingham (BC) Sum Rule at low Q^2
 - Extract the generalized longitudinal–transverse spin polarizability δ_{LT} to give a test for Chiral Perturbation Theory (χ PT)
 - Improve the calculation of Proton Hyperfine Splitting
 - Proton charge radius from μP Lamb shift disagrees with $e\text{P}$ scattering result

Existing Data

SLAC



JLab SANE

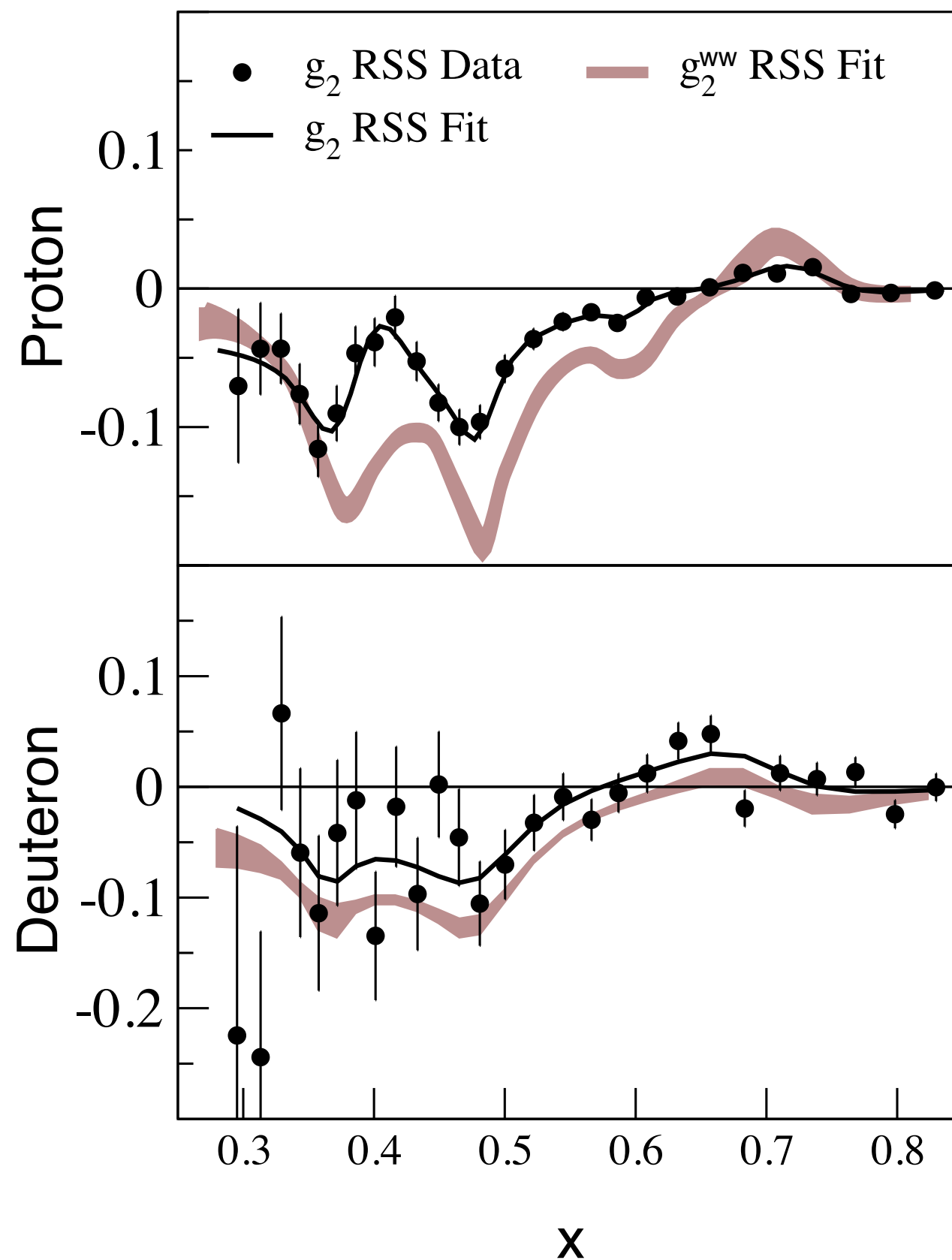


Generated by James Maxwell on 08/08/2011

SLAC: $Q^2 \sim 5 \text{ GeV}^2$

JLab SANE: $Q^2 \sim 3 \sim 6 \text{ GeV}^2$

Existing Data



JLab RSS: $Q^2 \sim 1.3 \text{ GeV}^2$

K. Slifer et al, arXiv:0812.0031

BC Sum Rule

- BC Sum Rule:

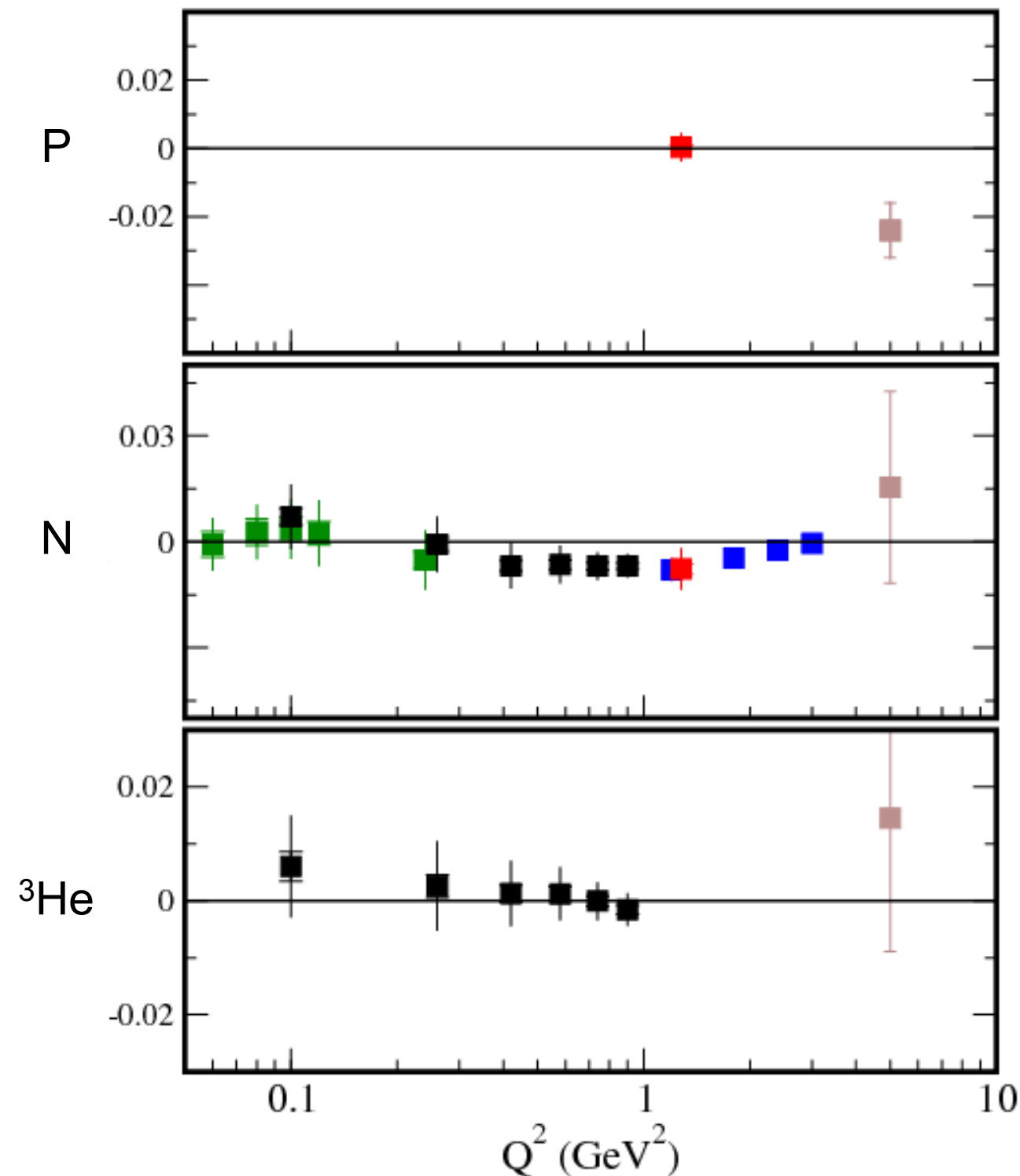
$$\int_0^1 g_2(x, Q^2) dx = 0$$

H. Burkhardt and W. N. Cottingham, *Annals. Phys.*, 56(1970)453

- BC Sum Rule will fail if g_2 :
 - exhibits non-Regge behavior at low x
 - exhibits a delta function singularity at $x=0$

R. L. Jaffe and X.-D. Ji, *Phys. Rev. D*, 43(1991)724

BC Sum Rule



- SLAC E155x
- Hall C RSS
- Hall A E94-010
- Hall A E97-110 (preliminary)
- Hall A E01-012 (preliminary)

- BC satisfied within errors for Neutron and ^3He
- Mostly unmeasured for proton

Generalized Longitudinal-Transverse Polarizability

- Start from forward spin-flip doubly-virtual Compton scattering (VVCS) amplitude g_{TT} and g_{LT}

$$\text{Re}[g_{TT}^{\text{non-pole}}(\nu, Q^2)] = \frac{\nu}{2\pi^2} \mathcal{P} \int_{\nu_\pi}^{\infty} \frac{d\nu' K}{\nu'^2 - \nu^2} \sigma_{TT}(\nu', Q^2)$$

$$\text{Re}[g_{LT}^{\text{non-pole}}(\nu, Q^2)] = \frac{1}{2\pi^2} \mathcal{P} \int_{\nu_\pi}^{\infty} \frac{d\nu' \nu' K}{\nu'^2 - \nu^2} \sigma_{LT}(\nu', Q^2)$$

- g_{TT} and g_{LT} can be expanded in power series of ν

$O(\nu^3)$ term of g_{TT} leads to the generalized forward spin polarizability γ_0

$$\begin{aligned} \gamma_0(Q^2) &= \frac{1}{2\pi^2} \int_{\nu_\pi}^{\infty} \frac{K(\nu, Q^2)}{\nu} \frac{\sigma_{TT}(\nu, Q^2)}{\nu^3} d\nu \\ &= \frac{16\alpha M^2}{Q^6} \int_0^{x_0} x^2 [g_1 - \frac{4M^2}{Q^2} x^2 g_2] dx \end{aligned}$$

$O(\nu^2)$ term of g_{LT} leads to the generalized longitudinal-transverse polarizability δ_{LT}

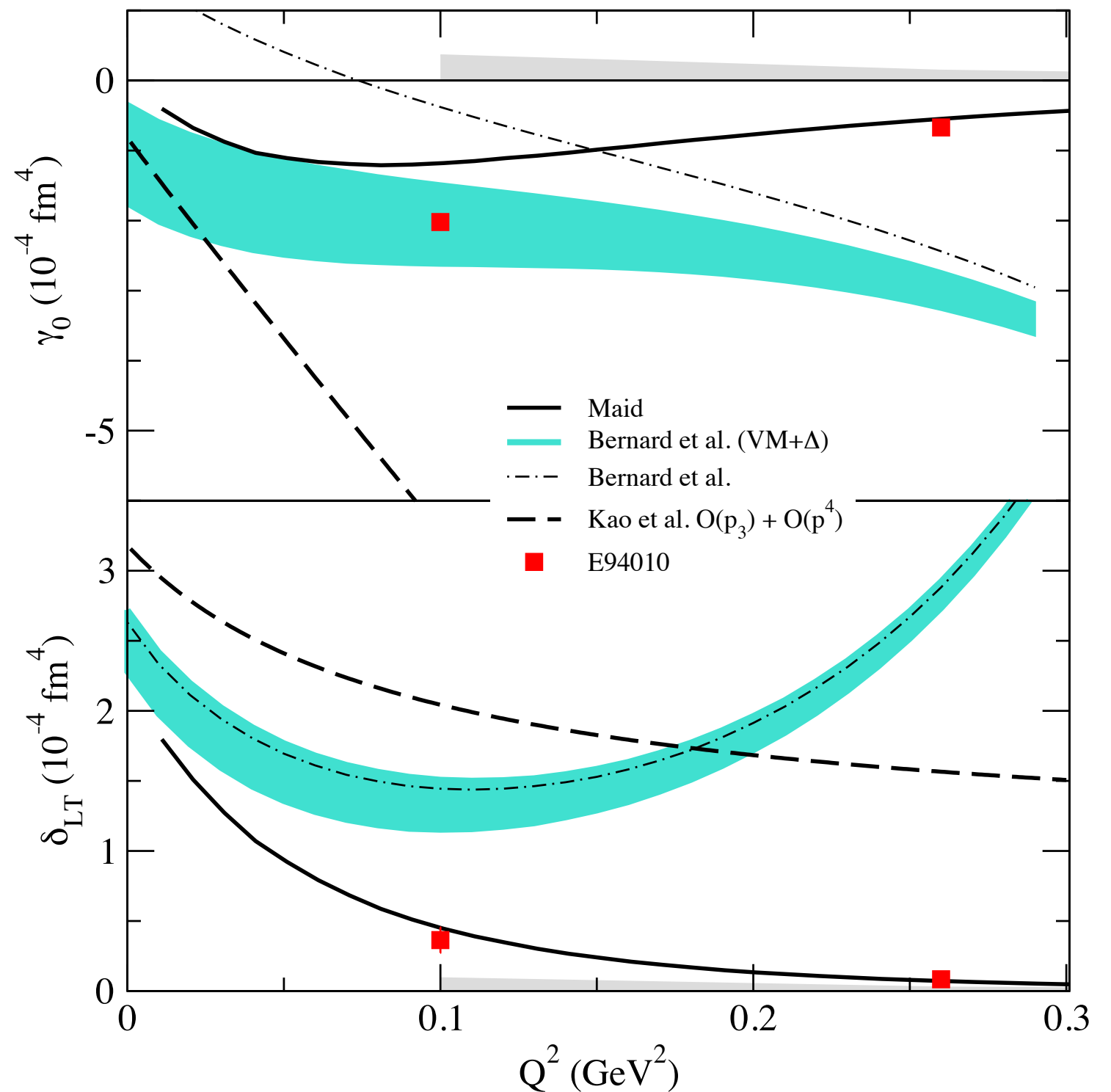
$$\begin{aligned} \delta_{LT}(Q^2) &= \frac{1}{2\pi^2} \int_{\nu_\pi}^{\infty} \frac{K(\nu, Q^2)}{\nu} \frac{\sigma_{LT}(\nu, Q^2)}{Q\nu^2} d\nu \\ &= \frac{16\alpha M^2}{Q^6} \int_0^{x_0} x^2 [g_1 + g_2] dx \end{aligned}$$

δ_{LT} puzzle

- At low Q^2 , the generalized polarizabilities have been evaluated with NLO χ PT calculations:
 - Relativistic Baryon χ PT (V. Bernard, T. Hemmert and Ulf-G. Meissner, [Phys. Rev. D, 67\(2003\)076008](#))
 - Heavy Baryon χ PT (C.W. Kao, T. Spitzenberg and M. Vanderhaeghen, [Phys. Rev. D, 67\(2003\)016001](#))
- One issue in the calculation is how to properly include the nucleon resonance contributions, especially the Δ resonance
 - γ_0 is sensitive to resonances
 - δ_{LT} is insensitive to the Δ resonance
- δ_{LT} should be more suitable than γ_0 to serve as a testing ground for the chiral dynamics of QCD

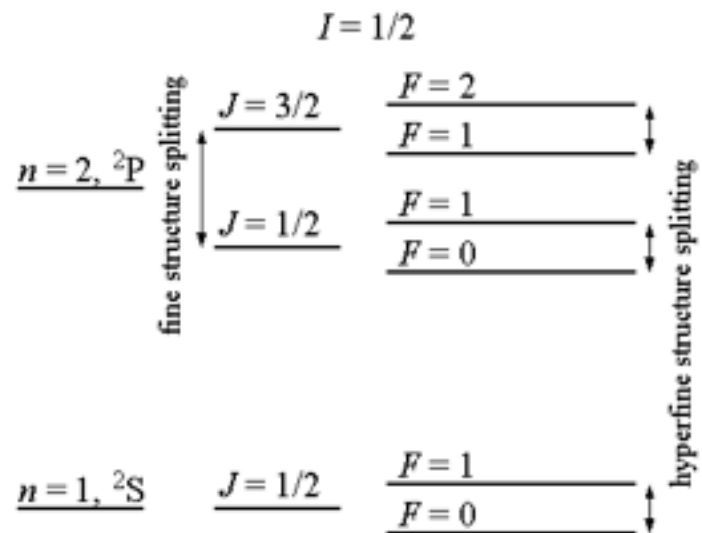
δ_{LT} puzzle

- Neutron Data shows a large deviation from the χ PT calculations
- No proton data yet
- This experiment will provide a test with proton data



Hydrogen Hyperfine Structure

- Hydrogen hyperfine splitting in the ground state has been measured to a relative high accuracy of 10^{-13}



$$\Delta E = 1420.4057517667(9)\text{MHz}$$

$$= (1 + \delta)E_F$$

$$\delta = (\delta_{\text{QED}} + \delta_R + \delta_{\text{small}}) + \Delta_S$$

- Δ_S is the proton structure correction and has the largest uncertainty

$$\Delta_S = \Delta_Z + \Delta_{\text{pol}}$$

- Δ_Z can be determined from elastic scattering, which is $-41.0 \pm 0.5 \times 10^{-6}$
- Δ_{pol} involves contributions of the inelastic part (exciting state), and can be extracted to 2 terms corresponding to 2 different spin-dependent structure function of proton

Hydrogen Hyperfine Structure

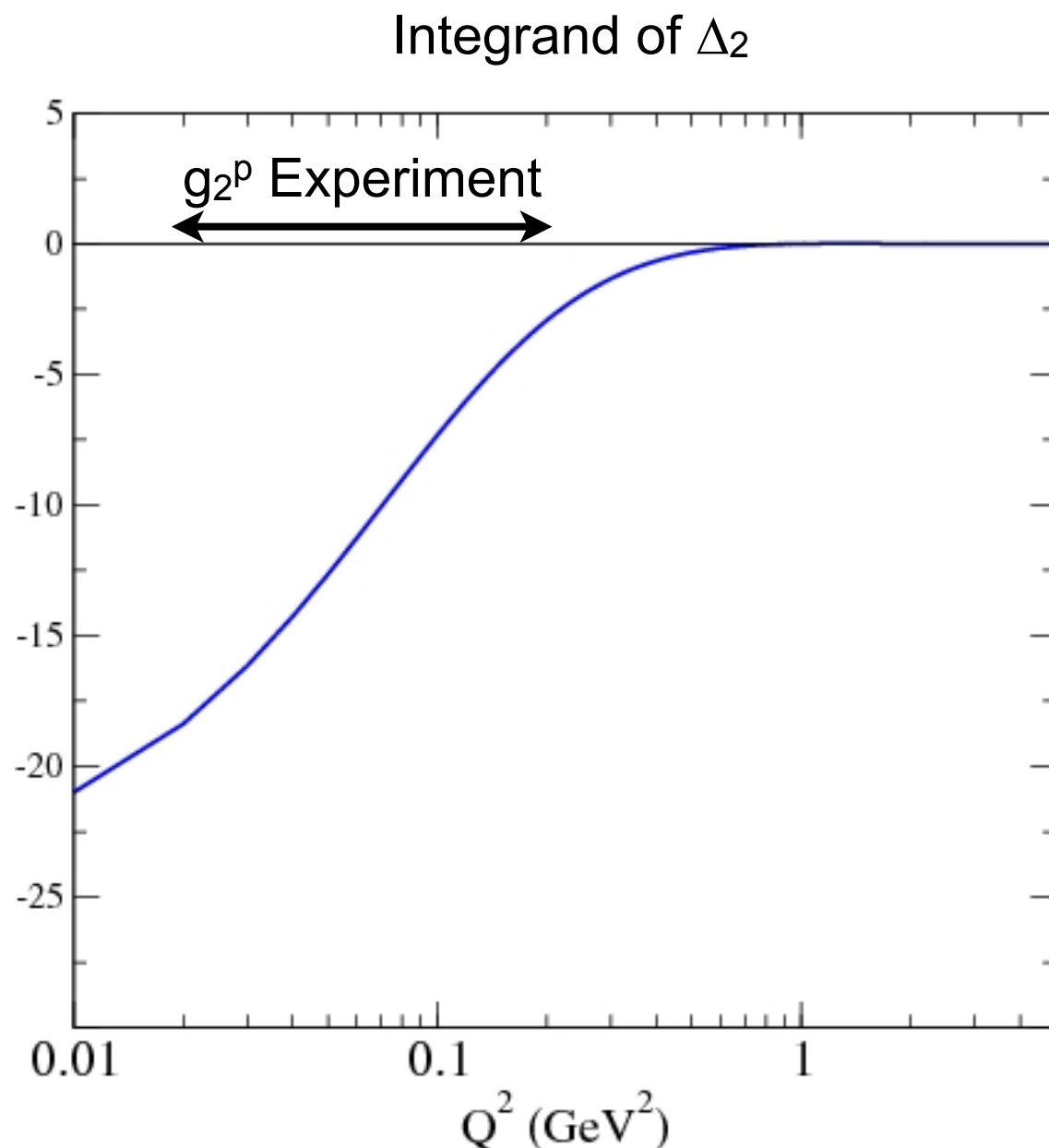
$$\Delta_{\text{pol}} = \frac{\alpha m_e}{\pi g_p m_p} (\Delta_1 + \Delta_2)$$

$$\Delta_2 = -24m_p^2 \int_0^\infty \frac{dQ^2}{Q^4} B_2(Q^2)$$

$$B_2(Q^2) = \int_0^{x_{th}} dx \beta_2(\tau) g_2(x, Q^2)$$

$$\beta_2(\tau) = 1 + 2\tau - 2\sqrt{\tau(\tau + 1)}$$

- B_2 is dominated by low Q^2 part
- g_2^p is unknown in this region, so there may be huge error when calculating Δ_2
- This experiment will provide a constraint



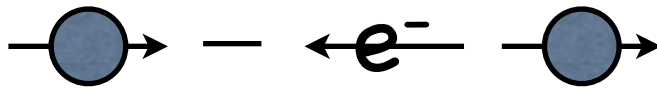
Size of the Proton

- 2 ways to measure:
 - energy splitting of the $2S_{1/2}$ - $2P_{1/2}$ level (Lamb shift)
 - scattering experiment
- The results do not match when using muonic hydrogen
 - $\langle R_p \rangle = 0.84184 \pm 0.00067 \text{ fm}$ by Lamb shift in muonic hydrogen
 - $\langle R_p \rangle = 0.87680 \pm 0.0069 \text{ fm}$ CODATA world average
- The main uncertainties originate from the proton polarizability and different values of the Zemach radius
 - This experiment will reduce the uncertainty of proton polarizability

Primary Motivation

Measure proton g_2 structure function at $0.02 < Q^2 < 0.2 \text{ GeV}^2$ region with an uncertainty of 5-7%

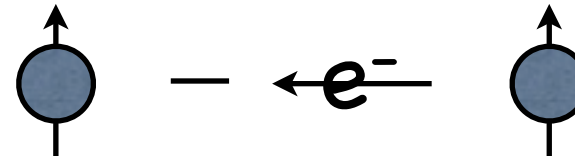
How to get g_2

$$\Delta\sigma_{\parallel} = \text{---} \text{e}^{-} \text{---} \text{---} \text{---} \text{e}^{-} \text{---}$$


$$= \frac{d^2\sigma^{\uparrow\uparrow}}{d\Omega dE'} - \frac{d^2\sigma^{\downarrow\uparrow}}{d\Omega dE'}$$

JLab Hall B experiment EG4
measured this quantity and
extracted g_1^p at low Q^2

$$= \frac{4\alpha^2 E'}{M\nu Q^2 E} [(E + E' \cos \theta) g_1 - 2Mx g_2]$$

$$\Delta\sigma_{\perp} = \text{---} \text{e}^{-} \text{---} \text{---} \text{---} \text{e}^{-} \text{---}$$


$$= \frac{d^2\sigma^{\uparrow\Rightarrow}}{d\Omega dE'} - \frac{d^2\sigma^{\downarrow\Rightarrow}}{d\Omega dE'}$$

g2p experiment will measure
this, combining the EG4 g_1^p
data to get g_2^p at low Q^2

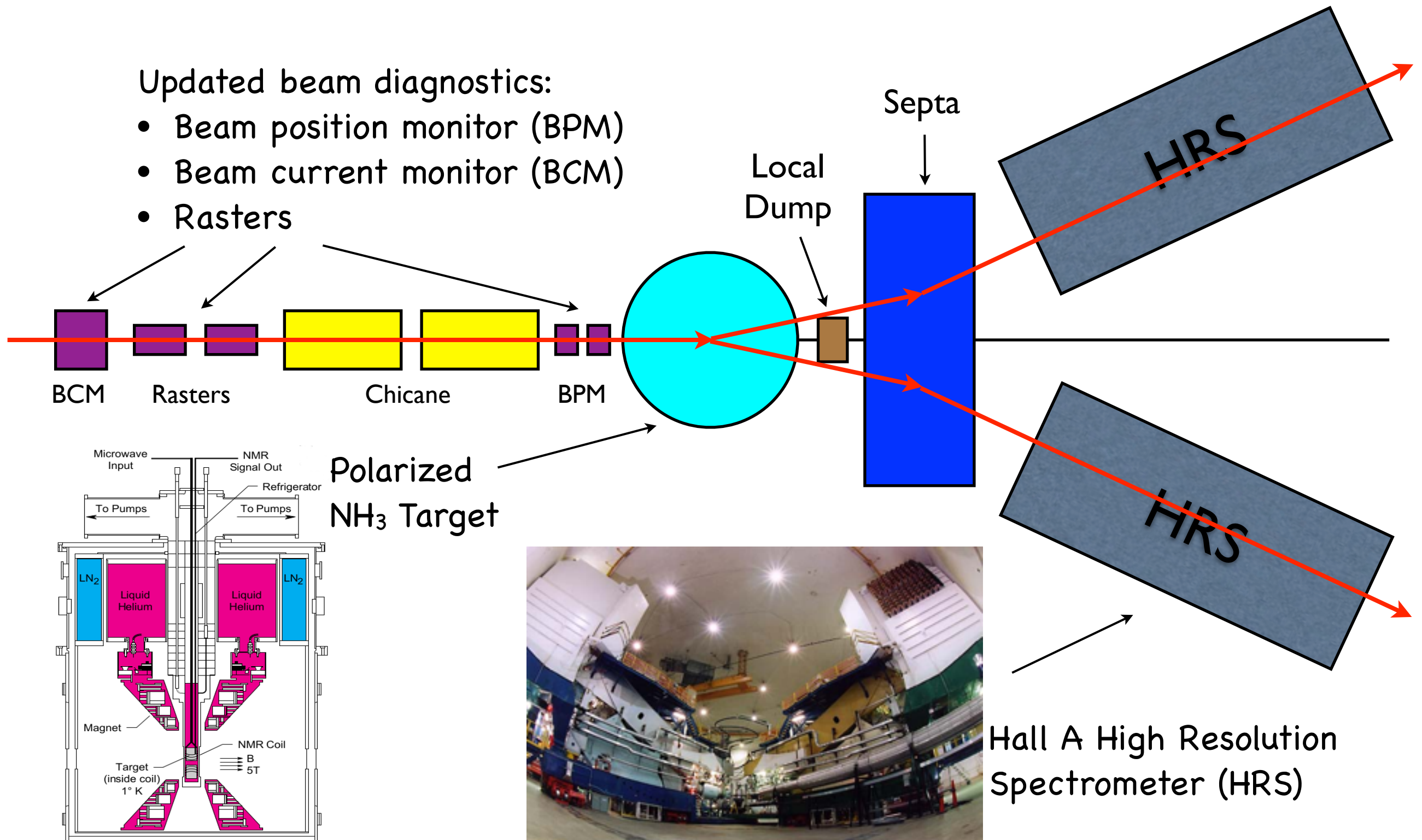
$$= \frac{4\alpha^2 E'^2}{M\nu Q^2 E} \sin \theta [g_1 + \frac{2E}{\nu} g_2]$$

Experiment Setup

Jefferson Lab Hall A

Updated beam diagnostics:

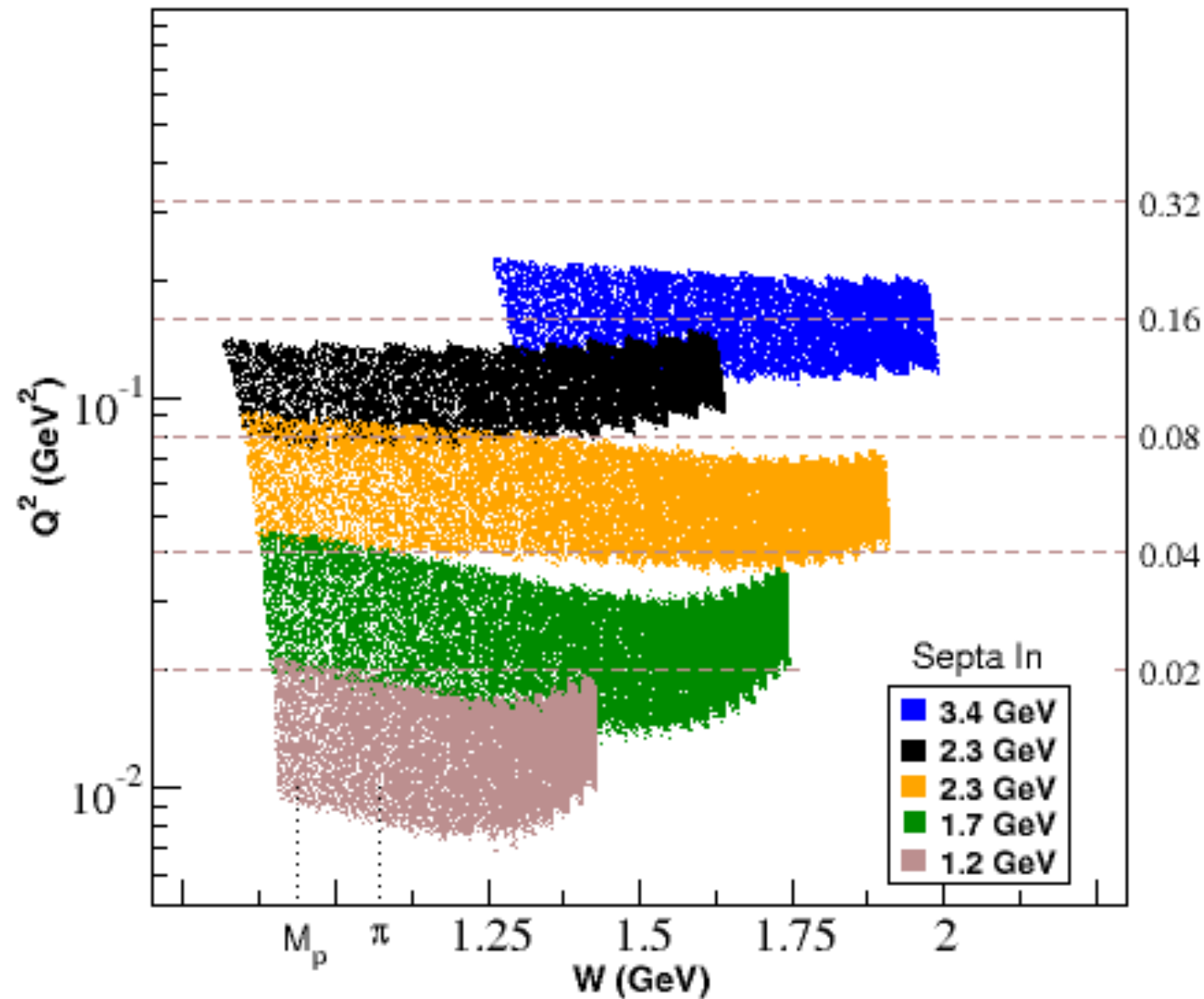
- Beam position monitor (BPM)
- Beam current monitor (BCM)
- Rasters



Experiment Setup

- Challenge: lowest possible Q^2
- Small scattering angle ($\sim 6^\circ$)
 - Use septa magnet to detect forward scattering
- Polarized NH_3 target: 2.5T~5T magnetic field
 - Use Chicane to provide an incident angle
 - Outgoing beam is not straight: use local dump
- Low current polarized beam
 - Upgrades to existing Beam Diagnostics to work at 50 nA

Kinematics Coverage



$$M_p < W < 2 \text{ GeV}$$

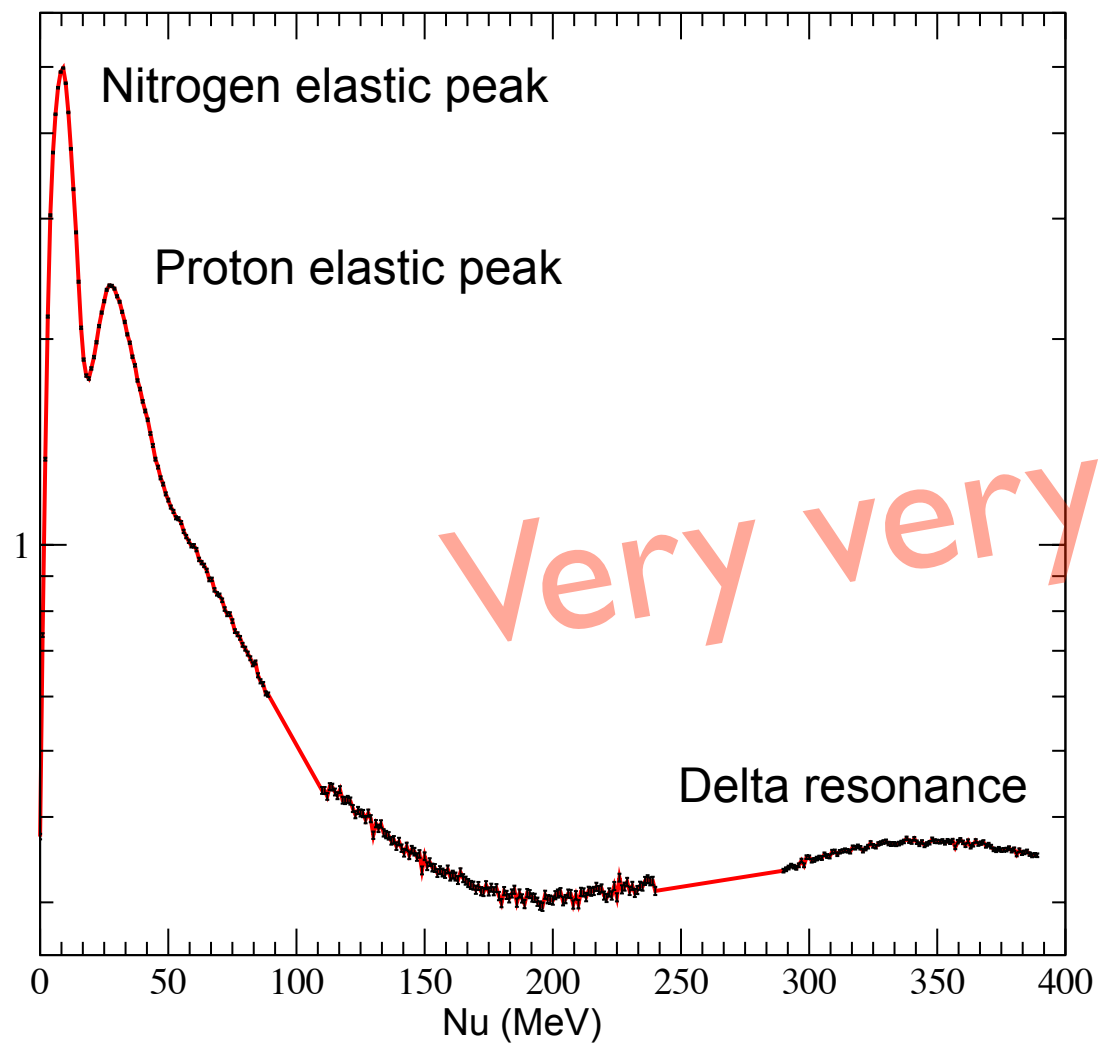
$$0.02 < Q^2 < 0.2 \text{ GeV}^2$$

- The Experiment was conducted at JLab Hall A successfully from 3/2/2012 to 5/18/2012
- Statistics:

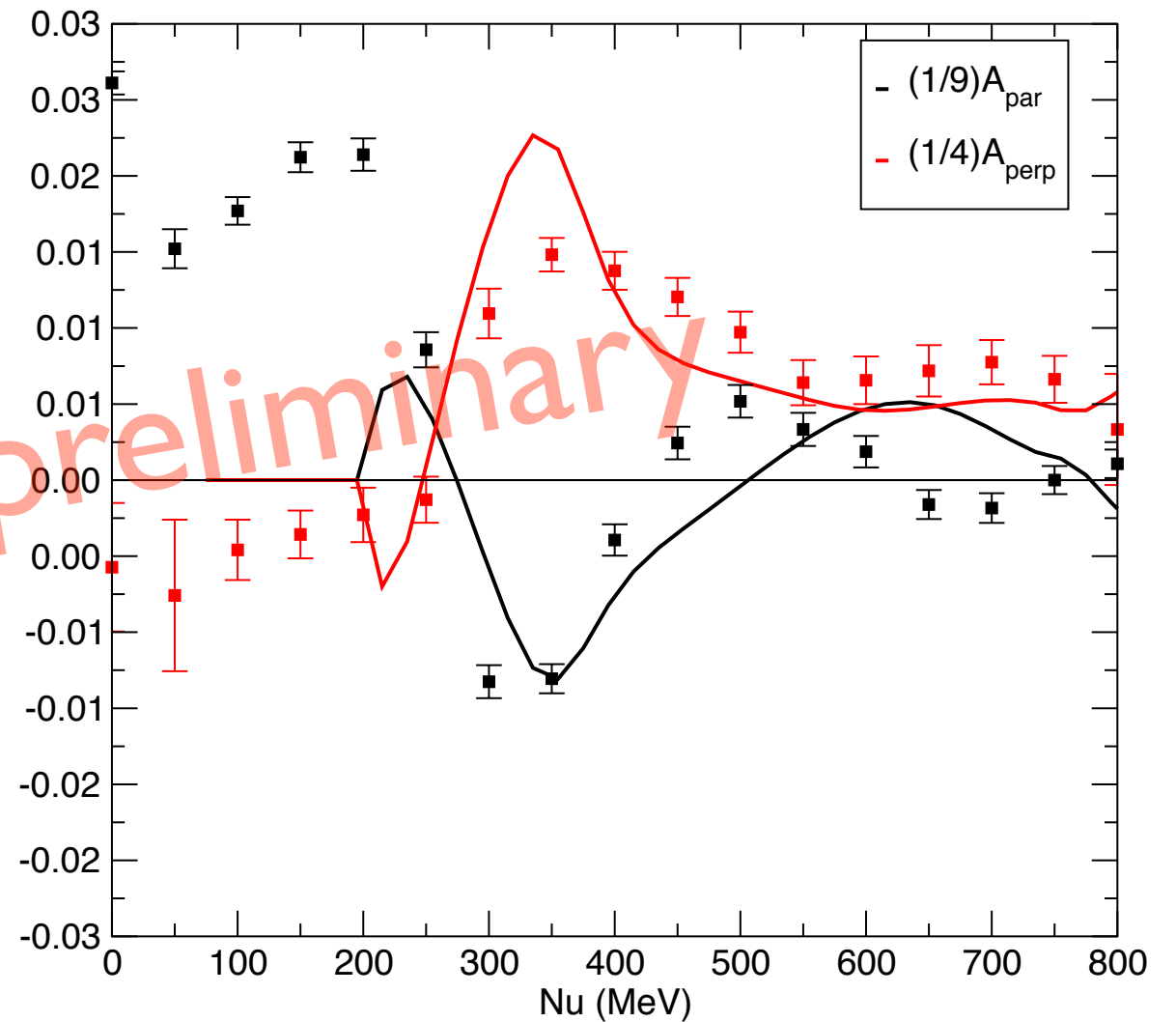
Beam Energy / GeV	Target Field /T	Recorded trigger
2.254	2.5	3.80E+09
1.706	2.5	3.20E+09
1.158	2.5	4.00E+09
2.254	5.0	7.00E+08
3.352	5.0	4.00E+08

Online results

E=2254MeV Normalized Yield



E=2254MeV Asymmetry



$$\Delta\sigma_{\perp} = \sigma_{\text{total}} \cdot A_{\perp}$$

Conclusion

- We managed to accomplish most of our physics goals
- New instruments are demonstrated working well during the experiment
- Will provide the an accurate measurement of g_2 in low Q^2 region
- Will also extract the fundamental quantities δ_{LT} to provide a test of χ PT calculations

Thanks