#### hadron spectroscopy and QCD

#### Jozef Dudek





hadron spectrum collaboration

### current & future experiments

high statistics novel production



## current & future experiments

high statistics novel production



### where the field was ... meson spectrum

state-of-the-art determinations of the excited meson spectrum in lattice QCD



actually appeared **after** first HSC results)



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## where the field was ... hybrid mesons

in 2009 flux-tube and bag models (1980s) were still state-of-the-art







#### where the field was ... hybrid mesons







#### distillation

efficiently evaluate a large number of correlation functions compute quark annihilation where needed

#### large basis of hadron operators

began with meson operator basis  $\, \bar{\psi} \Gamma \overleftrightarrow{D} \ldots \overleftrightarrow{D} \psi$ 

'subduced' into the irreps of the cubic symmetry

found a workaround for the breakdown of rotational symmetry

(up to three derivatives)

#### variational solution

'diagonalize' a matrix of correlation functions

extract many excited states

s 
$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \, \mathcal{O}_j^{\dagger}(0) | 0 \rangle$$
  
 $C(t)v^{\mathfrak{n}} = \lambda^{\mathfrak{n}}(t) \, C(t_0)v^{\mathfrak{n}}$ 





#### excited isovector meson spectrum





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#### excited isovector meson spectrum



• spectrum does not change qualitatively

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## isovector hybrid mesons

• 'super'-multiplet of hybrid mesons roughly 1.3 GeV above the ho



utilized overlaps with characteristic operators to identify state make-up

• these states have a dominant overlap onto  $\ ar{\psi}\Gamma[D,D]\psi\ \sim [qar{q}]_{m{8}_c}\otimes B_{m{8}_c}$ 



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 $(0,1,2)^{-+},1$ 

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#### exotic hybrid quark mass dependence







#### isoscalar meson spectrum

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## isoscalar meson spectrum

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rest of the lattice community still struggling with  $\eta, \eta'$  alone



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## charmonium

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this work lead by our Dublin collaborators





## chromo-magnetic gluonic excitation

• lightest set of hybrid mesons appear to contain a  $1^{+-}$  gluonic excitation

quarks in  
an S-wave 
$$\begin{bmatrix} q\bar{q}_{\mathbf{8}_{\mathbf{c}}} \begin{bmatrix} {}^{1}\!S_{0} \end{bmatrix} G_{\mathbf{8}_{\mathbf{c}}}^{\star} \begin{bmatrix} B \end{bmatrix} \end{bmatrix}_{\mathbf{1}_{\mathbf{c}}} \to 1_{\text{hyb.}}^{--} \\ \begin{bmatrix} q\bar{q}_{\mathbf{8}_{\mathbf{c}}} \begin{bmatrix} {}^{3}\!S_{1} \end{bmatrix} G_{\mathbf{8}_{\mathbf{c}}}^{\star} \begin{bmatrix} B \end{bmatrix} \end{bmatrix}_{\mathbf{1}_{\mathbf{c}}} \to (0, 1, 2)_{\text{hyb.}}^{-+}$$



- some models have similar systematics
  - bag model also has 1<sup>+-</sup> lowest in energy
  - 1<sup>+-</sup> in a Coulomb-gauge approach





## chromo-magnetic gluonic excitation

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$$\begin{array}{l} \begin{array}{l} \begin{array}{l} \begin{array}{l} \begin{array}{l} \begin{array}{l} \begin{array}{l} \left[ q\bar{q}_{8_{c}} \left[ ^{1}S_{0} \right] G_{8_{c}}^{\star} \left[ B \right] \right]_{1_{c}} \rightarrow 1_{hyb.}^{--} \\ \end{array} \\ \begin{array}{l} \left[ q\bar{q}_{8_{c}} \left[ ^{3}S_{1} \right] G_{8_{c}}^{\star} \left[ B \right] \right]_{1_{c}} \rightarrow (0,1,2)_{hyb.}^{-+} \\ \end{array} \\ \end{array} \\ \begin{array}{l} \begin{array}{l} \begin{array}{l} \begin{array}{l} \begin{array}{l} \left[ q\bar{q}_{8_{c}} \left[ ^{1}P_{1} \right] G_{8_{c}}^{\star} \left[ B \right] \right]_{1_{c}} \rightarrow (0,1,2)_{hyb.}^{++} \\ \end{array} \\ \end{array} \\ \begin{array}{l} \left[ q\bar{q}_{8_{c}} \left[ ^{1}P_{1} \right] G_{8_{c}}^{\star} \left[ B \right] \right]_{1_{c}} \rightarrow (0,1,2)_{hyb.}^{++} \\ \end{array} \\ \end{array} \\ \begin{array}{l} \begin{array}{l} \begin{array}{l} \left[ q\bar{q}_{8_{c}} \left[ ^{1}P_{1} \right] G_{8_{c}}^{\star} \left[ B \right] \right]_{1_{c}} \rightarrow (0,1,2)_{hyb.}^{++} \\ \end{array} \\ \end{array} \\ \begin{array}{l} \left[ q\bar{q}_{8_{c}} \left[ ^{3}P_{0,1,2} \right] G_{8_{c}}^{\star} \left[ B \right] \right]_{1_{c}} \rightarrow (0,1^{3},2^{2},3)_{hyb.}^{+-} \end{array} \end{array} \end{array} \end{array} \end{array} \end{array}$$

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  - bag model also has  $1^{+-}$  lowest in energy
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## excited baryons

• a 'super'-multiplet of hybrid baryons



spectrum from large basis of baryon operators

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$$\epsilon_{abc} \left( D^{n_1} \frac{1}{2} (1 \pm \gamma_0) \psi \right)^a \left( D^{n_2} \frac{1}{2} (1 \pm \gamma_0) \psi \right)^b \left( D^{n_3} \frac{1}{2} (1 \pm \gamma_0) \psi \right)^c$$

PRD84 074508 (2011) PRD85 054016 (2012)



## chromo-magnetic excitation

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lowest gluonic excitation in QCD now determined ?



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## what other properties ?



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## excited spectra summary of progress

# using operators $\overline{\psi}\Gamma\overleftrightarrow{D}\ldots\overleftrightarrow{D}\psi$ and $\epsilon_{abc}\left(D^{n_1}\frac{1}{2}(1\pm\gamma_0)\psi\right)^a\left(D^{n_2}\frac{1}{2}(1\pm\gamma_0)\psi\right)^b\left(D^{n_3}\frac{1}{2}(1\pm\gamma_0)\psi\right)^c$

- obtained unprecedented, and mainly still unmatched, detailed excited state spectra
- isolated the spectrum of hybrid mesons and baryons, both exotic and non-exotic
- inferred the probable nature of the lightest gluonic excitation

(2009 - 2013)

- demonstrated the possibility of extracting current transitions between excited states





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# using operators $\bar{\psi}\Gamma\overleftrightarrow{D}\ldots\overleftrightarrow{D}\psi$ and $\epsilon_{abc}\left(D^{n_1}\frac{1}{2}(1\pm\gamma_0)\psi\right)^a\left(D^{n_2}\frac{1}{2}(1\pm\gamma_0)\psi\right)^b\left(D^{n_3}\frac{1}{2}(1\pm\gamma_0)\psi\right)^c$

- obtained unprecedented, and mainly still unmatched, detailed excited state spectra
- isolated the spectrum of hybrid mesons and baryons, both exotic and non-exotic
- inferred the probable nature of the lightest gluonic excitation

(2009 - 2013)

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- demonstrated the possibility of extracting current transitions between excited states

but this comes close to exhausting what can usefully be learned from using just these operators

we learn a lot from these results, but need to move to methods that can reproduce the true physics of excited states

resonances !





#### excited hadrons are resonances

#### even the simplest excited hadrons are not really states of a definite single mass

PHYSICAL REVIEW D

VOLUME 7, NUMBER 5

1 MARCH 1973

 $\pi\pi$  Partial-Wave Analysis from Reactions  $\pi^+ p \to \pi^+ \pi^- \Delta^{++}$  and  $\pi^+ p \to K^+ K^- \Delta^{++}$  at 7.1 GeV/c<sup>+</sup>

S. D. Protopopescu,\* M. Alston-Garnjost, A. Barbaro-Galtieri, S. M. Flatté,‡

J. H. Friedman, § T. A. Lasinski, G. R. Lynch, M. S. Rabin, || and F. T. Solmitz Lawrence Berkeley Laboratory, University of California, Berkeley, California 94720 (Received 25 September 1972)







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#### excited hadrons are resonances



but how can scattering amplitudes be accessed in lattice QCD?





## one-dim quantum mechanics

• consider scattering of two identical bosons







## 'scattering' in a finite-volume

• consider scattering of two identical bosons



• apply periodic boundary conditions

$$\frac{\psi(-L/2) = \psi(L/2)}{\frac{d\psi}{dz}(-L/2)} \left\{ \frac{d\psi}{dz}(L/2) \right\} \frac{pL}{2} + \delta(p) = n\pi$$



discrete energy spectrum





## 3+1 dim field theory in a cubic volume

known

Lüscher:

 $\cot \delta_{\ell}(E) = \mathcal{M}_{\ell}(E,L)$ 

[ modulo some subtleties regarding *l*-mixing ]



## what operator basis is required ?

supplement large  $\bar{\psi}\Gamma\overleftrightarrow{D}\ldots\overleftrightarrow{D}\psi$  basis with meson-meson-like operators

e.g. 
$$\mathcal{O}_{\pi\pi}^{|\vec{p}|} = \sum_{\hat{p}} C(\hat{p}) \mathcal{O}_{\pi}(\vec{p}) \mathcal{O}_{\pi}(-\vec{p})$$
 where  $\mathcal{O}_{\pi}(\vec{p}) = \sum_{\vec{x}} e^{i\vec{p}\cdot\vec{x}} \bar{\psi} \Gamma \psi(\vec{x})$ 

now need to evaluate diagrams like



**distillation** can handle the annihilation lines





### spectra in moving frames









#### $\pi\pi P$ -wave phase-shift



180
### $\pi\pi P$ -wave phase-shift



### $\pi\pi P$ -wave phase-shift

• reducing the pion mass moves mass, width in the right direction ...



# coupling to external currents

- e.g. GlueX/CLAS12 use photons to excite resonances can this be studied ?
- finite-volume formalism applied to three-point functions





## coupling to external currents

- e.g. GlueX/CLAS12 use photons to excite resonances can this be studied ?
- finite-volume formalism applied to three-point functions
- first explicit application has appeared  $\gamma^{\star}\pi o 
  ho o \pi\pi$  PRL 115 242001 (2015)





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## excited isovector meson spectrum

• what about decays of higher excited states ?

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## coupled-channel meson-meson scattering

• more challenging analysis problem

e.g. in a **two-channel** process, **three** unknowns specify the S-matrix at each energy

our solution: parameterize the energy dependence of the S-matrix and describe the finite-volume spectra by varying parameters





## coupled-channel meson-meson scattering

• more challenging analysis problem

e.g. in a **two-channel** process, **three** unknowns specify the S-matrix at each energy

our solution: parameterize the energy dependence of the S-matrix and describe the finite-volume spectra by varying parameters

- first attempt, coupled-channel  $\pi K/\eta K$  scattering
- need to compute the finite-volume spectra ... lots of Wick contractions ...







### $\pi K/\eta K$ lattice QCD spectra



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### $\pi K/\eta K$ lattice QCD spectra

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### $\pi K/\eta K$ lattice QCD spectra





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## $\pi K/\eta K$ coupled-channel scattering

• describe all the finite-volume spectra

$$\chi^2 / N_{\rm dof} = \frac{49.1}{61 - 6} = 0.89$$





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 $m_{\pi} \sim 391 \,\mathrm{MeV}$ 

### versus experimental scattering

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• S-matrix poles as least model-dependent characterization of resonances



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• S-matrix poles as least model-dependent characterization of resonances



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• S-matrix poles as least model-dependent characterization of resonances





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• S-matrix poles as least model-dependent characterization of resonances



... but no strong channel-coupling here ...



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# $\pi\eta/K\overline{K}$ scattering and the $a_0$ (980)

- sharp experimental enhancement at  $K\overline{K}$  threshold
- usually in 'less-simple' production processes

• amplitude models typically give

e 
$$\frac{g^2(K\overline{K})}{g^2(\pi\eta)} \sim 1$$

e.g.  $p\overline{p} \rightarrow \pi \pi \eta$ 

 $\phi \rightarrow \gamma \pi \eta$ 





# $\pi\eta/K\overline{K}$ scattering and the $a_0$ (980)

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 $\phi \rightarrow \gamma \pi \eta$ 

- sharp experimental enhancement at  $K\overline{K}$  threshold
- usually in 'less-simple' production processes

• amplitude models typically give  $\frac{g^2(KK)}{q^2(\pi n)} \sim 1$ 

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# $\pi \eta / K \overline{K}$ scattering and the $a_0$ (980)

• sharp experimental enhancement at  $K\overline{K}$  threshold

has a 'partner' resonance,  $f_0(980)$ , in  $\pi\pi$  at  $K\overline{K}$  threshold

• often argued that these states don't fit into the  $q\overline{q}$  picture

hidden strange tetraquarks ? *KK* molecules ?





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hidden strange tetraquarks ? *KK* molecules ?

Progress of Theoretical Physics Supplement No. 168, 2007

Can Experiment Distinguish Tetraquark Scalars, Molecules and  $\overline{q}q$  Mesons?

Michael PENNINGTON\*)

Institute for Particle Physics Phenomenology, Durham University, Durham DH1 3LE, UK

#### ... and can theory ?





# $\pi\eta/K\overline{K}/\pi\eta'$ lattice QCD spectra

• used a basis of  $\pi\eta/K\overline{K}/\pi\eta'$  like operators and  $\overline{\psi}\Gamma\overleftrightarrow{D}\ldots\overleftrightarrow{D}\psi$ 





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# $\pi\eta/K\overline{K}$ scattering amplitudes







# $\pi\eta/K\overline{K}$ scattering amplitudes

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 such a strong effect at threshold likely indicates a nearby singularity



## pole singularities in two-channels

• unitarity implies four Riemann sheets in this case



# 'BW-like' resonance below KK threshold





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# 'BW-like' resonance above KK threshold

#### 'Flatté form'

$$t_{ij}(s) = \frac{g_i g_j}{m^2 - s - ig_1^2 \rho_1(s) - ig_2^2 \rho_2(s)}$$







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"a bump in each channel"



fits to experimental data tend to exhibit this structure





### single-pole resonance on sheet IV





# amplitude from lattice spectrum



this fit from a *K*-matrix parameterization

$$t_{ij}^{-1}(E) = K_{ij}^{-1}(E) + \delta_{ij} I_i(E)$$
$$K_{ij}(E) = \frac{g_i g_j}{m^2 - E^2} + \gamma_{ij}$$



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 $m_{\pi} \sim 391 \,\mathrm{MeV}$ 

### singularity structure





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bit esoteric isn't it ... ?



## pole distributions and molecular states ?

Pole counting and resonance classification

D. Morgan Rutherford Appleton Laboratory, Chilton, Didcot, Oxon, OX11 0QX, UK

Received 14 January 1992

'confined' state coupled to decay continuum  $\rightarrow$  Breit-Wigner like (two poles)

molecular state from long-range potential  $\rightarrow$  one pole





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#### other molecule diagnostics ?

couple to an external current e.g.  $\phi \rightarrow \gamma(\pi \eta, K\overline{K})$ or  $(\pi \eta, KK) \rightarrow \gamma(\pi \eta, K\overline{K})$  or other currents ...

and extract form-factors from the residue of the pole





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```
examples of the
interesting convergence of
lattice QCD,
S-matrix ideas,
and phenomenology
```





# highlights

• application of lattice methods has **reinvigorated theoretical hadron spectroscopy** 

A have a tool for phenomenology directly connected to QCD

→ put ideas of gluonic excitations / hybrid hadrons on a firmer footing

→ finite-volume methods give access to scattering amplitudes

 $\rightarrow$  can now really calculate the physical observables

• the Hadron Spectrum Collaboration is the world-leader in these efforts

 $\rightarrow$  nobody else is doing 'spectrum' calculations of this scope




# (near) future

### scattering

scalar mesons at lower  $m_{\pi}$ 

axial and other mesons in e.g.  $\pi\rho$ ,  $\pi\omega$ 

hybrid meson decays at heavier quark masses

simple baryon resonances

more photon transitions

three-body scattering

## phenomenology

tetraquark operator basis

molecules ? pole distributions ? form-factors ?

and haven't mentioned charmonium sector ... XYZ's ...



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## timeline

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2009	dynamical anisotropic lattices, distillation
2010	highly excited isovector meson spectrum
2011	highly excited isoscalar meson spectrum highly excited baryon spectrum phenomenology of hybrid mesons
2012	hybrid baryon spectrum $\pi\pi$ scattering, isospin=2 highly excited charmonium spectrum
2013	$\pi\pi$ scattering, isospin=1, $\rho$ resonance coupled-channel formalism
2014	coupled-channel $\pi K$ , $\eta K$ scattering
2015	excited meson radiative transitions $\gamma\pi \rightarrow \pi\pi$ and the $\rho \rightarrow \pi\gamma$ transition
2016	(to appear) coupled-channel $\pi\eta$ , $K\overline{K}$ scattering

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## hadron spectrum collaboration

JEFFERSON LAB	TRINITY COLLEGE, DUBLIN	CAMBRIDGE UNIVERSITY
Jozef Dudek Robert Edwards	Mike Peardon Sinead Ryan	Christopher Thomas
David Richards	TATA, MUMBAI	U. OF MARYLAND
	Nilmani Mathur	Steve Wallace

## BARYON SPECTRUM

 $\begin{array}{ll} \mbox{PRD84 074508 (2011)} & (N, \Delta)^{\star} \\ \mbox{PRD85 054016 (2012)} & (N, \Delta)_{\rm hyb} \\ \mbox{PRD87 054506 (2013)} & (N \dots \Xi)^{\star} \\ \mbox{PRD90 074504 (2014)} & \Omega^{\star}_{ccc} \\ \mbox{PRD91 094502 (2015)} & \Xi^{\star}_{cc} \end{array}$ 

### HADRON SCATTERING

PRD83 071504 (2011) $\pi\pi I = 2$ PRD86 034031 (2012) $\pi\pi I = 2$ PRD87 034505 (2013) $\pi\pi I = 1, \rho$ PRL113 182001 (2014) $\pi K, \eta K$ PRD91 054008 (2015) $\pi \pi, K \bar{K}$ PRD92 094502 (2015) $\pi\pi, K \bar{K}$ 

#### MESON SPECTRUM

PRL103 262001 (2009)I = 1PRD82 034508 (2010) $I = 1, K^*$ PRD83 111502 (2011)I = 0JHEP07 126 (2011) $c\bar{c}$ PRD88 094505 (2013)I = 0JHEP05 021 (2013) $D, D_s$ 

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#### "TECHNOLOGY"

PRD79 034502 (2009)latticesPRD80 054506 (2009)distillationPRD85 014507 (2012) $\vec{p} > 0$ 

#### MATRIX ELEMENTS

 $\begin{array}{ll} \mbox{PRD91 114501 (2015)} & \mbox{$M' \to \gamma M$} \\ \mbox{PRD90 014511 (2014)} & \mbox{$f_{\pi^{\star}}$} \\ \mbox{PRL115 242001 (2015)} & \mbox{$\gamma^{\star} \pi \to \pi \pi$} \end{array}$ 



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