

Motivation

- Generalized Parton Distributions (GPDs) provide 3D imaging of the nucleon, unifying Parton distributions and form factors (FF).
- Timelike Compton Scattering (TCS) is a process sensitive to GPDs via interference with Bethe-Heitler (BH)[1].
- Offers unique access to the *real part* of Compton Form Factors (CFFs).

Experimental Setup

- TCS description: $\gamma p \rightarrow \gamma^* p' \rightarrow e^+ e^- p'$
- Real photon beam produced via bremsstrahlung off 11.5 GeV electron beam in Hall D Jlab.
- γ interact with the proton target and the product are detected by the GlueX Detector
- Excellent vertex and momentum resolution allows precise invariant mass reconstruction.

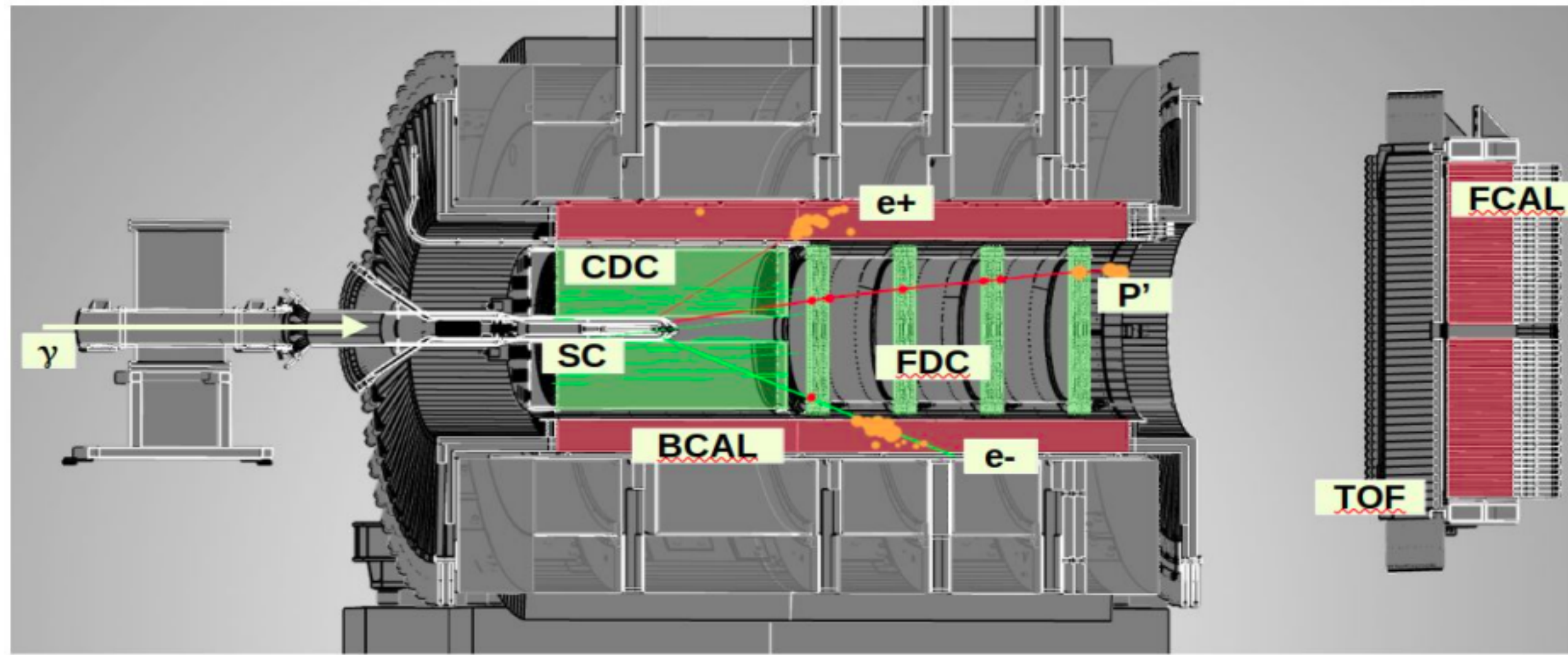


Figure 1. GlueX Detector Schematic.

TCS and BH Interference [2]

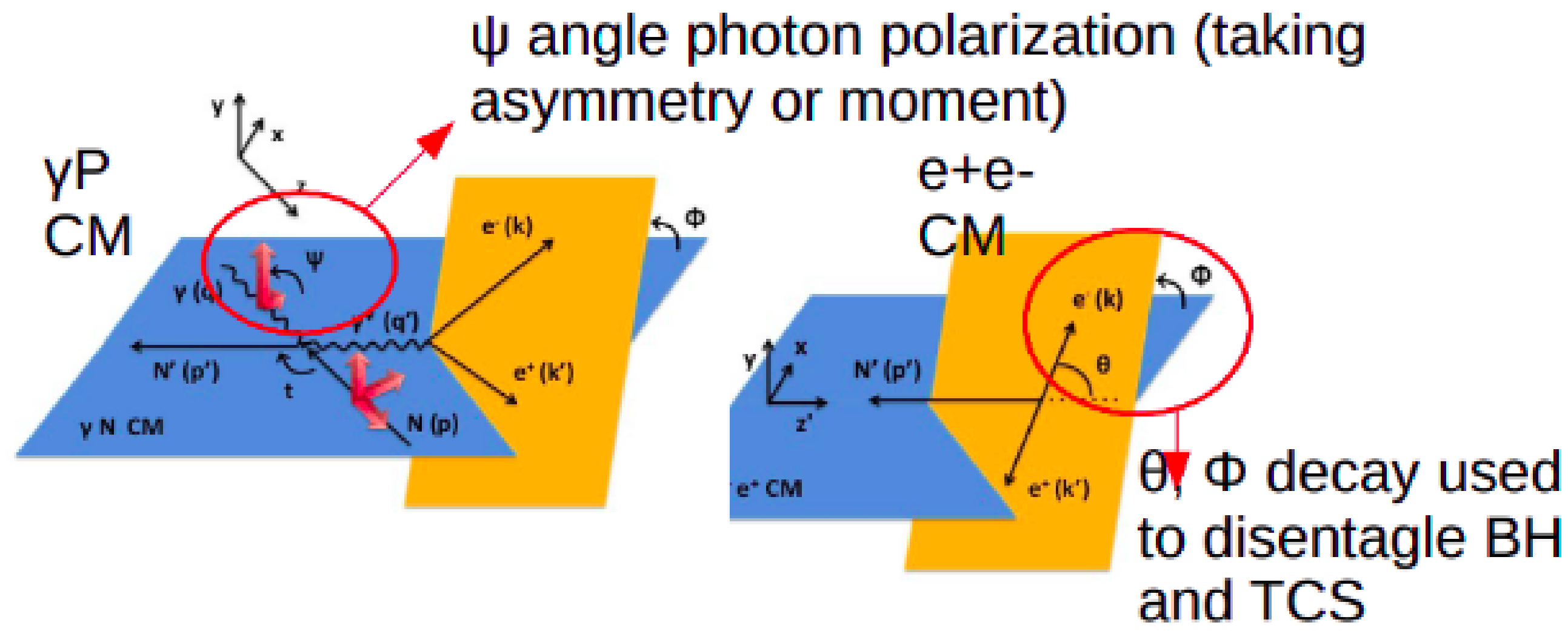


Figure 2. Frame of reference for studying TCS and BH.

The **Analysis Goal** is to extract Compton Form Factors from asymmetries. The interference term $\propto \text{Re}(\mathcal{H})$ and can be isolated via angular asymmetries.

This enable model-independent constraint of $H(x, \xi, t)$ at skewness $\xi = \frac{Q^2}{2s-Q^2}$, parton momentum fraction x and momentum transfer t .

A **step towards that goal** is to study the χ^2 dependence of some of the TCS/BH event selection kinematic terms.

Events spectrum

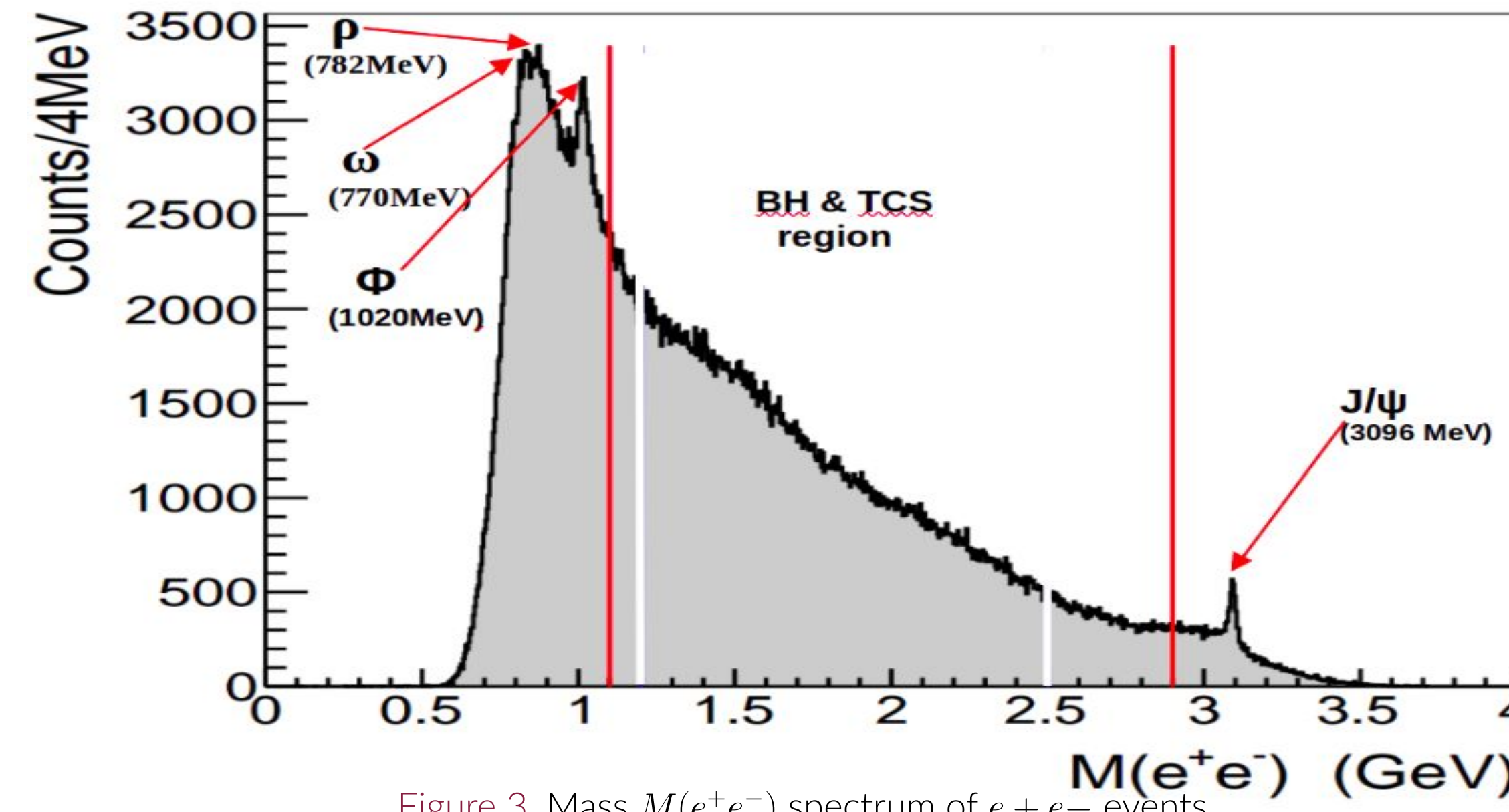


Figure 3. Mass $M(e^+e^-)$ spectrum of e^+e^- events.

BH & TCS region: $1.2 \leq M(e^+e^-) \leq 2.5$ GeV Lies away from many meson resonances.

χ_{kin}^2 dependence of Event Selection Variables

Cuts Applied: the cuts variables are ideally zero for Exclusive events.

$$M_{miss}^2 = (P_\gamma^\mu + p_p^\mu - p_{\gamma^*}^\mu - p_{p'}^\mu)^2, \quad |M_{miss}^2| \leq 0.04 \text{ GeV}^2$$

$$\Delta P_\perp = \vec{P}_{p'}_\perp - \vec{P}_{\gamma^*}_\perp, \quad |\Delta P_\perp| \leq 0.05 \text{ GeV}$$

$$\Delta \Phi = \Phi_{\vec{P}_{p'}} - \Phi_{\vec{P}_{\gamma^*}}, \quad |\Delta \Phi| \leq 0.15 \text{ rad}$$

$$P_T^{imb} = |\vec{P}_{\perp \gamma^*} + \vec{P}_{\perp p'}|, \quad P_T^{imb} \leq 0.125 \text{ GeV}/c$$

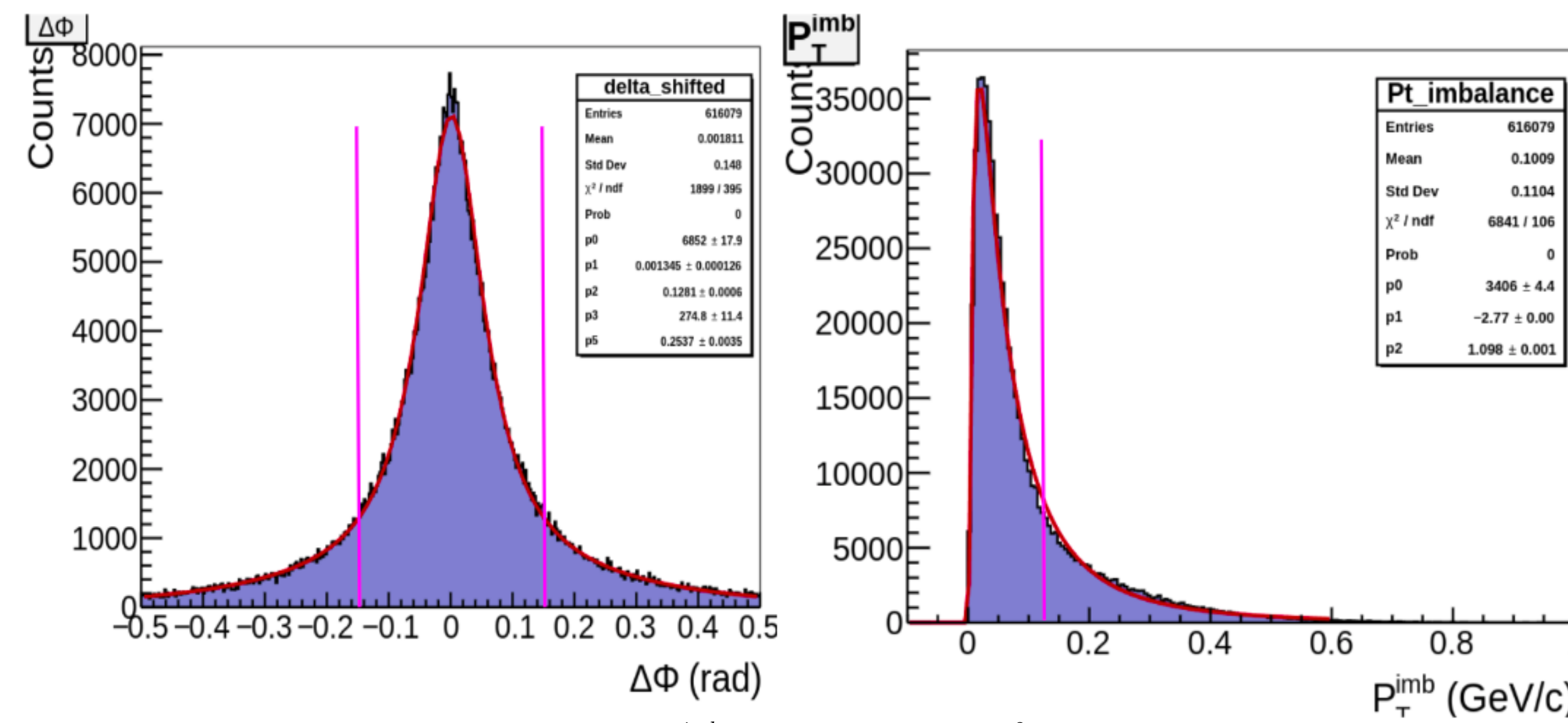


Figure 4. $\Delta \Phi$ & P_T^{imb} counts for a given χ_{kin}^2 cut value

Fitting function for Exclusive Events Selection Kinematic Variable

$$f(x) = \underbrace{A_L \cdot \frac{(\Gamma/2)^2}{(x - \mu)^2 + (\Gamma/2)^2}}_{\text{Lorentzian core}} + \underbrace{A_G \cdot \exp\left(-\frac{1}{2} \left(\frac{x - \mu}{\sigma_G}\right)^2\right)}_{\text{Gaussian tails}}$$

$$f(x) = \frac{1}{x \sigma_{\ln} \sqrt{2\pi}} \exp\left(-\frac{(\ln x - \mu_{\ln})^2}{2\sigma_{\ln}^2}\right)$$

Where: x : the variable ($\Delta P_\perp, \Delta P_T^{imb}, \Delta \Phi, \& M_{miss}^2$).

$A_L \equiv P[0]$: Amplitude of the Lorentzian core, μ : shared mean,

$\Gamma \equiv P[2]$: FWHM of the Lorentzian,

$A_G \equiv P[5]$: Amplitude of the Gaussian tail & σ_G : width of the Gaussian. μ_{\ln} : mean of $\ln x$ & σ_{\ln} : standard deviation of $\ln x$

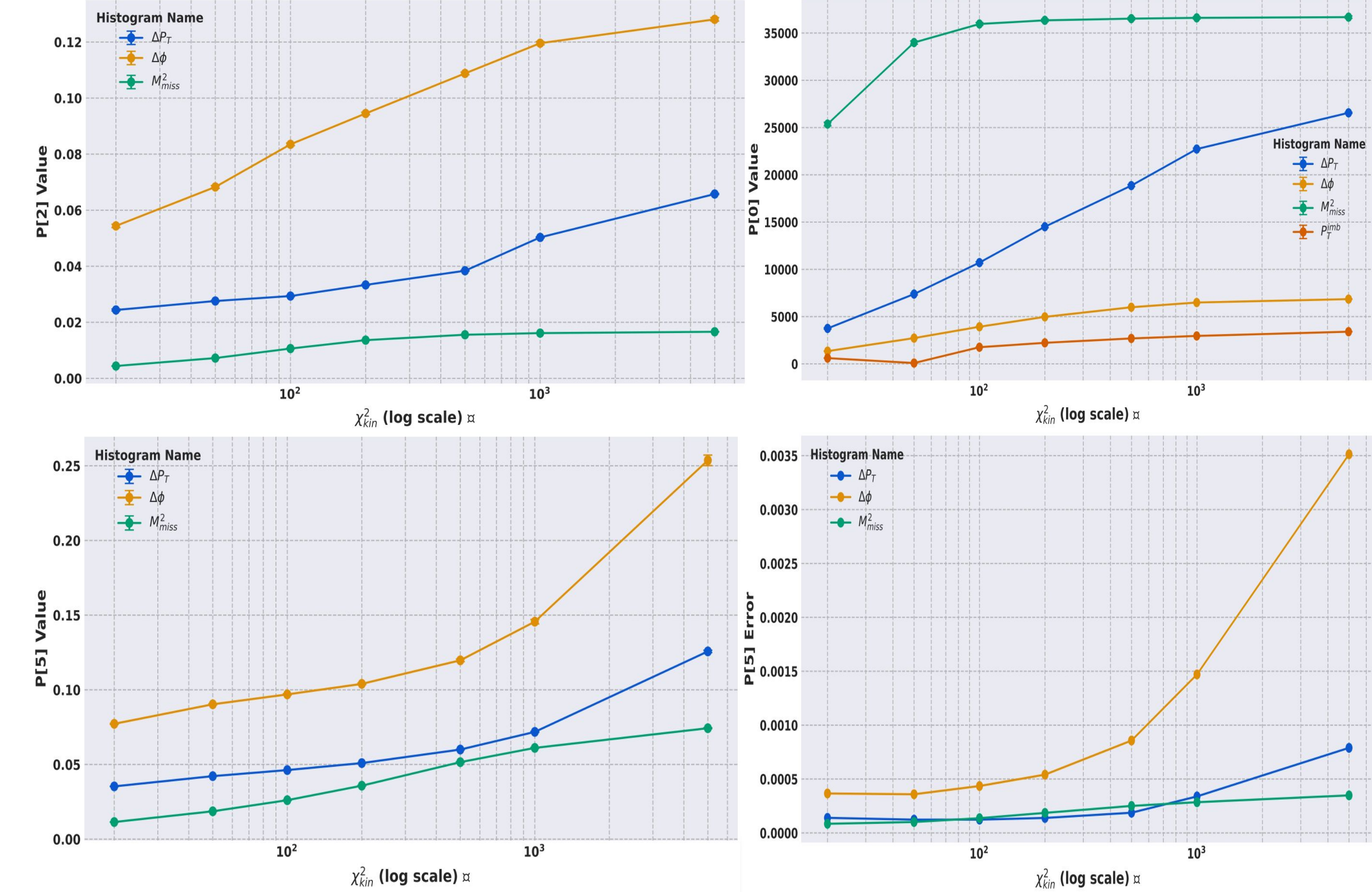


Figure 5. $\Delta P_T, \Delta P_T^{imb}, \Delta \Phi, \& M_{miss}^2$ dependence on χ_{kin}^2

- The above plots shows that the four variables have dependence on χ_{kin}^2
- It indicates that the optimum cut for the χ_{kin}^2 should be 100.0 since the M_{miss}^2 amplitude began to drop after that value.
- At χ_{kin}^2 below 100.0, the rate of change of width of the variables decrease & stabilize.

Outlook

- Impliment the new χ_{kin}^2 cut for optimum event selections.
- Remove regions with very high BH events that suppress TCS events.
- Continue background suppression studies.

Acknowledgment

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References

- [1] P Chatagnon, S Nicolai, S Stepanyan, MJ Amarian, G Angelini, WR Armstrong, H Atac, C Ayerbe Gayoso, NA Baltzell, L Barion, et al. First measurement of timelike compton scattering. *Physical review letters*, 127(26):262501, 2021.
- [2] M Boër, M Guidal, and M Vanderhaeghen. Timelike compton scattering off the proton and generalized parton distributions. *The European Physical Journal A*, 51(8):103, 2015.