

# Quarks Through the Looking-Glass – New Measurement of Parity Violation in Electron-Quark Scattering

(The Jefferson Lab Hall A Collaboration)

D. Wang,<sup>1</sup> K. Pan,<sup>2</sup> R. Subedi,<sup>1,\*</sup> X. Deng,<sup>1</sup> Z. Ahmed,<sup>3</sup> K. Allada,<sup>4</sup> K. A. Aniol,<sup>5</sup> D. S. Armstrong,<sup>6</sup> J. Arrington,<sup>7</sup> V. Bellini,<sup>8</sup> R. Beminiwaththa,<sup>9</sup> J. Benesch,<sup>10</sup> F. Benmokhtar,<sup>11</sup> A. Camsonne,<sup>10</sup> M. Canan,<sup>12</sup> G. D. Cates,<sup>1</sup> J.-P. Chen,<sup>10</sup> E. Chudakov,<sup>10</sup> E. Cisbani,<sup>13</sup> M. M. Dalton,<sup>1</sup> C. W. de Jager,<sup>10,1</sup> R. De Leo,<sup>14</sup> W. Deconinck,<sup>6</sup> A. Deur,<sup>10</sup> C. Dutta,<sup>4</sup> L. El Fassi,<sup>15</sup> J. Erler,<sup>16</sup> D. Flay,<sup>17</sup> G. B. Franklin,<sup>11</sup> M. Friend,<sup>11</sup> S. Frullani,<sup>13</sup> F. Garibaldi,<sup>13</sup> A. Giusa,<sup>8</sup> A. Glamazdin,<sup>18</sup> S. Golge,<sup>12</sup> K. Grimm,<sup>19</sup> K. Hafidi,<sup>7</sup> O. Hansen,<sup>10</sup> D. W. Higinbotham,<sup>10</sup> R. Holmes,<sup>3</sup> T. Holmstrom,<sup>20</sup> R. J. Holt,<sup>7</sup> J. Huang,<sup>2</sup> C. E. Hyde,<sup>12,21</sup> C. M. Jen,<sup>3</sup> D. Jones,<sup>1</sup> Hoyoung Kang,<sup>22</sup> P. King,<sup>9</sup> S. Kowalski,<sup>2</sup> K. S. Kumar,<sup>23</sup> J. H. Lee,<sup>6,9</sup> J. J. LeRose,<sup>10</sup> N. Liyanage,<sup>1</sup> E. Long,<sup>24</sup> D. McNulty,<sup>23,†</sup> D. J. Margaziotis,<sup>5</sup> F. Meddi,<sup>25</sup> D. G. Meekins,<sup>10</sup> L. Mercado,<sup>23</sup> Z.-E. Meziani,<sup>17</sup> R. Michaels,<sup>10</sup> M. Mihovilovic,<sup>26</sup> N. Muangma,<sup>2</sup> K. E. Myers,<sup>27,‡</sup> S. Nanda,<sup>10</sup> A. Narayan,<sup>28</sup> V. Nelyubin,<sup>1</sup> Nuruzzaman,<sup>28</sup> Y. Oh,<sup>22</sup> D. Parno,<sup>11</sup> K. D. Paschke,<sup>1</sup> S. K. Phillips,<sup>29</sup> X. Qian,<sup>30</sup> Y. Qiang,<sup>30</sup> B. Quinn,<sup>11</sup> A. Rakhman,<sup>3</sup> P. E. Reimer,<sup>7</sup> K. Rider,<sup>20</sup> S. Riordan,<sup>1</sup> J. Roche,<sup>9</sup> J. Rubin,<sup>7</sup> G. Russo,<sup>8</sup> K. Saenboonruang,<sup>1,§</sup> A. Saha,<sup>10,¶</sup> B. Sawatzky,<sup>10</sup> A. Shahinyan,<sup>31</sup> R. Silwal,<sup>1</sup> S. Sirca,<sup>26</sup> P. A. Souder,<sup>3</sup> R. Suleiman,<sup>10</sup> V. Sulkosky,<sup>2</sup> C. M. Suter,<sup>8</sup> W. A. Tobias,<sup>1</sup> G. M. Urciuoli,<sup>25</sup> B. Waidyawansa,<sup>9</sup> B. Wojtsekhowski,<sup>10</sup> L. Ye,<sup>32</sup> B. Zhao,<sup>6</sup> and X. Zheng<sup>1</sup>

<sup>1</sup>University of Virginia, Charlottesville, Virginia 22904, USA

<sup>2</sup>Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA

<sup>3</sup>Syracuse University, Syracuse, New York 13244, USA

<sup>4</sup>University of Kentucky, Lexington, Kentucky 40506, USA

<sup>5</sup>California State University, Los Angeles, Los Angeles, California 90032, USA

<sup>6</sup>College of William and Mary, Williamsburg, Virginia 23187, USA

<sup>7</sup>Physics Division, Argonne National Laboratory, Argonne, Illinois 60439, USA

<sup>8</sup>Istituto Nazionale di Fisica Nucleare, Dipt. di Fisica dell'Univ. di Catania, I-95123 Catania, Italy

<sup>9</sup>Ohio University, Athens, Ohio 45701, USA

<sup>10</sup>Thomas Jefferson National Accelerator Facility, Newport News, Virginia 23606, USA

<sup>11</sup>Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA

<sup>12</sup>Old Dominion University, Norfolk, Virginia 23529, USA

<sup>13</sup>INFN, Sezione di Roma, gruppo Sanità and Istituto Superiore di Sanità, I-00161 Rome, Italy

<sup>14</sup>Università di Bari, I-70126 Bari, Italy

<sup>15</sup>Rutgers, The State University of New Jersey, Newark, New Jersey 07102, USA

<sup>16</sup>Instituto de Física, Universidad Nacional Autónoma de México, 04510 México D.F., Mexico

<sup>17</sup>Temple University, Philadelphia, Pennsylvania 19122, USA

<sup>18</sup>Kharkov Institute of Physics and Technology, Kharkov 61108, Ukraine

<sup>19</sup>Louisiana Technical University, Ruston, Louisiana 71272, USA

<sup>20</sup>Longwood University, Farmville, Virginia 23909, USA

<sup>21</sup>Clermont Université, Université Blaise Pascal, CNRS/IN2P3, Laboratoire de Physique Corpusculaire, FR-63000 Clermont-Ferrand, France

<sup>22</sup>Seoul National University, Seoul 151-742, South Korea

<sup>23</sup>University of Massachusetts Amherst, Amherst, Massachusetts 01003, USA

<sup>24</sup>Kent State University, Kent, Ohio 44242, USA

<sup>25</sup>INFN, Sezione di Roma and Sapienza - Università di Roma, I-00161 Rome, Italy

<sup>26</sup>Institut Jožef Stefan, 3000 SI-1001 Ljubljana, Slovenia

<sup>27</sup>George Washington University, Washington, District of Columbia 20052, USA

<sup>28</sup>Mississippi State University, Starkeville, Mississippi 39762, USA

<sup>29</sup>University of New Hampshire, Durham, New Hampshire 03824, USA

<sup>30</sup>Duke University, Durham, North Carolina 27708, USA

<sup>31</sup>Yerevan Physics Institute, Yerevan 0036, Armenia

<sup>32</sup>China Institute of Atomic Energy, Beijing, 102413, P. R. China

Parity symmetry, or mirror-image symmetry, implies that flipping left and right does not change the laws of physics. Violation of parity symmetry in the subatomic weak force was discovered in 1957, and parity violation in electron scattering played a key role in establishing, and now testing, the standard model of particle physics. One particular set of the quantities accessible through measurements of parity-violating observables is the electron-quark effective weak couplings, called  $C_{2q}$ 's, measured directly only once in the past 40 years. We report here a precise measurement of a specific  $C_{2q}$  coupling combination that is five times better than previous data. These results are the first evidence at more than 95% confidence level that the  $C_{2q}$ 's are non-zero as

predicted by the electroweak theory of particle physics, and lead to new constraints on interactions beyond the standard model, particularly on those caused by flipping the quark chirality.

Symmetry permeates nature and is fundamental to all laws of physics. One such example is the mirror symmetry, also called “parity symmetry”: a physical system or process is said to respect this symmetry if it behaves the same way when reflected in a mirror. Laws for electromagnetism, gravity and the subatomic strong force respect parity symmetry. But the subatomic weak force does not, as was first observed by C.S. Wu [1]. Historically, the observation of parity violation played an important role in establishing the standard model of particle physics. We report here on a new measurement of the parity-violating asymmetry in electron-quark scattering and extract electron-quark weak couplings. Results presented here improve the precision of the vector-electron axial-vector-quark coupling combination  $2C_{2u} - C_{2d}$  by a factor of five, are in agreement with theoretical predictions, and set constraints on new interactions beyond the standard model. In today’s particle physics research led by colliders such as the LHC, our results provide specific chirality information on electroweak theory that is difficult to obtain at high energies. The measurement is relatively free of ambiguity in its interpretation, and opens the door to even more precise measurements in the future.

In parity-violating electron scattering (PVES) experiments, one measures an asymmetry that can be expressed as

$$A_{PV} = \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-}, \quad (1)$$

where  $\sigma_+(\sigma_-)$  are the cross sections for scattering longitudinally polarized electrons that are right-handed (left-handed), meaning their spins are parallel (antiparallel) to the electron’s momentum. For deep inelastic scattering (DIS) from nuclear targets (DIS is defined as scattering in which the electron interacts with a single quark, independent of the surrounding quarks and gluons), this asymmetry can be written in a largely model-independent way as [2]

$$A_{PV} = \frac{G_F Q^2}{4\sqrt{2}\pi\alpha} [a_1(x, Q^2)Y_1(x, y, Q^2) + a_3(x, Q^2)Y_3(x, y, Q^2)], \quad (2)$$

where  $G_F$  is the Fermi constant,  $\alpha$  is the fine structure constant,  $Q^2 \equiv -q^2$  with  $q$  the four-momentum transferred from the electron to the target,  $x$  is the Bjorken scaling variable and describes the fraction of momentum carried by the quark struck by the electron,  $y = (E - E')/E$  is the fractional energy loss of the electron with  $E(E')$  the incident (scattered) electron energy,  $Y_{1,3}$  are kinematic factors, and the variables  $a_{1,3}$  are related to the subatomic structure of the target. (See Methods section for a complete description.) The first experiment (SLAC E122) to detect parity violation in electron scattering, led by C.Y. Prescott [3, 4], provided results that strongly favored a model attributed to Weinberg, Salam and Glashow [5–7], establishing it as the keystone of the now

highly successful standard model of particle physics. PVES has subsequently been used as a sensitive probe to study diverse physics ranging from physics beyond the Standard Model [8, 9] to the structure of both nuclei [10] and the nucleon [11].

In so-called tree-level scattering where the electron exchanges only a single photon or a single  $Z$  boson with the target, very simple expressions for  $a_{1,3}$  in Eq. (2) emerge for electron DIS from deuterium:

$$a_1 = \frac{6}{5} (2C_{1u} - C_{1d}), a_3 = \frac{6}{5} (2C_{2u} - C_{2d}). \quad (3)$$

The use of the deuterium target simplifies interpretation because it has equal numbers of up and down valence quarks. Here  $C_{1u(1d)}$  and  $C_{2u(2d)}$  are the effective weak couplings between the electrons and the up (down) quark. The subscripts 1 and 2 refer to whether the coupling to the electron or quark is vector or axial-vector in nature:  $C_{1u(d)}$  is the axial-vector-electron vector-quark (AV) coupling, i.e. it probes parity violation caused by flipping of the electron chirality;  $C_{2u(d)}$  is the vector-electron axial-vector-quark (VA) coupling that is sensitive to parity violation due to flipping of the quark chirality. In the standard model the  $C_1$  and the  $C_2$  couplings are proportional to the quark vector and axial weak charges, respectively. In testing the standard model, it is important to determine all four  $C_{1u,1d,2u,2d}$  as accurately as possible, since new interactions could manifest itself in either set of couplings. Experimentally, one could extract both  $2C_{1u} - C_{1d}$  and  $2C_{2u} - C_{2d}$  by measuring asymmetries at different  $Y_{1,3}$  values in the DIS regime. However, a precise determination of  $2C_{2u} - C_{2d}$  is difficult because of its small value ( $-0.095$ ), as opposed to  $2C_{1u} - C_{1d} = -0.719$ , in the standard model.

The new measurement reported here was performed using the electron beam at Thomas Jefferson National Accelerator Facility (JLab), in Virginia, USA. A 100- $\mu$ A, nearly 90%-longitudinally-polarized electron beam was incident on a 20-cm-long liquid deuterium target held at a temperature of 22 K. Scattered particles were detected in a pair of spectrometers that determine the momentum and the direction of the detected particles to high precisions [12]. To directly access  $C_{2u,2d}$ , the kinematics were chosen so that the bulk of the detected electrons emerged from the target after undergoing a DIS interaction. In contrast, all other PVES experiments after SLAC E122 were performed outside the DIS regime, thus could not provide clean information on  $C_{2q}$ .

The size of the asymmetry expected for this measurement is at the level of  $10^{-4}$ . The major challenge comes from the combination of the high electron event rate, and the high pion background typical of DIS measurements. This was overcome by the use of a custom electronic and data acquisition (DAQ) system with built-in pion rejection capability [13]. The DAQ successfully counted electrons, event-by-event, at rates up to 600 kHz. The relative uncertainty in the measured asymme-

tries due to pion background was below  $5 \times 10^{-4}$ , and that due to counting deadtime was below 0.4%. The leading systematic effect comes from normalizing by the electron beam polarization, which had a relative uncertainty at the level of (1.2 – 1.8)%. On the other hand, beam stability was not a significant issue because of recent advances in the monitoring and feedback control of the beam, a direct outcome of some of the earlier PVES studies [8–11].

The high intensity of the JLab beam allowed the completion of the experiment in just under two months. A total of 35.4 million electrons were counted at two DIS kinematics. The asymmetry measured at  $E = 6.067$  GeV,  $\langle x \rangle = 0.241$ ,  $Y_1 = 1.0$ ,  $Y_3 = 0.44$  and  $\langle Q^2 \rangle = 1.085$  (GeV/c)<sup>2</sup> is

$$A_{\text{exp}} = [-91.1 \pm 3.1(\text{stat.}) \pm 3.0(\text{syst.})] \times 10^{-6}, \quad (4)$$

where  $\langle x \rangle$  and  $\langle Q^2 \rangle$  are averaged over the spectrometer acceptance. This is to be compared with the standard model (SM) expectation of  $A_{\text{SM}} = -87.74 \times 10^{-6}$ . To allow an extraction of  $C_{1u,1d}$  and  $C_{2u,2d}$ , it is necessary to express the asymmetry in terms of these couplings. This was calculated using the MSTW2008 leading-order parametrization [14] of parton distribution functions (PDF). For the kinematics above, it gives  $A_{\text{SM}} = (1.156 \times 10^{-4}) [(2C_{1u} - C_{1d}) + 0.348(2C_{2u} - C_{2d})]$ . The second DIS kinematic point was  $E = 6.067$  GeV,  $\langle x \rangle = 0.295$ ,  $Y_1 = 1.0$ ,  $Y_3 = 0.69$ ,  $\langle Q^2 \rangle = 1.901$  (GeV/c)<sup>2</sup>, and the result is

$$A_{\text{exp}} = [-160.8 \pm 6.4(\text{stat.}) \pm 3.1(\text{syst.})] \times 10^{-6}. \quad (5)$$

The standard model expectation is  $A_{\text{SM}} = -158.90 \times 10^{-6}$ , with the coupling sensitivity  $A_{\text{SM}} = (2.022 \times 10^{-4}) [(2C_{1u} - C_{1d}) + 0.594(2C_{2u} - C_{2d})]$ . Details of the standard model calculation and the uncertainty due to PDF fits are given in the Methods section.

Using the most recent world data for the coupling  $C_{1u,1d}$  [15], obtained from PVES and cesium atomic parity violation experiments [16–19], a simultaneous fit of  $2C_{1u} - C_{1d}$  and  $2C_{2u} - C_{2d}$  to our results and to the asymmetries from SLAC E122 was performed, yielding

$$(2C_{2u} - C_{2d})|_{Q^2=0} = -0.145 \pm 0.066(\text{total}), \quad (6)$$

where the total uncertainty includes those from PDF fits and higher-order radiative corrections. The zero- $Q^2$  values  $C_{2u,2d}|_{Q^2=0}$  are called  $g_{VA}^{eu,ed}$  (and similarly  $C_{1u,1d}|_{Q^2=0}$  are  $g_{AV}^{eu,ed}$ ) in Ref. [20], where certain electroweak radiative corrections have been applied such that the values in Eq. (6) can be compared directly to results from other precision experiments using different kinds of processes. The values for  $C_{2u,2d}|_{Q^2=0}$  differ from the non-zero  $Q^2$  values accessed in this experiment by 0.002–0.003 for both the up and the down quark, at both kinematics.

In the standard model, the  $C_{1,2}$  are related to the weak mixing angle  $\theta_W$ , an important ingredient of the Weinberg-Salam-Glashow theory. The weak mixing angle extracted

from our asymmetry results, evolved to the mass of the  $Z$  boson in the modified minimal subtraction ( $\overline{\text{MS}}$ ) scheme, is  $\hat{s}_Z^2 \equiv \sin^2 \theta_W(Q^2=M_Z^2, \overline{\text{MS}}) = 0.2299 \pm 0.0043$ . This can be compared to the latest standard-model fit to world data  $\hat{s}_Z^2 = 0.23126 \pm 0.00005$ .

The result in Eq. (6) is to be compared with the standard model prediction  $2C_{2u} - C_{2d}|_{Q^2=0} = -0.0949$  as shown in Fig. 1. Our results have greatly improved the uncertainty

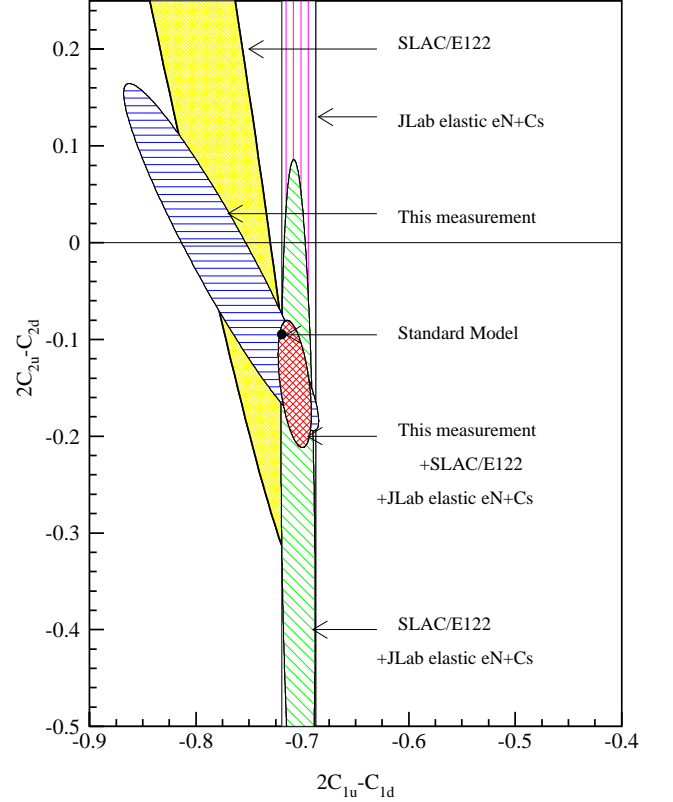


FIG. 1: Results for  $(2C_{1u} - C_{1d})|_{Q^2=0}$  and  $(2C_{2u} - C_{2d})|_{Q^2=0}$  obtained from this experiment alone (blue horizontal-line-hatched ellipse) compared with SLAC E122 (yellow ellipse) [3, 4]. The latest data on  $C_1$  [15] (obtained from PVES and Atomic Cs [16–19]) is shown as the magenta vertical-line-hatched band. The green slanted-line-hatched ellipse shows the combined result of SLAC E122 and the latest  $C_1$ , while the red diamond-line-hatched ellipse shows the combined result of SLAC E122, this experiment, and the latest  $C_1$ . The standard model value is shown as the black dot. Note that the same scale is used for the two axes to illustrate the significant difference in our knowledge of the  $C_1$  and  $C_2$  couplings.

on the effective vector-electron axial-quark weak couplings  $C_{2u,2d}$  and are in good agreement with the standard model prediction. This is also the first direct measurement of the coupling combination  $2C_{2u} - C_{2d}$  that shows a deviation from zero. We note that evidence for nonzero values of the  $C_{2u,2d}$ , perhaps a different combination from what we measured, may have also been observed in experiments measuring the nucleon axial form factors [21]. However, extraction of  $C_{2u,2d}$  from the nucleon axial form factor is model-dependent, while

in DIS the electron probes quarks unambiguously. The directness of our approach is an important feature that makes it possible to reach significantly higher accuracy in the future.

A comparison of the present result with the standard model predictions can be used to set mass limits  $\Lambda$  below which new interactions are unlikely to occur. For the case of electron and quark compositeness and contact interactions, we used the convention of Ref. [22] and the procedure in Ref. [23]. The limit for the constructive (destructive) interference contribution to the standard model is:

$$(\Lambda)^\pm = v \left[ \frac{8\sqrt{5}\pi}{|(2C_{2u} - C_{2d})_{Q^2=0}|^\pm} \right]^{1/2}, \quad (7)$$

where  $|(2C_{2u} - C_{2d})_{Q^2=0}|^\pm$  is the difference between the standard model value and the upper (lower) confidence bound extracted from the data,  $v = \sqrt{2}/(2G_F) = 246.22$  GeV is the Higgs vacuum expectation value setting the electroweak scale, and the  $\sqrt{5}$  is a normalization factor taking into account the coefficients of the  $C_{2u,2d}$  in the denominator. For a 95% confidence level, we extracted

$$\Lambda^+ = 5.8 \text{ TeV and } \Lambda^- = 4.6 \text{ TeV}, \quad (8)$$

for constructive and destructive beyond-the-standard-model physics. Figure 2 illustrates these limits. The limits set by  $C_{1u,1d}$  are determined mostly by previous PVES and cesium atomic-parity-violation results, but this experiment has clearly improved the limits set by  $C_{2u,2d}$ .

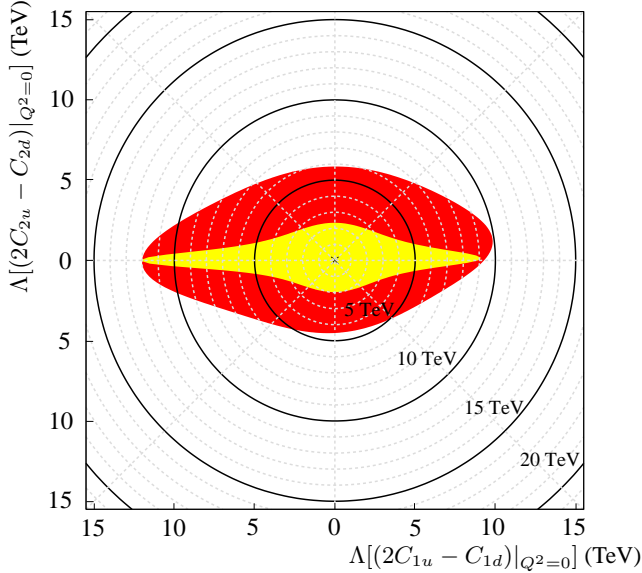


FIG. 2: Mass exclusion limits on the electron and quark compositeness and contact interactions obtained from the zero- $Q^2$  values of  $2C_{1u} - C_{1d}$  and  $2C_{2u} - C_{2d}$  at the 95% confidence level. The yellow contour shows the limit obtained from SLAC E122 asymmetry results [3, 4] combined with the best  $C_1$  values [15]. The red contour shows the limit with our new results added.

The strength of our results reported here is that they isolate a well-defined combination of the electron-quark contact

interaction. We note that mass limits on the electron-quark contact interactions have been published by ZEUS [24] and H1 [25] at HERA. They find  $\Lambda^+ = 3.3$  TeV and  $\Lambda^- = 3.2$  TeV [24],  $\Lambda^+ = 3.8$  TeV and  $\Lambda^- = 3.6$  TeV [25] on the vector-electron axial-vector-quark term. Similar limits have been published by ATLAS [26] at LHC in the left-left isoscalar model, which are  $\Lambda^+ = 9.5$  TeV and  $\Lambda^- = 12.1$  TeV [27]. The HERA and the LHC measurements are sensitive to several different vector and axial-vector coupling combinations, thus their limits were obtained with the assumption that besides the particular chirality combination used in the model, all other contact interactions are zero. This assumption is unnecessary for the extraction of mass limits from our results.

In summary, we have measured to a high precision the parity-violating asymmetry of electron-deuteron deep inelastic scattering and improved our knowledge of the effective electron-quark weak coupling combination  $2C_{2u} - C_{2d}$  by a factor of five. The results provide the first direct and unambiguous evidence that  $2C_{2u} - C_{2d}$  differs from zero at more than the 95% confidence level. They significantly improve the limits on certain types of new interactions in which it is the chirality of the quarks that is responsible for the observed parity violation.

## METHODS SUMMARY

The experiment was carried out at the Thomas Jefferson National Accelerator Facility (JLab). A longitudinally polarized electron beam scattered from an unpolarized deuterium target. Scattered particles were detected by the high resolution spectrometer pair equipped with a custom electronic and data acquisition system which separated electrons from background particles [13]. From the detected counts  $C$  and the beam intensity  $I$  integrated over periods of stable beam helicity, we computed a ratio  $\sigma = \frac{C}{I}$  that is proportional to the scattering probability, and from these ratios we computed the parity-violating asymmetry  $A_{\text{exp}}$  between right- and left-handed electrons. Two kinds of corrections were then made to the asymmetries: an overall normalization factor and a possible systematic shift due to false asymmetries arising from backgrounds or helicity correlations in the beam parameters. The normalization factors include the beam polarization (measured by a Møller and a Compton polarimeter), measurements of scattered-electron kinematics, electromagnetic radiative corrections, and effects from two-photon exchange between the electron and target. The false-asymmetry corrections were all very small compared to our statistical error and included an evaluation of helicity correlations in beam current, position, and energy and backgrounds such as pions, scattering from the target aluminum windows, or rescattering inside the spectrometers. A summary of all corrections and the asymmetry results are presented in Table I of Supplementary Information.

To calculate the standard model expectation of the mea-

sured asymmetry and its sensitivity to  $2C_{1u} - C_{1d}$  and  $2C_{2u} - C_{2d}$ , we used parton distribution functions (PDFs) to calculate the structure functions in  $a_{1,3}$ . Three PDF fits were used. Results of the calculation are shown in Table II of Supplementary Information. The variation among all three fits is below  $10^{-6}$  for the asymmetry; therefore the uncertainty in the extracted  $C_{1,2}$  due to PDF fits is quite small. Effects from interactions among quarks inside the target, called “higher-twist effects”, were evaluated by comparing the measured asymmetries at the two DIS kinematics reported here. We found that our results on  $C_{2u,2d}$  are largely not affected by this effect at the present precision.

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\* now at Richland College, Dallas County Community College District, Dallas, Texas 75243, USA.

† now at Idaho State University, Pocatello, Idaho 83201, USA.

‡ now at Rutgers, The State University of New Jersey, Newark, New Jersey 07102, USA.

§ now at Kasetsart University, Bangkok 10900, Thailand.

¶ deceased.

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**Supplementary Information** is linked to the online version of the paper at [www.nature.com/nature](http://www.nature.com/nature).

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**Competing Interests** The authors declare that they have no competing financial interests.

**Correspondence** Correspondence and requests for materials should be addressed to X. Zheng (email: xz5y@virginia.edu).